



Sectional curvatures on Weyl manifolds with a special metric connection

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Received: 27.03.2018

Accepted/Published Online: 29.11.2018

Final Version: 18.01.2019

Abstract: In this paper, Weyl manifolds, denoted by $WS(g, w, \pi, \mu)$, having a special a semisymmetric recurrent-metric connection are introduced and the uniqueness of this connection is proved. We give an example of $WS(g, w, \pi, \mu)$ with a constant scalar curvature. Furthermore, we define sectional curvatures of $WS(g, w, \pi, \mu)$ and prove that any isotropic Weyl manifold $WS(g, w, \pi, \mu)$ is locally conformal to an Einstein manifold with a semisymmetric recurrent-metric connection, $EWS(g, w, \pi, \mu)$.

Key words: Weyl manifold, semisymmetric connection, recurrent-metric connection, generalized Bianchi identities, sectional curvature

1. Introduction

Linear connections are defined on manifolds to establish the parallel transport of vector fields along any curve in the manifolds, so that infinitesimally close tangent spaces are connected to each other. Riemannian manifolds are defined by a linear metric connection or by the Levi-Civita connection, and this connection is uniquely defined once the metric tensor is determined or given and the metric tensor keeps the geometrical information about the space.

The concept of the semisymmetric linear connection in a differentiable manifold without metricity condition was introduced by Friedmann and Schouten in 1924 (see [4, p. 214]). Later, Hayden introduced the idea of metric connection with torsion on a Riemannian manifold in 1932 [5]. Afterwards, Yano considered the semisymmetric metric connection on a Riemannian manifold in 1970 [18].

Spaces with metric, nonmetric, torsion-free, or torsionful connections have wide applications in theories of gravity as well as differential geometry [16].

This paper is devoted to the study of Weyl manifolds endowed with a semisymmetric recurrent-metric connection, which we denote by $WS(g, w, \pi, \mu)$. We derive some relations involving the curvature tensor of a semisymmetric recurrent-metric connection. Moreover, we give an example of $WS(g, w, \pi, \mu)$ spaces with constant scalar curvature.

We also define the sectional curvature of Weyl manifolds with a semisymmetric recurrent-metric connection $WS(g, w, \pi, \mu)$ and we prove that any isotropic Weyl manifold with a semisymmetric recurrent-metric connection can be locally conformal to an Einstein manifold with semisymmetric recurrent-metric connection $EWS(g, w, \pi, \mu)$.

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2010 *AMS Mathematics Subject Classification*: Primary 53A30, 53A40

2. Preliminaries

In this section, we first give some preliminary concepts related to Weyl spaces and semisymmetric spaces.

Two Riemannian metrics g and \bar{g} are conformal if they coincide up to a factor that is positive function, i.e. $\bar{g} = e^{2\lambda}g$. This is an equivalence relation, each class G being called a conformal structure. A Weyl structure is a map $w : G \rightarrow \Lambda^1(W)$ satisfying $w(e^{2\lambda}g) = w(g) + 2d\lambda$. A manifold with a Weyl structure is called a Weyl manifold, denoted by $W(g, w)$.

In [3], it was proved that for a Weyl manifold $W(g, w)$, there exists a unique torsion-free connection ∇ that preserves the conformal class G . Preserving the conformal class means that for any $g \in G$ there exists 1-form w such that

$$\nabla g = 2w \otimes g. \tag{2.1}$$

Equation (2.1) can be expressed in local coordinates as

$$\nabla_k g_{ij} = 2w_k g_{ij}. \tag{2.2}$$

Here, w is a 1-form called a complementary covector field.

Under the renormalization of the metric tensor g ,

$$\bar{g} = \Omega^2 g, \quad (\Omega > 0), \tag{2.3}$$

the 1-form w is transformed by the law

$$\bar{w} = w + d \ln \Omega, \tag{2.4}$$

so that

$$\nabla_k \bar{g}_{ij} = 2\bar{w}_k \bar{g}_{ij}. \tag{2.5}$$

Here, Ω is a positive scalar differentiable function defined on $W(g, w)$ (see [6] and [10, p. 152]).

The relation between the Weyl connection ∇ and the Riemannian connection ∇^g is

$$\nabla_X Y = \nabla_X^g Y - w(X)Y - w(Y)X + g(X, Y)\psi, \tag{2.6}$$

where X, Y are vector fields on $W(g, w)$ and ψ is the dual vector field to w such that $w(X) = g(X, \psi)$.

In local coordinates, (2.6) can be given by

$$\Gamma_{ji}^h = \left\{ \begin{matrix} h \\ ji \end{matrix} \right\} - (w_j \delta_i^h + w_i \delta_j^h - w^h g_{ji}), \tag{2.7}$$

where Γ_{ji}^h are the coefficients of the Weyl connection and

$$\left\{ \begin{matrix} h \\ ji \end{matrix} \right\} = \frac{1}{2} g^{hm} (\partial_j g_{mi} + \partial_i g_{mj} - \partial_m g_{ji}), \tag{2.8}$$

are the coefficients of the Levi-Civita connection ∇^g see [2, p. 81], and [10, p. 154].

