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Asymptotic for a second-order evolution equation with convex potential and vanishing damping term

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Abstract: In this short note, we recover by a different method the new result due to Attouch, Chbani, Peyrouqet, and Redont concerning the weak convergence as $t \to +\infty$ of solutions x(t) to the second-order differential equation

$$x''(t) + \frac{K}{t}x'(t) + \nabla\Phi(x(t)) = 0,$$

where K > 3 and Φ is a smooth convex function defined on a Hilbert space \mathcal{H} . Moreover, we improve their result on the rate of convergence of $\Phi(x(t)) - \min \Phi$.

Key words: Dynamical systems, asymptotically small dissipation, asymptotic behavior, energy function, convex function, convex optimization

1. Introduction and statement of the result

Let \mathcal{H} be a real Hilbert space with inner product and norm respectively denoted by $\langle .,. \rangle$ and $\|.\|$. In a very recent work [1], Attouch et al. considered the following second-order differential equation:

$$x''(t) + \gamma(t)x'(t) + \nabla\Phi(x(t)) = 0, \tag{1.1}$$

where $\gamma(t) = \frac{K}{t}$ with K as a nonnegative constant and $\Phi : \mathcal{H} \to \mathbb{R}$ is a convex continuously differentiable function. By developing a method due to Su et al. [5], they proved the following result:

Theorem 1.1 (Attouch, Chbani, Peypouquet, and Redont) Assume that K > 3 and the set $\arg \min \Phi \equiv \{x \in \mathcal{H} : \Phi(x) \leq \Phi(y) \ \forall y \in \mathcal{H}\}$ is nonempty. Let $x : [t_0, +\infty[\to \mathcal{H}]$ be a solution to (1.1). Then x(t) converges weakly in \mathcal{H} as $t \to +\infty$ to some element of $\arg \min \Phi$. Moreover, the energy function

$$W(t) \equiv \frac{1}{2} \|x'(t)\|^2 + \Phi(x(t)) - \min \Phi$$
 (1.2)

satisfies $W(t) = O(t^{-2})$ as $t \to +\infty$.

In this note, we establish, by using a different method, a slightly improved version of the previous theorem. Precisely, we prove the following result.

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Theorem 1.2 Assume that K > 3 and $\arg \min \Phi \neq \emptyset$. Let $x : [t_0, +\infty[\to \mathcal{H} \text{ be a solution to } (1.1)]$. Then x(t) converges weakly in \mathcal{H} as $t \to +\infty$ to some element of $\arg \min \Phi$. Moreover, $W(t) = \circ(t^{-2})$ as $t \to +\infty$.

Remark 1.1 In [3], we studied the asymptotic behavior as $t \to +\infty$ of the solution to Equation (1.1) when the damping term $\gamma(t)$ behaves, for t large enough, like $\frac{K}{t^{\alpha}}$ with K > 0 and $\alpha \in [0,1[$. We proved that if $\arg \min \Phi \neq \emptyset$ then every solution to (1.1) converges weakly in \mathcal{H} to some element of $\arg \min \Phi$. Hence, Theorem 1.1 and Theorem 1.2 extend this result to the limit case corresponding to $\alpha = 1$.

2. Proof of Theorem 1.2

We will prove Theorem 1.2 in a more general setting. Indeed, we will assume that the damping term γ in Equation (1.1) is a real function defined on $[t_0, +\infty[$ that belongs to the class $W_{loc}^{1,1}([t_0, +\infty[, \mathbb{R})])$ and satisfies:

There exists
$$K > 3$$
 such that $\gamma(t) \ge \frac{K}{t} \ \forall t \ge t_0,$ (2.1)

and

$$\int_{t_0}^{+\infty} \left[(t\gamma(t))' \right]_+ dt < +\infty, \tag{2.2}$$

where $[(t\gamma(t))']_{+} \equiv \max\{(t\gamma(t))', 0\}$ is the positive part of $(t\gamma(t))'$.

Typical examples of functions γ satisfying (2.1) and (2.2) are $\gamma(t) = \frac{K}{a+t}$ with $a \in \mathbb{R}$ and K > 3.

Proof [Proof of Theorem 1.2] We will use a modified version of a method introduced by Cabot and Frankel in [2] and recently developed in [3].

