Cardiac Passive Acoustic Localization: Cardiopal

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Abstract

A novel non-invasive system is proposed as an adjunct diagnostic tool for cardiac disorders, which is based on processing of the heart sounds acquired using a specially designed 2-D passive acoustic array. In addition to the acoustic array, the system consists of a personal computer, specially developed instrumentation and interface hardware and an adaptive array processing scheme. A signal model is constructed in accordance with the basic assumptions made for very near-field low frequency sources. The multiple signal characterization (MUSIC) method together with the model is used for the localization of assumed point sources in the heart. Locations of the sources are estimated by applying 2-D searches.

The effectiveness of the system is demonstrated with extensive recordings at different SNR levels on the phantoms constructed using small sound sources. The system is also tested on human subjects. Different "images" corresponding to the various phases of the heart beat are obtained. It is explained that extraction of 3-D "images" is also possible with the same array data by evaluating a distance parameter relating the sources to the origin of the array.

1. Introduction

The heart, due to its blood pumping action, is one of the most powerful acoustic sources within the body. It emits low frequency vibrations with a rather high dynamic range. Recent experimental and theoretical studies have emphasized the importance of distinct structures of the heart tissue in the genesis of the heart sounds. For example, rapidly changing tension of the myocardium, acceleration and deceleration of blood masses and valvular activities cause multi-component heart sounds during the cardiac cycle. They act as the driving forces on the vibratory modes of the surrounding tissue and the myocardium. Moreover, the turbulent flow due to a regurgitation or a stenosis can cause deterministic vibrations of the vessel walls and the myocardium. When these sounds are recorded on the chest wall, they are also conditioned by the transmission characteristics of the myocardium, the lung and the thoracic tissue [1-7].

Monitoring the entire region of the heart may require many sensors. Vermarien et al. [8] and Okada et al. [9] successfully obtained low-frequency vibratory maps on the chest wall with contact type transducers. Ramachandran et al. [10] used a capacitance type distance meter to reconstruct 2-D patterns of the cardiac activity. During various phases of the ECG, they displayed the out-of-plane displacements caused by major cardiac events in a perspective view. Hök et al. [11] applied a holographic method with pulsed ruby laser to provide a measurement without any contact. Another optical technique based on laser speckle interferometry has been developed by Ramachandran et al. [12]. They analyzed the displacements during P, QRS and T

segments of the ECG. Recently, several investigators have directed their efforts to achieve a spatial mapping of the acoustical energy emitted by the heart. In these investigations, the heart sounds are acquired from certain conventional landmark points on the chest. They are then processed using coherent averaging and interpolating the acoustic energy among the cardiac sensors [1, 4].

In the present study, a novel system was developed for spatial localization and 2-D mapping of the assumed acoustic sources of the heart. The system is called Cardiac Passive Acoustic Localizer or CARDIOPAL, for short. It consists of a specially designed multi-sensor probe in the form of a planar microphone array, precision amplifiers, filters and A/D converters, interface circuitry, a PC and software. With this system different "images" corresponding to the various phases of the heart beat; e.g., closure of the mitral and tricuspid valves, ejection of blood flow in the systole, closure of the aortic and pulmonary valves, early and late diastole, were obtained. The details of the system and the method used are described in the following sections.

2. The 2-D Sensor Array and the Instrumentation System

The multi-sensor probe developed is a 105 mm x 105 mm x 40 mm planar acoustic array made with 16 miniature electret microphones mounted on a 4 x 4 grid. A soft *polyurethane* layer was used to couple the rigid surface of the collimator of the array to the chest wall. The data acquisition hardware consists of balanced microphone amplifiers, 80-1000 Hz band-pass analog filters, Analog-to-Digital Converters (ADC) and an interface to a PC. During data acquisitions all microphone channels are sampled *simultaneously* and quantized using 12-bits and the acquired data is transferred to the PC via Direct-Memory-Access (DMA).

In the present study, from the signal processing point of view, it is crucial to have an instrumentation system for the array with sufficiently low amplifier noise, thus an adequate dynamic range, and a high enough Common-Mode-Rejection-Ratio (CMRR). The present system has an average noise factor of 1.0232 and 51.6 dB dynamic range (measured when an acoustically isolated microphone is coupled to a channel) and average 86 dB (at 50 Hz) CMRR. These specifications of the instrumentation system are quite sufficient for accurate measurement of the heart sounds. Further details about the microphone array and the instrumentation system are given by Bahadırlar [13].