The curvature tensor W of ∇ is given by

$$W(X, Y)Z = \nabla_X \nabla_Y Z - \nabla_Y \nabla_X Z - \nabla_{[X, Y]} Z. \tag{2.9}$$

Using (2.6) in (2.9), the curvature tensor of $W(g, w)$ is obtained:

$$W(X, Y)Z = R(X, Y)Z - s(X, Z)Y + s(Y, Z)X + s(Y, X)Z - s(Y, X)Z + g(Y, Z)\bar{S}X - g(X, Z)\bar{S}Y, \quad (2.10)$$

for any vector fields X, Y, Z where R denotes the curvature tensor of the Riemannian connection ∇^g and s is the tensor field of type $(0, 2)$ defined by

$$s(X, Y) = (\nabla_X w)(Y) + w(X)w(Y) - \frac{1}{2}w(\psi)g(X, Y), \quad (2.11)$$

and \bar{S} is the tensor field of type $(1, 1)$ defined by

$$g(\bar{S}X, Y) = s(X, Y). \quad (2.12)$$

In local coordinates, using the curvature tensor of $W(g, w)$,

$$W_{kji}^h = \partial_k \Gamma_{ji}^h - \partial_j \Gamma_{ki}^h - \Gamma_{kt}^t \Gamma_{jt}^h + \Gamma_{ji}^t \Gamma_{kt}^h, \quad (2.13)$$

we obtain the components of curvature tensor W_{kji}^h as

$$W_{kji}^h = R_{kji}^h - w_{ki}\delta_j^h + w_{ji}\delta_k^h + (w_{jk} - w_{kj})\delta_i^h + g^{hs}(w_{ks}g_{ji} - w_{js}g_{ki}), \quad (2.14)$$

with

$$w_{ji} = \nabla_j w_i + w_j w_i - \frac{1}{2}w_t w^t g_{ji}, \quad (2.15)$$

and R_{kji}^h represents the Riemannian curvature tensor with respect to the Levi-Civita connection.

The curvature tensor and covariant curvature tensor, the Ricci tensor, and the scalar curvature of Weyl space are defined through parallel transportation of vector fields v , respectively, by

$$W_{kji}^h = W_{kji}^m g_{ml}, \quad (2.16)$$

$$W_{ji} = g^{kl} W_{kji}^h = W_{lji}^l, \quad (2.17)$$

$$W = g^{ij} W_{ij}. \quad (2.18)$$

By transvecting (2.14) with g_{hm} , we get

$$W_{kjim} = R_{kjim} - w_{ki}g_{jm} + w_{ji}g_{km} + (w_{jk} - w_{kj})g_{im} + w_{km}g_{ji} - w_{jm}g_{ki}, \quad (2.19)$$

which is called the covariant curvature tensor of $W(g, w)$.

The curvature tensor, the covariant curvature tensor, and the Ricci tensor of $W(g, w)$ satisfy the following symmetry properties (see [10, p. 157]):

$$W_{kji}^h = -W_{jkil}, \quad (2.20)$$

$$W_{kji}^h + W_{kjli}^h = 2g_{il}(\nabla_j w_k - \nabla_k w_j) = 4g_{il}\nabla_{[j}w_{k]}. \quad (2.21)$$

By using symmetries of the curvature tensor of Weyl space, we obtain identities that are similar to identities held in Riemannian spaces. The following identities are known as the first and second Bianchi identities for Weyl spaces, respectively [6, 11]:

$$W_{kji}{}^l + W_{jik}{}^l + W_{ikj}{}^l = 0, \tag{2.22}$$

$$\nabla_m W_{kji}{}^l + \nabla_k W_{jmi}{}^l + \nabla_j W_{mki}{}^l = 0. \tag{2.23}$$

Furthermore, the Ricci tensor of the Weyl manifold is computed in terms of Ricci curvature R_{kj} of Riemannian space as

$$W_{kj} = R_{kj} + (n - 2)w_{kj} + (w_{kj} - w_{jk}) + w_{st} g^{st} g_{kj}, \tag{2.24}$$

where w_{kj} is defined in (2.15). It should also be noted that the Ricci tensor of a Weyl manifold is not symmetric; its symmetric and antisymmetric parts are given as follows (see [2], p.82):

$$W_{(kj)} = R_{kj} + \frac{1}{2}(n - 2)[\nabla_j w_k + \nabla_k w_j + 2w_k w_j - 2w_t w^t g_{jk}] + g_{jk} \nabla_t w^t, \tag{2.25}$$

$$W_{[kj]} = n \nabla_{[k} w_{j]}. \tag{2.26}$$

From (2.24), we obtain the scalar curvature of Weyl space:

$$W = R + 2(n - 1)\nabla_j w^j - (n - 1)(n - 2)w_j w^j, \tag{2.27}$$

where R is the Riemannian scalar curvature and w is the complementary covector defined in (2.2).

In the next section, we give some definitions and properties of manifolds with semisymmetric connection, and we also construct a new special connection on a Weyl manifold.

3. Semisymmetric recurrent-metric connection on Weyl manifolds

In the literature, the idea of semisymmetric connection was introduced by [1, 13, 15, 17, 18] and curvature-related properties were studied widely therein. Let M be an n -dimensional, ($n > 2$) differentiable manifold. A linear connection ∇^* on M , whose coefficients are $\Gamma^*_{jk}{}^i$, is said to be semisymmetric if the torsion tensor T of ∇^* satisfies the relation

$$T(X, Y) = \pi(Y)X - \pi(X)Y, \tag{3.1}$$

where π is a 1-form, and X, Y are smooth vector fields on M . In local coordinates, (3.1) can be written as

$$T_{jk}{}^i = \Gamma^*_{jk}{}^i - \Gamma^*_{kj}{}^i = \pi_k \delta_j^i - \pi_j \delta_k^i. \tag{3.2}$$

In addition, if a semisymmetric connection has the recurrency condition

$$\nabla_X^* g = 2 \mu(X)g \tag{3.3}$$

in local coordinates, (3.3) can be written as

$$\nabla^*_k g_{ij} = 2\mu_k g_{ij}, \tag{3.4}$$

and then the connection ∇^* is said to be a semisymmetric recurrent-metric connection and μ is called the recurrent covariant vector field [7, 8].