Let $x^* \in \arg \min \Phi$ and define the function $h: [t_0, +\infty[\to \mathbb{R}^+ \text{ by } h(t) = \frac{1}{2} \|x(t) - x^*\|^2$. By differentiating, we have

$$h'(t) = \langle x'(t), x(t) - x^* \rangle,$$

$$h''(t) = ||x'(t)||^2 + \langle x^{''}(t), x(t) - x^* \rangle.$$

Combining these last equalities and using Equation (1.1), we get

$$h''(t) + \gamma(t)h'(t) = \|x'(t)\|^2 + \langle \nabla \Phi(x(t)), x^* - x(t) \rangle.$$
 (2.3)

Using now the convexity inequality

$$\Phi(x^*) \ge \Phi(x) + \langle \nabla \Phi(x), x^* - x \rangle, \tag{2.4}$$

and the definition (1.2) of the energy function W, we obtain

$$W(t) \le \frac{3}{2} \|x'(t)\|^2 - h''(t) - \gamma(t)h'(t). \tag{2.5}$$

On the other hand, in view of (1.1),

$$W'(t) = \langle x'(t), x(t) \rangle + \langle \nabla \Phi(x(t)), x'(t) \rangle$$
$$= -\gamma(t) \|x'(t)\|^{2}.$$

Hence

$$(t^{2}W(t))' = 2tW(t) - t^{2}\gamma(t) \|x'(t)\|^{2}.$$
(2.6)

Using now assumption (2.1), we get

$$\frac{3}{2}t \|x'\|^{2} \leq \frac{3}{2K}t^{2}\gamma(t) \|x'(t)\|^{2}
= \frac{3}{K}tW(t) - \frac{3}{2K}(t^{2}W(t))'.$$
(2.7)

Multiplying (2.5) by t and using inequality (2.7), we obtain

$$(1 - \frac{3}{K})tW(t) + \frac{3}{2K} (t^2W(t))' \le -th''(t) - t\gamma(t)h'(t).$$

Integrating this last inequality on $[t_0, t]$, we get after simplification

$$(1 - \frac{3}{K}) \int_{t_0}^t sW(s)ds + \frac{3}{2K} (t^2W(t)) \le C_0 - th'(t) + (1 - t\gamma(t))h(t) + \int_{t_0}^t (s\gamma(s))'h(s)ds,$$
(2.8)

where $C_0 = \frac{3}{2K} (t_0^2 W(t_0)) + t_0 h'(t_0) - h(t_0)$.

Let $\varepsilon > 0$ such that $K > 3 + 3\varepsilon$. By using (2.1), we obtain from the inequality (2.8)

$$(1 - \frac{3}{K}) \int_{t_0}^{t} sW(s)ds + \frac{3}{2K} (t^2W(t)) + \varepsilon h(t) \le C_0 - th'(t) - (K - 1 - \varepsilon)h(t) + \int_{t_0}^{t} [(s\gamma(s))']_{+} h(s)ds.$$

Using now the fact that

$$t |h'(t)| \le t ||x'(t)|| ||x(t) - x^*||$$

 $\le 2\sqrt{t^2W(t)}\sqrt{h(t)},$

and applying the elementary inequality

$$\forall a > 0 \forall b, x \in \mathbb{R}, -ax^2 + bx \le \frac{b^2}{4a}$$

with $x = \sqrt{h(t)}$, we get

$$A \int_{t_0}^{t} sW(s)ds + Bt^2W(t) + \varepsilon h(t) \le C_0 + \int_{t_0}^{t} \left[(s\gamma(s))' \right]_{+} h(s)ds, \tag{2.9}$$

where $A = 1 - \frac{3}{K}$ and $B = \frac{3}{2K} - \frac{1}{K-1-\varepsilon}$.