3. Directional Response of the Array

The construction of the microphone array effectively limits the omnidirectional response of the air-coupled microphones and confines its sensitive region to the orifice of the coupling layer. However, the quiescent directional response of the array is not suitable to localize vibrating structures within the heart. In other words, one cannot differentiate these sites merely by steering the mainlobe of the microphone array. This is due to the fact that the mainlobe is almost semi-spherical on the chest side, even for the upper frequency of the instrumentation system (1000 Hz). The dimensions of the array are very small and the frequency of interest is rather low [13]. Fortunately, with the help of advanced array processing techniques, the degrees of freedom of the array can be utilized effectively to obtain spatial information beyond that obtained from the quiescent directional response of the array [14]. Furthermore, by making judicious assumptions about the hypothetical acoustic sources and about the propagation of the acoustic energy emitted from the heart tissue to the array, spatial differentiation can be achieved.

4. Assumptions about the Heart Sounds

The genesis and propagation of the heart sounds through the transmission system from the heart to the chest surface is a very complex phenomenon and the mechanisms involved are still controversial. The oldest theory, called the valvular theory states that the heart sounds are actually produced by the opening and the closure of the four cardiac valves [1]. Another widely accepted cardiohemic theory relates the heart sounds to the vibration of the entire cardiac structures acting as a single body [1, 2]. The theory contains a mechanism for the genesis of the heart sounds, incompatible with the valvular theory since it emphasizes the importance of the whole complex of the heart tissues, and acceleration and deceleration of the blood in the heart. The cardiohemic theory excludes the valve vibration as a cause of the heart sounds. Combining the observations of the previous theories, a modern theory named the multi-degree of freedom theory states that intracardiac sounds originate from distinct cardiac structures [1]. This theory hypothesizes that the cardiac structures differs in morphology, therefore the individual structures contribute to the heart sounds on the basis of their different vibratory modes. Although the valvular and the cardiohemic theories are partly valid, recent investigations strongly support the multi-degree of freedom theory [1, 3, 7].

The heart-thorax acoustic system acts as a slowly time-invariant low-pass filter. Therefore the intracardiac sounds are also conditioned by this acoustic system. In addition, the frequency content of the sounds on the chest surface does not change significantly between different sites on the chest, whereas the intensity changes considerably and characteristically [1, 7]. Bearing in mind these findings and the multi-degree of freedom theory, the following assumptions for the production and propagation of the heart sounds can be adopted to form a signal model to be occupied in our array processing approach:

- The vibrating tissue is small compared with the size of the microphone array and acts as a point source. Therefore, the sound wave from the tissue is propagated as a spherical wave in this very near field.
- There is no scattering or reflection of the sound wave within the tissue between the heart and the chest surface.
- The intervening tissue is homogeneous in terms of sound absorption characteristics. The sound velocity is assumed to be constant and can be taken as c = 1530 m/sec.
- The sound sources are stationary within the observation interval of interest.

Considering the difficulty in establishing of a strong theory for the genesis and propagation of the heart sounds, the assumptions may be viewed as controversial. Nevertheless, they will promote the array processing scheme described in the next section and the performance of the method will be discussed in sections 6 and 7 with this signal model and the array data acquired from the human subjects.

5. The Signal Model and Acoustic Source Localization

The MUSIC algorithm is a subspace-based estimation technique in array signal processing and relies on the geometrical properties of an assumed signal model [15, 16]. If the model accurately represents the measured data, the algorithm can result in "superresolution"; a capability not limited by the array size, but by the data collection time and/or the SNR level of the measured data.



Figure 1. A point source in the near-field of the 4x4-element planar array.