In this work, we use the notion of a semisymmetric recurrent metric connection for Weyl manifolds. Let $\bar{\nabla}$ be a linear connection with coefficients $\bar{\Gamma}^i_{jk}$ on a Weyl manifold $W(g, w)$ satisfying (3.2). If the following relation also holds on $W(g, w)$,

$$\bar{\nabla}_X g(Y, Z) = 2(w + \mu)(X)g(Y, Z), \tag{3.5}$$

in local coordinates, (3.5) is represented by

$$\bar{\nabla}_k g_{ij} = \nabla_k g_{ij} + 2\mu_k g_{ij} = 2(w_k + \mu_k)g_{ij}, \tag{3.6}$$

and then $W(g, w)$ is called a Weyl manifold with a semisymmetric recurrent-metric connection denoted by $WS(g, w, \pi, \mu)$.

From (2.2), we have

$$\begin{aligned} \nabla_k g_{ij} &= \partial_k g_{ij} - g_{hj} \Gamma^h_{ki} - g_{ih} \Gamma^h_{kj} \\ &= 2w_k g_{ij}, \end{aligned} \tag{3.7}$$

and from (3.6), more explicitly,

$$\begin{aligned} \bar{\nabla}_k g_{ij} &= \partial_k g_{ij} - g_{hj} \bar{\Gamma}^h_{ki} - g_{ih} \bar{\Gamma}^h_{kj} \\ &= 2(w_k + \mu_k)g_{ij}. \end{aligned} \tag{3.8}$$

By using (3.8), we have

$$\bar{\nabla}_k g^{ij} = -2(w_k + \mu_k)g^{ij}. \tag{3.9}$$

Here, we will examine the existence and uniqueness of the semisymmetric recurrent-metric connection $\bar{\nabla}$ on a Weyl manifold and will prove the following theorem.

Theorem 3.1 *Let $WS(g, w, \pi, \mu)$ be an n -dimensional Weyl manifold equipped with the semisymmetric recurrent-metric connection $\bar{\nabla}$ associated with 1-forms w, π , and μ satisfying (2.2), (3.2), and (3.4), respectively. Then there exists a unique connection $\bar{\nabla}$ on $WS(g, w, \pi, \mu)$ given by*

$$\bar{\nabla}_X Y = \nabla_X Y - \mu(X)Y - \mu(Y)X + g(X, Y)\xi + \pi(Y)X - g(X, Y)\eta, \tag{3.10}$$

where ξ and η are dual vector fields such that

$$\mu(X) = g(X, \xi), \quad \pi(X) = g(X, \eta). \tag{3.11}$$

Proof Let $\bar{\nabla}$ be a semisymmetric recurrent metric connection and ∇ be a Weyl connection. We have

$$(\bar{\nabla}_X g)(Y, Z) = \bar{\nabla}_X g(Y, Z) - g(\bar{\nabla}_X Y, Z) - g(Y, \bar{\nabla}_X Z) \tag{3.12}$$

and

$$(\nabla_X g)(Y, Z) = \nabla_X g(Y, Z) - g(\nabla_X Y, Z) - g(Y, \nabla_X Z) \tag{3.13}$$

for any vector fields X, Y , and Z .

We put

$$\bar{\nabla}_X Y = \nabla_X Y + U(X, Y), \tag{3.14}$$

where U is a tensor field of type $(1, 2)$ defined as the difference of the connections.

Using (3.1) and (3.14) it is obtained that

$$\begin{aligned} T(X, Y) &= \bar{\nabla}_X Y - \bar{\nabla}_Y X - [X, Y], \\ &= U(X, Y) - U(Y, X). \end{aligned} \tag{3.15}$$

From (3.12), (3.13), (3.14), and (3.15) we get

$$g(U(X, Y), Z) + g(U(X, Z), Y) = -2\mu(X)g(Y, Z). \tag{3.16}$$

By using (3.15), and permuting vector fields X, Y , and Z for T , we get

$$\begin{aligned} g(T(X, Y), Z) &= g(U(X, Y), Z) - g(U(Y, X), Z), \\ g(T(Z, X), Y) &= g(U(Z, X), Y) - g(U(X, Z), Y), \\ g(T(Z, Y), X) &= g(U(Z, Y), X) - g(U(Y, Z), X). \end{aligned} \tag{3.17}$$

From (3.16) and (3.17) we obtain

$$\begin{aligned} g(T(X, Y), Z) + g(T(Z, X), Y) + g(T(Z, Y), X) &= 2g(U(X, Y), Z) + 2\mu(X)g(Y, Z) \\ &\quad + 2\mu(Y)g(Z, X) - 2\mu(Z)g(X, Y). \end{aligned} \tag{3.18}$$

Defining the tensor \acute{T} of type $(1, 2)$ as

$$g(T(Z, X), Y) = g(\acute{T}(X, Y), Z), \tag{3.19}$$

equation (3.18) can be written as

$$\begin{aligned} g(U(X, Y), Z) &= \frac{1}{2}[g(T(X, Y), Z) + g(\acute{T}(X, Y), Z) + g(\acute{T}(Y, X), Z)] \\ &\quad - \mu(X)g(Y, Z) - \mu(Y)g(Z, X) + \mu(Z)g(X, Y). \end{aligned} \tag{3.20}$$

Thus, we find

$$U(X, Y) = \frac{1}{2}[T(X, Y) + \acute{T}(X, Y) + \acute{T}(Y, X)] - \mu(X)Y - \mu(Y)X + g(X, Y)\xi, \tag{3.21}$$

where $\mu(X) = g(X, \xi)$.