Since $K > 3 + 3\varepsilon$, the constants A and B are positive; then

$$\varepsilon h(t) \le C_0 + \int_{t_0}^t \left[(s\gamma(s))' \right]_+ h(s) ds.$$

Hence, by using Gronwall's inequality and the assumption (2.2), we deduce that the function h is bounded, more precisely, we get

$$\sup_{t \geq t_0} h(t) \leq \frac{C_0}{\varepsilon} \exp(\frac{1}{\varepsilon} \int_{t_0}^{+\infty} \left[(s\gamma(s))' \right]_+ ds).$$

Therefore, we infer from (2.9) that

$$\sup_{t>t_0} t^2 W(t) < +\infty, \tag{2.10}$$

$$\int_{t_0}^{+\infty} sW(s)ds < +\infty. \tag{2.11}$$

Combining (2.6) and (2.11) yields that the positive part $\left[(t^2W(t))'\right]_+$ of $(t^2W(t)')$ belongs to $L^1([t_0, +\infty[, \mathbb{R}]; hence <math>m := \lim_{t \to +\infty} t^2W(t)$ exists. This limit m must be equal to 0, since otherwise $tW(t) \simeq \frac{m}{t}$ as $t \to +\infty$, which contradicts (2.11). It remains to prove the weak convergence of x(t) as $t \to +\infty$. Let us note that (2.10) implies that $\Phi(x(t)) \to \min \Phi$ as $t \to +\infty$. Hence by using the weak lower semicontinuity of the function Φ , we deduce that if $x(t_n) \to \bar{x}$ weakly in \mathcal{H} with $t_n \to +\infty$ then $\Phi(\bar{x}) \le \min \Phi$, which is equivalent to $\bar{x} \in \arg \min \Phi$. On the other hand, from the convex inequality (2.4) we deduce that $\langle \nabla \Phi(x), x^* - x \rangle \le 0$ for every $x \in \mathcal{H}$. Then Equation (2.3) implies

$$h''(t) + \gamma(t)h'(t) \le ||x'(t)||^2$$
.

Multiply this last equation by $e^{\Gamma(t,t_0)}$, where $\Gamma(t,s)=\int_s^t \gamma(\tau)dt$, and integrate between t_0 and t, and we obtain

$$h'(t) \le e^{-\Gamma(t,t_0)} h'(t_0) + \int_{t_0}^t e^{-\Gamma(t,\tau)} \|x'(\tau)\|^2 d\tau.$$
 (2.12)

In view of the assumption (2.1), a simple calculation gives

$$\forall s \ge t_0, \quad \int_{0}^{+\infty} e^{-\Gamma(t,s)} dt \le \frac{s}{K-1}.$$

Hence by using (2.12) and Fubini Theorem, we get

$$\int_{t_0}^{+\infty} [h'(t)]_+ dt \le \frac{t_0 |h'(t_0)|}{K - 1} + \frac{1}{K - 1} \int_{t_0}^{+\infty} \tau \|x'(\tau)\|^2 d\tau.$$

Thanks to (2.11), the right-hand side of the last inequality is finite; thus $\int_{t_0}^{+\infty} [h'(t)]_+ dt < +\infty$, which implies that $\lim_{t\to +\infty} h(t)$ exists. Hence, for every $x^* \in \arg\min \Phi$, the limit of $||x(t) - x^*||$ as $t\to +\infty$ exists. Therefore, Opial's lemma [4], which we recall below, guarantees the required weak convergence of x(t) in $\mathcal H$ to some element of $\arg\min \Phi$.

Lemma 2.1 (Opial's lemma) Let $x:[t_0,+\infty[\to \mathcal{H}. \ Assume \ that \ there \ exists \ a \ nonempty \ subset \ S \ of \ \mathcal{H}$ such that:

- i) If $t_n \to +\infty$ and $x(t_n) \rightharpoonup x$ weakly in \mathcal{H} , then $x \in S$.
- ii) For every $z \in S$, $\lim_{t \to +\infty} ||x(t) z||$ exists.

Then there exists $z_{\infty} \in S$ such that $x(t) \rightharpoonup z_{\infty}$ weakly in \mathcal{H} as $t \to +\infty$.

3. Conclusion

In this paper, we have proved that if the damping term $\gamma(t)$ behaves at infinity like $\frac{K}{t}$ with K > 3, then every solution x(t) of the equation (1.1) converges weakly as $t \to +\infty$ to a minimizer of Φ and the energy function W(t) is $\circ(t^{-2})$. However, two important questions remain open. The first one is on the behavior of the solution x(t) in the limit case K = 3 and the second one is about the effect of the constant K on the convergence rate of the associated energy function W(t).

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