In the present study, considering the special geometry of the microphone array, a signal model was defined in accordance with the assumptions given in Section 4. Figure 1 shows a point source in the near field of the array. Its coordinates and its range to the array origin are used as the principal parameters of the signal model in Eq. (1) given below. Assuming that there are K point sources in the near field of the array, the model for the signal $X_{mn}(t)$ at the microphone in location (m, n) can be written as

$$X_{mn}(t) = \sum_{k=1}^{K} S_k(t) a_{mn}(t, r_{ok}, \theta_k, \phi_k) + n_{mn}(t), \ 1 \le m \le M, 1 \le n \le N;$$
(1)

where $S_k(t)$ is the amplitude of the k'th sound source, $a_{mn}(t, r_{ok}, \theta_k, \phi_k)$ is the array steering parameter depending on the location of the k'th source, $n_{mn}(t)$ is spatially white Gaussian noise, M and N are the number of microphones in the x and the y axes, respectively. Assuming the near-field sources to be stable in their locations during the observation time, the array steering parameter can be given as:

$$a_{mn}(r_{ok},\theta_k,\phi_k) = \left(\frac{r_{ok}}{d_{mn}(r_{ok},\theta_k,\phi_k)}\right)^2$$

$$\exp\left[-(ad_m(r_k,\theta_k,\phi_k)/r_k + ik_0(r_k,-d_m(r_k,\theta_k,\phi_k)))\right]$$
(2)

$$\left[\left(p_{mn}(r, \delta_{k}, \delta_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \delta_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}) \right) \right] = \left[\left(p_{mn}(r, \delta_{k}, \phi_{k}) \right] = \left[\left(p_{$$

$$d_{mn}(r_{ok},\theta_k,\phi_k) = \left[r_{ok}^2 + x_m^2 + y_n^2 - 2r_{ok}\sin\theta_k(x_m\cos\theta_k + y_n\sin\phi_k) \right]^{1/2}$$
(3)

$$x_m = [m - (M+1)/2] d_x, \quad y_n = [n - (N+1)/2] d_y, \quad \kappa_0 = 2\pi f_0/(c/d_x), \ d_x = d_y, \quad k_0 = 2\pi f_0/(c/d_x), \ d_x = d_y, \quad k_0 = 2\pi f_0/(c/d_x), \ d_x = d_y, \quad k_0 = 2\pi f_0/(c/d_x), \ d_y =$$

where r_{ok} is the distance between the k'th source and the array origin, $d_{mn}(r_{ok}, \theta_k, \phi_k)$ is the distance from the same source to the (m, n)'th array element, θ_k and ϕ_k are the elevation and azimuth angles, respectively (see Figure 1), κ_0 is the wave number depending on the frequency of the source (f_0) and on the speed of the sound (c) in the medium, d_x and d_y are the inter-element distances in the respective axes on the array and ρ is the exponential loss factor of the medium.

The steering parameter $a_{mn}(r_{ok}, \theta_k, \phi_k)$ consists of three factors, two of which correspond to the assumptions given in the preceding section and the other is due to the delay/advance of the wave with respect to the unmeasured wave at the array center. The first factor represents the intensity loss due to the

spherical wave expanding about a point source. The intensity level is inversely proportional to the second power of the distance from the source. This is due to the spherical surface of the symmetrically expanding wave. The second factor is due to the energy loss per unit volume in a lossy homogenous medium. If the loss is a constant fraction of the wave energy, the intensity loss caused by the medium is an exponential function of the distance to the source [17, 18].

The array output can be written as an $M \times N$ column vector \mathbf{X} (t) from (1) by arranging the outputs of the rows one after the other. In this case the $MN \times MN$ array output covariance matrix can be given as

$$\mathbf{R} = \mathbf{E} \left[\mathbf{X}(t) \mathbf{X}^{\mathrm{H}}(t) \right] \tag{4}$$

where $(\cdot)^{H}$ represents the complex conjugate transpose (Hermitian). If we assume the spatial noise to be statistically independent among the channels and from the source signals, the covariance matrix can be written as

$$\mathbf{R} = \mathbf{A}\mathbf{R}_{\mathbf{S}}\mathbf{A}^{\mathrm{H}} + \sigma^{2}\mathbf{I}.$$
 (5)

where $\mathbf{R}_{\mathbf{S}}$ represents the source covariance matrix of size $K \times K$, \mathbf{A} is the steering matrix of $MN \times K$ elements correspondingly defined by the steering parameters $a_{mn}(r_{ok}, \theta_k, \phi_k)$, σ^2 is the variance of the spatial noise and \mathbf{I} is the identity matrix. If the sources are incoherent, then $\mathbf{R}_{\mathbf{S}}$ remains nonsingular and has rank K. \mathbf{A} is also of rank K due to the linear independence of its columns. Therefore this implies that the $\mathbf{A}\mathbf{R}_{\mathbf{S}}\mathbf{A}^{\mathrm{H}}$ is also of rank K.