From (3.1) and (3.19) we have

$$\begin{aligned} g(T(Z, X), Y) &= g(\pi(X)Z, Y) - g(\pi(Z)X, Y), \\ &= g(\acute{T}(X, Y), Z). \end{aligned} \tag{3.22}$$

From (3.19), and (3.22), we reach

$$g(\acute{T}(X, Y), Z) = \pi(X)g(Z, Y) - g(Z, \eta)g(X, Y), \tag{3.23}$$

which implies

$$\acute{T}(X, Y) = \pi(X)Y - g(X, Y)\eta, \tag{3.24}$$

where $\pi(X) = g(X, \eta)$.

Hence, (3.21) turns into

$$U(X, Y) = \pi(Y)X - \mu(X)Y - \mu(Y)X + g(X, Y)\xi - g(X, Y)\eta. \tag{3.25}$$

Then (3.14) becomes

$$\bar{\nabla}_X Y = \nabla_X Y - \mu(X)Y - \mu(Y)X + g(X, Y)\xi + \pi(Y)X - g(X, Y)\eta,$$

which completes the proof. □

Also, equation (3.10) is obtained in local coordinates as

$$\bar{\Gamma}_{ik}^l = \left\{ \begin{matrix} l \\ ik \end{matrix} \right\} - (w_i \delta_k^l + w_k \delta_i^l - w^l g_{ik}) + (\lambda_k \delta_i^l - \mu_i \delta_k^l - \lambda^l g_{ik}), \tag{3.26}$$

where $\lambda_k = \pi_k - \mu_k$.

The following subsection is devoted to presentation of curvature tensors of Weyl manifolds with the semisymmetric recurrent-metric connection, $WS(g, w, \pi, \mu)$, in local coordinates in detail. The covariant curvature tensor, the Ricci tensor, and the scalar curvature of $WS(g, w, \pi, \mu)$ will be denoted by \bar{R}_{kjim} , \bar{R}_{ji} , and \bar{R} , respectively.

3.1. The curvature tensor of $WS(g, w, \pi, \mu)$

Theorem 3.2 *The curvature tensor of a Weyl manifold with the semisymmetric recurrent-metric connection $WS(g, w, \pi, \mu)$ has the following properties:*

$$(i) \quad \bar{R}_{kjim} = W_{kjim} + Q_{kjim} - \alpha_{ij}g_{mk} + \alpha_{ik}g_{mj} - \alpha_{mk}g_{ij} + \alpha_{mj}g_{ik}, \tag{3.27}$$

where W_{kjim} is the covariant curvature tensor of $W(g, w)$, and

$$Q_{kji}^h = \lambda_{ki}\delta_j^h - \lambda_{ji}\delta_k^h + \lambda_{jl}g^{lh}g_{ki} - \lambda_{kl}g^{lh}g_{ji} + \delta_i^h(\nabla_j\mu_k - \nabla_k\mu_j), \tag{3.28}$$

or by transvecting with metric tensor g_{hm}

$$Q_{kjim} = g_{mj}\lambda_{ki} - g_{mk}\lambda_{ji} + \lambda_{jm}g_{ki} - \lambda_{km}g_{ji} + 2g_{mi}\nabla_{[j}\mu_{k]}, \tag{3.29}$$

and

$$\alpha_{ij} = \lambda_i w_j + \lambda_j w_i - w_l \lambda^l g_{ij}, \tag{3.30}$$

$$\lambda_{ki} = \nabla_k \lambda_i - \lambda_k \lambda_i + \frac{1}{2} g_{ki} \lambda_t \lambda^t, \tag{3.31}$$

and also, λ , w , and μ are 1-forms as given in (3.2), (2.2), and (3.4), respectively.

$$(ii) \quad \bar{R}_{ji} = W_{ji} + Q_{ji} - (n - 2)(\lambda_j w_i + \lambda_i w_j - 2g_{ji} w_t \lambda^t), \tag{3.32}$$

where W_{ji} is given by (2.24), and

$$Q_{ji} = Q_{kjim} g^{km} = (n - 2)[- \nabla_j \lambda_i + \lambda_i \lambda_j - g_{ji} \lambda_t \lambda^t] - g_{ji} \nabla_t \lambda^t + 2 \nabla_{[j} \mu_{i]}. \tag{3.33}$$

$$(iii) \quad \bar{R} = R + 2(n - 1)(\nabla_t w^t - \nabla_t \lambda^t) - (n - 1)(n - 2)(w_t - \lambda_t)(w^t - \lambda^t), \tag{3.34}$$

where R is the Riemannian scalar curvature and holds the relation $R = R_{ji} g^{ji}$.

Proof

- (i) In spaces with torsion, parallel transport of vector fields is defined by [18]. The curvature tensor of $WS(g, w, \pi, \mu)$ can be computed by using the Ricci identity for a covariant vector field v_i :

$$(\bar{\nabla}_k \bar{\nabla}_j - \bar{\nabla}_j \bar{\nabla}_k) v_i = -\bar{R}_{kji}{}^t v_t - T_{kj}{}^t \bar{\nabla}_t v_i, \tag{3.35}$$

where $\bar{R}_{kji}{}^h$ is the curvature tensor of $WS(g, w, \pi, \mu)$,

$$\bar{R}_{kji}{}^h = \partial_k \bar{\Gamma}_{ji}{}^h - \partial_j \bar{\Gamma}_{ki}{}^h + \bar{\Gamma}_{ji}{}^t \bar{\Gamma}_{kt}{}^h - \bar{\Gamma}_{ki}{}^t \bar{\Gamma}_{jt}{}^h, \tag{3.36}$$

and $T_{ij}{}^h$ is the torsion tensor of $WS(g, w, \pi, \mu)$,

$$T_{ij}{}^h = \bar{\Gamma}_{ij}{}^h - \bar{\Gamma}_{ji}{}^h = U_{ij}{}^h - U_{ji}{}^h. \tag{3.37}$$

Substituting coefficients of connections (3.26) in (3.36), and after some calculations, we obtain the curvature tensor of $WS(g, w, \pi, \mu)$ as

$$\begin{aligned} \bar{R}_{kji}{}^h &= W_{kji}{}^h + Q_{kji}{}^h - \delta_k^h (\lambda_j w_i + \lambda_i w_j - w_l \lambda^l g_{ij}) + \delta_j^h (\lambda_k w_i + \lambda_i w_k - w_l \lambda^l g_{ik}) \\ &\quad - g_{ij} (\lambda_k w^h + \lambda^h w_k - w_l \lambda^l \delta_k^h) + g_{ik} (\lambda_j w^h + \lambda^h w_j - w_l \lambda^l \delta_j^h), \end{aligned} \tag{3.38}$$

where $W_{kji}{}^h$ represents the curvature tensor of Weyl space defined in (2.14).

If we simplify our calculations we define the tensor $Q_{kji}{}^h$ as in (3.28) and λ_{ki} as in (3.31), respectively:

$$Q_{kji}{}^h = \delta_j^h \lambda_{ki} - \delta_k^h \lambda_{ji} + \lambda_{jl} g^{lh} g_{ki} - \lambda_{kl} g^{lh} g_{ji} + \delta_i^h (\nabla_j \mu_k - \nabla_k \mu_j),$$

$$\lambda_{ki} = \nabla_k \lambda_i - \lambda_k \lambda_i + \frac{1}{2} g_{ki} \lambda_t \lambda^t,$$

and multiplying (3.28) by the metric tensor g_{hm} , we obtain

$$Q_{kjim} = g_{mj} \lambda_{ki} - g_{mk} \lambda_{ji} + \lambda_{jm} g_{ki} - \lambda_{km} g_{ji} + 2g_{mi} \nabla_{[j} \mu_{k]},$$

and similarly, we get

$$Q_{ji} = Q_{kjim} g^{km} = (n-2)[- \nabla_j \lambda_i + \lambda_i \lambda_j - g_{ji} \lambda_t \lambda^t] - g_{ji} \nabla_t \lambda^t + 2 \nabla_{[j} \mu_{i]}.$$

From (3.29), we see that the following antisymmetry property holds for Q_{kjil} :

$$Q_{kjil} = -Q_{jkil}. \tag{3.39}$$

Multiplying (3.38) by metric tensor g_{hm} and using (3.29) and (2.19), we reach (3.27):

$$\bar{R}_{kjim} = W_{kjim} + Q_{kjim} - g_{mk} \alpha_{ij} + g_{mj} \alpha_{ik} - g_{ij} \alpha_{mk} + g_{ik} \alpha_{mj},$$

where

$$\alpha_{ij} = \lambda_i w_j + \lambda_j w_i - g_{ij} w_l \lambda^l.$$

(ii) Now let us examine the Ricci curvature and its symmetric properties for $WS(g, w, \pi, \mu)$. Multiplying (3.27) by g^{mk} , we get the Ricci tensor of $WS(g, w, \pi, \mu)$ as

$$\bar{R}_{ji} = W_{ji} + Q_{ji} - (n-2)(\lambda_j w_i + \lambda_i w_j) + 2(n-2)g_{ji} w_l \lambda^l, \tag{3.40}$$

where W_{ji} and Q_{ji} are given by (2.24) and (3.33), respectively.

It is seen that the Ricci tensor \bar{R}_{ji} of $WS(g, w, \pi, \mu)$ is not symmetric. The symmetric and antisymmetric parts of \bar{R}_{ji} can be calculated as

$$\begin{aligned} \bar{R}_{(ji)} &= W_{(ji)} - \frac{(n-2)}{2} [(\nabla_j \lambda_i + \nabla_i \lambda_j) - 2\lambda_i \lambda_j + 2g_{ji} \lambda_t \lambda^t + 2(\lambda_j w_i + \lambda_i w_j)] \\ &\quad + g_{ji} [2(n-2)w_t \lambda^t - \nabla_t \lambda^t] \end{aligned} \tag{3.41}$$

and

$$\bar{R}_{[ji]} = n \nabla_{[j} w_{i]} - (n-2) \nabla_{[j} \lambda_{i]} + 2 \nabla_{[j} \mu_{i]}. \tag{3.42}$$