Let $\lambda_1 \geq \lambda_2 \geq \ldots \geq \lambda_{MN}$ be the ordered eigenvalues of the array covariance matrix **R**. They are given by

$$\lambda_k = \begin{cases} \mu_k + \sigma^2 & k = 1, 2, \dots K \\ \sigma^2 & k = K + 1, K + 2, \dots, MN, \end{cases}$$
(6)

where μ_k is the eigenvalue of the source in the noiseless case. For any k > K, the following can be written from (6):

$$\mathbf{R}\mathbf{v}_{k} = \lambda_{k}\mathbf{v}_{k} = \sigma^{2}\mathbf{v}_{k}$$
$$\mathbf{R}\mathbf{v}_{k} = (\mathbf{A}\mathbf{R}_{S}\mathbf{A}^{\mathrm{H}} + \sigma^{2}\mathbf{I})\mathbf{v}_{k}$$
(7)

where \mathbf{v}_k is the eigenvector corresponding to the eigenvalue, λ_k of the covariance matrix. It is seen from (7) that the term $\mathbf{AR}_S \mathbf{A}^{\mathrm{H}} \mathbf{v}_k$ is equal to zero. Since \mathbf{R}_S is of rank K and \mathbf{A} has full column rank K, it follows from (7) that

$$\mathbf{A}^{\mathrm{H}}\mathbf{v}_{k} = 0, \quad K+1 \le k \le MN.$$

$$\tag{8}$$

A fundamental property of the eigenvectors of a correlation matrix is that they are orthogonal to each other. Hence the K eigenvectors span a subspace that is the orthogonal complement of the space spanned by the remaining MN - K eigenvectors. Therefore, it follows from (8) that the MN - K eigenvectors related to the noise subspace are orthogonal to the K columns of the steering matrix **A**, and the columns are related to the sources in the scene. Thus the peaks of the following function will correspond to the true locations of the sources:

$$P(r,\theta,\phi) = \frac{1}{\sum_{k=K+1}^{MN} \left| \mathbf{v}_k^{\mathrm{H}} \mathbf{a}(r,\theta,\phi) \right|^2}$$
(9)

where $\mathbf{a}(r, \theta, \phi)$ is the steering column vector of size MN and r represents the distance from a source to the array center. Clearly, the parameters of the function in (9) are the spherical coordinates and can be used to describe the location of a sound source in a volume, thus the classical MUSIC algorithm extends to a three-dimensional version in the present study [19-22]. Obviously, the function $P(r, \theta, \phi)$ is meaningful when $MN \ge K + 1$ and this limits the minimum number of microphones to resolve the total number of sources in a scene.

When the array covariance matrix \mathbf{R} is known exactly and the assumptions about the noise statistics hold, the peaks of the $P(r, \theta, \phi)$ are guaranteed to correspond to the true locations. Since the peaks or the eigenvectors are always distinct, this estimator can distinguish arbitrarily close acoustic sources. However, a low SNR in the measured signal and any mismatch between the model in (1) and the measured signal from the scene, seriously degrade the resolution of the estimator [15, 16, 19, 20]. In practice, the array covariance matrix \mathbf{R} can be estimated from the ensemble average of the array outputs as

$$\hat{\mathbf{R}} = \frac{1}{L} \sum_{i=1}^{L} \mathbf{X}(t_i) \mathbf{X}^{\mathrm{H}}(t_i)$$
(10)

where L is the number of the samples and $\mathbf{X}(t_i)$ is the array output vector obtained at the *i'th* sample. The K principal and MN - K noise eigenvectors (i.e., the signal and noise subspaces) will be determined from the estimated covariance matrix. The accuracy, with which the number of principal eigenvectors are determined, and the quality of the estimated subspace are the crucial factors for the success of the MUSIC method.

The quality of the estimated signal subspace increases with an increase in the number of samples. Similarly, the estimation bias of the source location is inversely proportional to the number of samples; $(\propto L^{-1})$. The increase in the noise variance also causes an increase in the bias; $(\propto \sigma^2)$, and the estimation variance has order L^{-2} dependence to the samples. Furthermore, the bias is strongly dependent on the separation of the close sources. Hence, two close sources may not be resolved if the estimation bias is high [15, 23]. In the asymptotic case (i.e., the number of samples is very large), for any number of sources, the estimation variance is very close to the Cramer-Rao lower bound as the noise variance approaches zero, i.e., as the SNR increases [23].