(iii) Transvecting the Ricci curvature tensor \bar{R}_{ji} in (3.40) by the metric tensor g^{ji} , we obtain

$$\bar{R} = \bar{R}_{ji} g^{ji}, \tag{3.43}$$

and then the scalar curvature of $WS(g, w, \pi, \mu)$ is

$$\bar{R} = R + 2(n-1)(\nabla_t w^t - \nabla_t \lambda^t) - (n-1)(n-2)(w_t - \lambda_t)(w^t - \lambda^t). \tag{3.44}$$

□

Now we examine the properties of the covariant curvature tensor of Weyl space with the semisymmetric recurrent-metric connection $WS(g, w, \pi, \mu)$. Using the properties of R_{kjil} and W_{kjil} , it can be seen that the curvature tensor of $WS(g, w, \pi, \mu)$ satisfies the following symmetry relations in the following propositions:

Proposition 3.3 *The curvature tensor of $WS(g, w, \pi, \mu)$ satisfies the following symmetry relations:*

$$(i). \quad \bar{R}_{kjim} = -\bar{R}_{jkim}, \tag{3.45}$$

$$(ii). \quad \bar{R}_{kjim} + \bar{R}_{kjmi} = 4g_{im}(\nabla_{[j}\mu_{k]} + \nabla_{[j}w_k]). \tag{3.46}$$

Proof

(i) Interchanging the indices k and j in equation (3.27), we have

$$\bar{R}_{kjim} + \bar{R}_{jkim} = W_{kjim} + W_{jkim} + Q_{kjim} + Q_{jkim},$$

and using (2.20) and (3.39) in the above equation, we get (3.45).

(ii) Using (2.20), (2.21), (3.29), and (3.30) in the equation of (3.27), we obtain (3.46).

$$\begin{aligned} \bar{R}_{kjim} + \bar{R}_{kjmi} &= W_{kjim} + W_{kjmi} + Q_{kjim} + Q_{kjmi} \\ &= W_{kjim} + W_{kjmi} + 2g_{im}(\mu_{jk} - \mu_{kj}) \\ &= 2g_{im}(\nabla_{[j}w_k] - \nabla_{[k}w_j]) + 2g_{im}(\mu_{jk} - \mu_{kj}) \\ &= 4g_{im}(\nabla_{[j}\mu_{k]} + \nabla_{[j}w_k]). \end{aligned}$$

Note that if w_k and μ_k are gradients or if w_k and μ_k have opposite signs, then

$$\bar{R}_{kjim} = -\bar{R}_{kjmi}.$$

In the following proposition, we introduce extended (generalized) first and second Bianchi identities for $WS(g, w, \pi, \mu)$.

□

Proposition 3.4 *The curvature tensor of $WS(g, w, \pi, \mu)$ satisfies the following first and second Bianchi identities for $WS(g, w, \pi, \mu)$, respectively:*

$$(i). \quad \bar{R}_{kji}{}^l + \bar{R}_{jik}{}^l + \bar{R}_{ikj}{}^l = 2(\delta_j^l \nabla_{[k}\pi_{i]} + \delta_i^l \nabla_{[j}\pi_{k]} + \delta_k^l \nabla_{[i}\pi_{j]}), \tag{3.47}$$

$$(ii). \quad (\bar{\nabla}_l \bar{R}_{kji}{}^t + \bar{\nabla}_k \bar{R}_{jli}{}^t + \bar{\nabla}_j \bar{R}_{lki}{}^t) = 2(\pi_l \bar{R}_{kji}{}^t + \pi_k \bar{R}_{jli}{}^t + \pi_j \bar{R}_{lki}{}^t). \tag{3.48}$$

Proof

(i) Using (3.38) and (3.31), and by changing indices k, j, i cyclically, we get

$$\bar{R}_{kji}{}^l + \bar{R}_{jik}{}^l + \bar{R}_{ikj}{}^l = Q_{kji}{}^l + Q_{jik}{}^l + Q_{ikj}{}^l. \tag{3.49}$$

On the other hand, using (3.31), we calculate $Q_{kji}{}^l + Q_{jik}{}^l + Q_{ikj}{}^l$ as

$$\begin{aligned} Q_{kji}{}^l + Q_{jik}{}^l + Q_{ikj}{}^l &= \delta_j^l(\lambda_{ki} - \lambda_{ik}) + \delta_k^l(\lambda_{ij} - \lambda_{ji}) + \delta_i^l(\lambda_{jk} - \lambda_{kj}) \\ &\quad + \delta_i^l(\mu_{jk} - \mu_{kj}) + \delta_k^l(\mu_{ij} - \mu_{ji}) + \delta_j^l(\mu_{ki} - \mu_{ik}), \end{aligned} \quad (3.50)$$

where

$$\mu_{kj} = \nabla_k \mu_j - \mu_j \mu_k + \frac{1}{2} g_{jk} \mu_t \mu^t. \quad (3.51)$$

Using (3.31), and arranging (3.50), we find

$$\begin{aligned} Q_{kji}{}^l + Q_{jik}{}^l + Q_{ikj}{}^l &= \delta_j^l[\nabla_k(\pi_i - \nabla_i \pi_k)] \\ &\quad + \delta_k^l[\nabla_i(\pi_j - \nabla_j \pi_i)] + \delta_i^l[\nabla_j(\pi_k - \nabla_k \pi_j)]. \end{aligned} \quad (3.52)$$

Using (3.52) in (3.49), we get

$$\bar{R}_{kji}{}^l + \bar{R}_{jik}{}^l + \bar{R}_{ikj}{}^l = 2(\delta_j^l \nabla_{[k} \pi_{i]} + \delta_i^l \nabla_{[j} \pi_{k]} + \delta_k^l \nabla_{[i} \pi_{j]}).$$

which is called the generalized first Bianchi identity.

Also, the covariant form is obtained:

$$\bar{R}_{kjim} + \bar{R}_{jikm} + \bar{R}_{ikjm} = 2(g_{jm} \nabla_{[k} \pi_{i]} + g_{im} \nabla_{[j} \pi_{k]} + g_{km} \nabla_{[i} \pi_{j]}).$$

(ii) Using the Ricci identity (3.35) and differentiating covariantly both sides of (3.35), we get

$$\begin{aligned} -\bar{\nabla}_l \bar{\nabla}_k \bar{\nabla}_j v_i + \bar{\nabla}_l \bar{\nabla}_j \bar{\nabla}_k v_i &= \bar{\nabla}_l (\bar{R}_{kji}{}^t) v_t + \bar{R}_{kji}{}^t \bar{\nabla}_l (v_t) \\ &\quad + \bar{\nabla}_l (T_{kj}{}^t) \bar{\nabla}_t v_i + T_{kj}{}^t (\bar{\nabla}_l \bar{\nabla}_t v_i), \end{aligned} \quad (3.53)$$

which is written in terms of the covariant derivative of the curvature tensor and torsion tensor.

Now, interchanging the indices l , k , and j in (3.53), using the components of the torsion tensor of (3.2), and by some tensor calculations, the Ricci identity (3.35) reduces to

$$\begin{aligned} T_{lk}{}^t (\bar{\nabla}_t \bar{\nabla}_j - \bar{\nabla}_j \bar{\nabla}_t) v_i + T_{jl}{}^t (\bar{\nabla}_t \bar{\nabla}_k - \bar{\nabla}_k \bar{\nabla}_t) v_i + T_{kj}{}^t (\bar{\nabla}_t \bar{\nabla}_l - \bar{\nabla}_l \bar{\nabla}_t) v_i \\ = -2\pi_k (\bar{R}_{lji}{}^t v_t + T_{lj}{}^t \bar{\nabla}_t v_i) - 2\pi_l (\bar{R}_{jki}{}^t v_t + T_{jk}{}^t \bar{\nabla}_t v_i) - 2\pi_j (\bar{R}_{kli}{}^t v_t + T_{kl}{}^t \bar{\nabla}_t v_i). \end{aligned} \quad (3.54)$$

Using (3.53), (3.54), and (3.2), we find

$$(\bar{\nabla}_l \bar{R}_{kji}{}^t + \bar{\nabla}_k \bar{R}_{jli}{}^t + \bar{\nabla}_j \bar{R}_{lki}{}^t) = 2(\pi_l \bar{R}_{kji}{}^t + \pi_k \bar{R}_{jli}{}^t + \pi_j \bar{R}_{lki}{}^t),$$

which is called the generalized second Bianchi identity for $WS(g, w, \pi, \mu)$.

□

Theorem 3.5 $WS(g, w, \pi, \mu)$ and $W(g, w)$ have the same curvature tensors if and only if the recurrent covariant vector field μ_k of $\bar{\nabla}$ defined by (3.4) is a gradient vector and the following equation holds:

$$\lambda_{ij} + \alpha_{ji} = 0, \tag{3.55}$$

where α_{ij} and λ_{ij} are as in (3.30) and (3.31), respectively.

Proof Let $WS(g, w, \pi, \mu)$ and $W(g, w)$ have the same curvature tensors:

$$\bar{R}_{kjim} = W_{kjim}.$$

Using (3.27), we have

$$Q_{kjim} = g_{mk}\alpha_{ij} - g_{mj}\alpha_{ik} + g_{ij}\alpha_{mk} - g_{ik}\alpha_{mj}. \tag{3.56}$$

Also, from (3.30), (3.31), and (3.29), we obtain the relation

$$g_{mj}(\alpha_{ik} + \lambda_{ki}) - g_{mk}(\alpha_{ij} + \lambda_{ji}) + g_{ik}(\alpha_{mj} + \lambda_{jm}) - g_{ij}(\alpha_{mk} + \lambda_{km}) + 2g_{mi}\nabla_{[j}\mu_{k]} = 0. \tag{3.57}$$

From (3.58), it follows that $\alpha_{ik} + \lambda_{ki} = 0$ and $\nabla_{[j}\mu_{k]} = 0$ simultaneously.

Conversely, using (3.27),

$$\bar{R}_{kjim} = W_{kjim} + Q_{kjim} - \alpha_{ij}g_{mk} + \alpha_{ik}g_{mj} - \alpha_{mk}g_{ij} + \alpha_{mj}g_{ik},$$

and substituting (3.29) in (3.27), we get

$$\begin{aligned} \bar{R}_{kjim} = & W_{kjim} + g_{mj}(\alpha_{ik} + \lambda_{ki}) - g_{mk}(\alpha_{ij} + \lambda_{ji}) + g_{ik}(\alpha_{mj} + \lambda_{jm}) \\ & - g_{ij}(\alpha_{mk} + \lambda_{km}) + 2g_{mi}\nabla_{[j}\mu_{k]}. \end{aligned} \tag{3.58}$$

By using the given assumptions of $\alpha_{ik} + \lambda_{ki} = 0$ and $\nabla_{[j}\mu_{k]} = 0$, we conclude that

$$\bar{R}_{kjim} = W_{kjim}. \tag{3.59}$$

□

Next, we give an example of 3-dimensional $WS(g, w, \pi, \mu)$ with a constant curvature in which components of the torsion tensor, complementary, and recurrency covector fields are chosen specially.