6. Experimental Study and Results

For testing the array data acquisition system and proposed source localization method experimental data was acquired. A phantom with three sound sources (small headphone capsules) was constructed to simulate the point sources and to create various scenarios in a relatively controllable environment. The small capsule which has a 2 mm orifice behaves like omnidirectional radiators due to the low frequency of interest (around 500 Hz) in the experiments. The aperture of a capsule is very small compared with the wavelength of the sound waves ($\lambda_{air} = 684 \text{ mm}$ at c = 342 m/s and f = 500 Hz). The angle of divergence of the waves depends on the sine of the ratio between the wavelength and half of the aperture [18]. In our case, obviously, the angle of divergence is very large, so that the capsules do not have any directional responses and emit spherical waves. The sound signals emitted by these sources were recorded in a silent room. Also, the same phantom was placed into a plastic box which was filled with a boiled *gelatin-water mixture* to mimic the tissue between the heart and the chest wall. Similar signal measurements were performed on the outer surface of the plastic box. The heart sounds were also recorded from human subjects while the subjects were supine and asked to hold their breaths. All measurements were performed at 4 kHz sampling frequency. As stated above the success of the acoustical source localization method totally depends on the SNR level of the measured array data, the stability of the sound producing sites during the observations and matching of the signal propagation model with the signals impinging onto the array. Therefore, the experimental data was used to evaluate the method in real circumstances.

To get array signals with various SNR levels, two sources were driven by 495 Hz and 505 Hz sinusoidal signals and the intensities were changed for each SNR. The frequencies of the sources chosen were slightly different to prevent possible coherence. The array signals were passed through a 400-600 Hz band-pass digital filter to intentionally confine the background noise into a bandwidth similar to that of the weak diastolic signals carrying diagnostic value. The SNR level was defined as

$$SNR(dB) = 10 \log \left[\frac{\sum_{i=1}^{L} \left[\mathbf{X}_{S}(t_{i}) - \bar{\mathbf{X}}_{S} \right]^{T} \left[\mathbf{X}_{S}(t_{i}) - \bar{\mathbf{X}}_{S} \right]}{\sum_{i=1}^{L} \left[\mathbf{X}_{N}(t_{i}) - \bar{\mathbf{X}}_{N} \right]^{T} \left[\mathbf{X}_{N}(t_{i}) - \bar{\mathbf{X}}_{N} \right]} \right]$$
(11)

where L is the number of samples, $\mathbf{X}_{S}(t_{i})$ was the array output vector sampled at *i*'th sample when the sources were active, $\mathbf{\bar{X}}_{S}$ was the mean vector of $\mathbf{X}_{S}(t_{i})$ calculated over L samples, and $\mathbf{X}_{N}(t_{i})$ was the array output vector sampled at *i*'th sample when the array output was only due to the instrumentation and the ambient noise, $\mathbf{\bar{X}}_{N}$ was the mean vector of $\mathbf{X}_{N}(t_{i})$ calculated over L samples. The SNR was estimated using L = 4000 samples.

The array signals were then segmented into L samples for each SNR level and the array covariance matrix in (10) was calculated over these samples. The MUSIC algorithm was applied by choosing the number of sources as K = 2. A source plane was defined on a 40x40 grid (with 2.5 mm x 2.5 mm grid size) at 5 cm distal to the array plane, using the parameter r in (9). The distance between the source plane and the array plane is the actual distance in the measurement setup. The speed of sound (c) in air was chosen as 342 m/sec and the exponential loss factor of the medium (ρ) was found by trial and error as approximately 0.45, using a signal with a moderate SNR level.



(a) (b) Figure 2. The estimates of two near-field sound sources at (a) 1.9 dB and (b) 27.3 dB SNR and using L=400 snapshots. The sources are 5 cm distal to the array plane.

Figure 2 presents the two estimates at SNR levels of 1.9 dB and 27.3 dB. The estimates clearly demonstrate that an increasing purity in the signals helps the MUSIC estimator to resolve the acoustic sources easily with this number of samples. Figure 3 shows the dependence of the estimates on the number of samples at a SNR level of 6.6 dB (L was assigned to 400 and 1600 samples, respectively). In the case of 400 samples, the estimator does not resolve the low-power source whereas at 1600 samples the MUSIC gives a good clue to the locations at this impurity level.



Figure 3. The estimates of two near-field sources at 6.6 dB SNR and using (a) L=400 and (b) L=1600 snapshots. The sources are 5 cm distal to the array plane.



Figure 4. Increasing the size of signal-subspace from K=2 to K=5 can improve the spatial spectrum at low SNR (1.9 dB).



Figure 5. Three factors in the array steering vector contribute with various degrees to source localization process in air: (a) Only the phase-shift factor, (b) only the absorption loss and (c) only the spherical spreading factor included in the steering vector. (d) all factors are included. (K=2, L=800, SNR=20.2 dB).

If the array covariance matrix is estimated from the real measurements, the ordered eigenvalues are all different with probability one [24]. In general they decrease monotonously to a minimum value, thus making it difficult to determine the size of the signal- or noise-subspaces [24, 25]. In the case where the SNR level is low, the condition may be more severe, because the high power noise influences the quality of the signal- and noise-subspaces and the condition in (8) may be violated. In this case, an overestimation in the size (K) of signal-subspace may help the MUSIC estimator to improve its resolving capability, since the estimator uses a uniform average of the norms of the projection of the signal-subspace results in spurious signal sources on the spatial spectra. Figure 4 shows the results of the MUSIC estimator in a sample size (L=1600) and at a low SNR (1.9 dB) when the size of the signal-subspace is increased to K = 5. Although the estimator

cannot indicate the low-power source in the case K = 2, it can resolve the two sources in the same sample number when the subspace size is K = 5. An improvement may still be obtained by increasing the number of samples to a point determined by the SNR level.

Figure 5 presents the contribution of each factor in the streering vector, $\mathbf{a}(r, \theta, \phi)$ to the estimates in air at an SNR level of 20.2 dB and L = 800. Once again, the factors are due to the time delay/advance (or phase-shift), spherical-spreading loss and absorption loss. Figure 5.(a) demonstrates that the phase-shift factor contributes less to the estimate in the scenario. Figure 5.(b) gives the estimate which is obtained by retaining only the absorption loss factor in the vector. As verified by a better estimate in Figure 5.(c) the spherical spreading loss dominates the other two factors in the near-field. Figure 5.(c) also confirms that the small headphone capsules possess the omnidirectional effect. Figure 5.(d) gives an augmented and the best estimate where the steering vector consists of all three factors.

Figure 6 shows the estimates when all three sources are activated. Although two close sources are not resolved in Figure 6.(a) the flat form of the larger peak seems to indicate the existence of the two sources. This form is usually observed in other scenarios when the number of samples is small at a high SNR level or when the SNR level is relatively low. Figure 6.(b) shows a gray-scaled version of the estimate in Figure 6.(a). This presentation may provide an easy inspection when the unknown scenarios are estimated by the present method.



(a) (b) Figure 6. The estimate and its gray-scaled version when three soruces are activated. The sources are 5 cm distal to the array plane and K=5, L=3200, SNR=23.2 dB.

Figure 7 presents the MUSIC estimates from the gelatin phantom at high SNR levels. The sources in the phantom are 7.5 cm distal to the array plane. The speed of sound is taken as the speed in water; c=1530 m/sec and the absorption loss is determined by trial and error as approximately $\rho = 2$. In this case the multi-path arrivals are stronger since the size of the plastic box is small compared the array size and the source distance, and the acoustic impedances of the gelatin, the plastic and air are quite different from each other. In contrast to the signals recorded from the phantom in air, it is observed that the signals from the gelatin phantom do not exhibit a systematic variation among the channels of the array. This demonstrates that strong echoes in the phantom complicate the scenario. Hence, the sources could not be resolved by the estimator, in spite of a relatively high SNR level. However, the estimate in Figure 7.(b) indicates the existence of two sources.



(a) (b) Figure 7. The estimates from the gelatin phantom at K=7, L=3200, and (a) at SNR=14.9 dB, (b) at SNR=19.9 dB.

Figure 8 shows the sound localization estimates during the various phases of the heart beats from a healthy subject. Part (f) of the Figure correlates the estimates to the time interval of a representative ECG signal. The heart sounds recorded by the microphone array are segmented by considering the main physiological activities at each heart beat. These are the first heart sound (S1), the systole, the second heart sound (S2), early diastole (S3, if it exists) and late diastole (S4, if it exists). The segment length (the number of samples) in the estimates is L = 500. The array covariance matrix for each phase of the heart beat is calculated and averaged over the segments from 4 - 5 available beats to increase the quality of the estimates. This may seem to violate our previous assumption that the sound sources must be stationary during the observation interval since it is well known that the dynamic activities in a phase of the cardiac cycle causes the heart sounds to be nonstationary even in relatively short periods. Nevertheless, in the present study, possible sound foci are assumed to be stationary during each phase of the beats. The steering parameter in the signal model includes a single frequency because of the time delay/advance factor. Therefore, the segments passed through a narrow band filter with corner frequency at 165 Hz and 235 Hz (70 Hz band width) to obtain nearly single frequency heart sound signals, and to practically conform to the single frequency property of the signal model. The filter is the FIR digital filter with 65 tap weights and constructed using the Chebychev window with the ripple level of -40dB at the stop-band region.

The dimensions of the search plane in Figure 8 are assigned for 125 mm x 125 mm for all estimates. Four parameters in the method are then chosen to extract the possible sources in each phase of the heart beat:

- The speed of sound is assigned to the speed in water; c = 1530 m/sec.
- The absorption loss is determined by trial and error approximately as $\rho = 8$.
- Number of sources in the scene is determined by trial and error approximately as K = 5.
- The distance of the search plane to the array is changed to get various planes of the estimates.







(c)





(e)

(f)

Figure 8. The sound localization estimates of a healthy subject and the electrocardiogram (ECG) showing the various phases of the heart beat: (a) the S1, (b) the systole, (c) the S2, (d) early diastole, (e) late diastole, (f) the symbolic ECG signal. (K=5, L=500, bandwidth= 165-235 Hz)

BAHADIRLAR, GÜLÇÜR: Cardiac Passive Acoustic Localization: Cardiopal

The single absorption loss factor in the present study originates from the fundamental assumption which assumes the intervening tissues between the heart and the chest surface are homogeneous in terms of sound absorption characteristics. However, the complex tissue structures make it difficult to extract a single factor representing the absorption characteristic of the tissues, and a quantitative measurement of this factor may need detailed in vitro studies. Therefore, in the present study the absorption loss is determined by trial and error. It is observed that increasing the loss shrinks the features recognized on the estimates and decreasing the value of the factor causes the features to widen together with a smoothing effect. An extreme change in the factor causes all existing features to disappear.

Many observations in the study show that ordered eigenvalues of the array covariance matrices, which correspond to the different phases of the heart beats, do not exhibit any distinct jump, rather they decrease monotonously to a small value. This basically implies that there are many sound sources in the scene, buried in the noise. This also does not conflict with the multi-degree of freedom theory accepted as a fundamental conception in the study. Increasing the number of sources to the maximum value simply gives distinct and sharp sound foci in the estimates, whereas a decreased value of K results in smoothed and widened sources which seem to be more easily correlated to the activities included in a cardiac phase under inspection. Therefore, the absorption loss and the number of sources are determined by making many observations on the various phases of the beats and by considering the activities in the corresponding phases. The most useful values given above are kept constant in the estimates of all cardiac phases.

The first heart sound, during which many cardiac activities take place in the heart, has the most complicated nature among the other heart sounds. Because of this, the estimate for the first heart sound in Figure 8.(a) is quite difficult to interpret. However, the higher acoustic energy in the figure appears to represent the closure of the mitral valve and early systolic activity in the left heart. Figure 8.(b) presents the estimate during the ejection of blood from the left ventricle to the aorta, i.e. the systolic phase. This figure exhibits higher acoustic energy on the left ventricle side and lower energy towards the aortic side. The acoustic energy due to the closure of the pulmonary and aortic semilunar valves during the end systole and the second heart sound are clearly seen as two foci in Figure 8.(c). During early diastole the rush of the blood from the left atrium to the left ventricle may be inferred from Figure 8.(d). Figure 8.(e) shows the acoustic energy during the active filling of the left ventricle by the contraction of the left atrium.

Normally, the vibrating sites of the heart do not exhibit very small and discrete sounding foci like ideal point sources, rather they are continuously distributed in a volume of heart tissue. Therefore, one can expect that the present array data and signal processing method can give a profile of the sounding volume projected onto the search plane and the profile may be considered to be composed of very close acoustic sources, not resolved due to the limits of the resolution. From this point of view, the estimate in Figure 8 can be regarded as a kind of "image" of the vibrating regions, smoothed by the system resolution. As mentioned above, the heart can be assumed to be a system with a multi-degree of freedom in which the heart sounds are composed of various vibrating structures having different frequency bands [1]. Thus, the estimates in Figure 8 are related to narrow-band vibrations. These ideas lead to the careful interpretations of Figure 8 when the underlying events in the heart are considered.

Results similar to those in Figure 8 were also obtained for two other human subjects; one healthy and the other diseased, with a weak aortic stenosis and widened aortic root, which were determined by echocardiography examination. Also, the estimates corresponding to the first and the second heart sounds of a prenatal baby (30 weeks old) were achieved. The closure of the four cardiac valves could be demonstrated by utilizing the different estimates with increasing distances of the search plane.



Figure 9. The second heart sound estimates from another healthy subject, with respect to the distance from the array plane: (a) 4 cm, (b) 5 cm, (c) 6 cm, (d) 7 cm, (e) 8 cm, (f) 9 cm. (K=5, L=500, bandwidth=100-400 Hz).

(f)

256

(e)

Figure 9 shows the estimates for the second heart sound (S2) of the other healthy subject, depending on increasing distances from the array plane. The sound energy seen on the left side of the estimates in the figure represents the closure of the pulmonary valve and the energy on the right shows the closure of the aortic valve. The figure indicates that the pulmonary valve comes closer to the array plane than the aortic valve. These findings do not conflict with the anatomical placements of the related structures. The estimates far from the array plane show a blurred version of the sound sources existing near the array as given in Figure 9.(e) and Figure 9.(f). This out-of-plane smoothing effect is also observed in the experiment with the phantoms when the search plane is placed at a distance other than the real distance of the sound sources of the phantoms.

These conclude the principal results obtained from the measured signals. Details and further results were given by Bahadırlar [13].

7. Conclusion

In this study, an array signal processing method is proposed for multi-channel heart sound signals. The method is totally original in this area and provides acoustic "images" of the heart corresponding to the various phases of the heart beat; e.g., closure of the mitral and tricuspid valves, ejection of blood flow during systole, closure of the aortic and pulmonary valves, early and late diastole, etc.

Because of physical limitations, the size of the array used is very small compared with the wavelength of the acoustic waves propagating through the thorax. This makes the directional response of the array quite wide and therefore the localization of the acoustic sources within the chest difficult. In spite of these difficulties, the proposed method has been successful in the localizations due to the use of "superresolution" array processing techniques together with a judicious choice of the signal model.

The performance of the proposed method depends totally on the amount of uncertainty in the data, i.e., a high SNR in the recorded signal is the key factor for improved resolution. Naturally, the resolution also depends on the number of sensors in the array and the quantization depth used in the analog to digital conversion process. Therefore, a system with more sensors, very low instrumentation noise and a higher quantization depth will have a better overall resolution. A further improvement in the resolution is possible if a flexible array is used instead of a rigid array. In a rigid array, some of the sensors may not be tightly coupled to the surface of the chest wall, if the contour of the array does not match closely that of the chest. This may seriously degrade the resolution. In a flexible array the coupling of the sensors could be much better. However, the flexible array must be equipped with the means for monitoring the positions of the individual sensors with respect to the reference coordinates used in the processing.

The system can also be used to obtain 3-D acoustic "images" of the heart, since the signal model with the streering vector defines a three dimensional volume in front of the array in which the localization process can be performed using a single estimate of the array covariance matrix and using 3-D search. In this sense, the 3-D estimation is a natural extension of the study. However, the 3-D estimation needs some additional, surface rendering and shading processes to depict the vibrating sites like physical objects. In addition interfering components from the off-place sources should be eliminated. Basic normalizing and thresholding operations defined over the search volume may be used to remove these interferences.

The processing can also be made independent of frequency by neglecting the minor contribution of the phase factor in the model, so that the method can directly be used for wide-band estimates of the heart sound signals.

This new imaging modality offered by the proposed system makes it possible to view the heart sounds

from a new perspective. The method, visualizing the sound foci of the human heart, employing a depth parameter and relating them to specific frequency bands, can yield a huge amount of valuable, diagnostic information from the array data recorded in a short time from a human subject. Our limited experience with the system shows that it has a good potential of becoming a useful tool for the diagnosis and prescreening of the cardiac disorders. The potential of this totally noninvasive and safe method must be further exploited through extensive comparative clinical investigations.

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