Example 3.6 Let us consider the three dimensional metric given as

$$ds^2 = \frac{dr^2}{1 - \kappa r^2} + r^2(d\theta^2 + \sin^2 \theta d\phi^2), \quad (r > 0, 0 \leq \theta < \pi, 0 \leq \phi < 2\pi), \tag{3.60}$$

where $1 - \kappa r^2 > 0$ and κ is an arbitrary constant.

The scalar curvature of (3.60) is obtained as $R = 6\kappa$. For $\kappa = 1, 0, -1$, space is called spherical, planar, and hyperbolic, respectively.

Here, all 1-forms w, π , and μ are represented with three components in spherical directions r, θ, ϕ , i.e. $w = (w_r, w_\theta, w_\phi)$, $\pi = (\pi_r, \pi_\theta, \pi_\phi)$. For this example, we choose the complementary vector w as $w = (0, w_\theta, 0)$,

covector fields π as $\pi = (0, \pi_\theta, 0)$, and recurrency form μ as $\mu = (0, \mu_\theta, 0)$, which are defined in equations (2.2), (3.2), and (3.4), respectively. Thus, for the metric (3.60), we find the connection coefficients, the Ricci curvature, and the scalar curvature of $WS(g, w, \pi, \mu)$ as follows:

$$\begin{aligned} \bar{\Gamma}_{rr}^r &= \frac{\kappa r}{1 - \kappa r^2}, & \bar{\Gamma}_{r\theta}^r &= -w_\theta + \mu_\theta, \\ \bar{\Gamma}_{\theta r}^r &= -(w_\theta + \mu_\theta), & \bar{\Gamma}_{\theta\theta}^r &= -r(1 - \kappa r^2), \\ \bar{\Gamma}_{\phi\phi}^r &= -r(1 - \kappa r^2) \sin^2 \theta, & \bar{\Gamma}_{rr}^\theta &= \frac{w_\theta - \mu_\theta}{r^2(1 - \kappa r^2)}, \\ \bar{\Gamma}_{\phi r}^\phi &= \bar{\Gamma}_{\theta r}^\theta = \bar{\Gamma}_{r\phi}^\phi = \bar{\Gamma}_{r\theta}^\theta = \frac{1}{r}, & \bar{\Gamma}_{\phi\theta}^\phi &= \cot \theta - w_\theta + \mu_\theta, \\ \bar{\Gamma}_{\theta\theta}^\theta &= -(w_\theta + \mu_\theta) = -\pi_\theta, & \bar{\Gamma}_{\phi\phi}^\theta &= -\sin \theta (\cos \theta + (\mu_\theta - w_\theta) \sin \theta), \\ \bar{\Gamma}_{\theta\phi}^\phi &= \cot \theta - \pi_\theta. \end{aligned} \tag{3.61}$$

Here, if we choose w_θ and μ_θ to be functions of radial coordinate r , and also π_θ to satisfy the relation $\pi_\theta = w_\theta(r) + \mu_\theta(r)$, then from (3.37), we can compute components of the torsion tensor for (3.62).

$$T_{r\theta}^r = T_{\phi\theta}^\phi = 2w_\theta, \quad T_{\theta r}^r = T_{\theta\phi}^\phi = -2w_\theta, \tag{3.62}$$

and the components of the Ricci tensor of $WS(g, w, \pi, \mu)$ are

$$\begin{aligned} \bar{R}_{rr} &= \frac{2\kappa r^2 + (w_\theta - \mu_\theta)(\cot \theta - w_\theta + \mu_\theta)}{r^2(1 - \kappa r^2)}, \\ \bar{R}_{r\theta} &= \frac{2rw'_\theta - w_\theta + \mu_\theta}{r}, \\ \bar{R}_{\theta r} &= -\frac{r(w'_\theta + \mu'_\theta) + w_\theta - \mu_\theta}{r}, \\ \bar{R}_{\theta\theta} &= 2\kappa r^2 + \cot \theta (w_\theta - \mu_\theta), \\ \bar{R}_{\phi\phi} &= \sin \theta (2\kappa r^2 \sin \theta + (w_\theta - \mu_\theta)[2 \cos \theta + (\mu_\theta - w_\theta) \sin \theta]), \end{aligned} \tag{3.63}$$

where prime ($'$) denotes the derivative with respect to r , and the scalar curvature of $WS(g, w, \pi, \mu)$ is obtained as

$$\bar{R} = \frac{1}{r^2} (6\kappa r^2 + 2(w_\theta - \mu_\theta)(2 \cot \theta - w_\theta + \mu_\theta)). \tag{3.64}$$

Particularly, in (3.64), by taking

$$w_\theta = \lambda_\theta = \mu_\theta = c_1 \sqrt{1 - \kappa r^2} + c_2 \left[-1 + \operatorname{artanh} \left(\frac{1}{\sqrt{1 - \kappa r^2}} \right) \right], \quad 1 - \kappa r^2 > 0, \tag{3.65}$$

where c_1, c_2 are any real constant, we obtain that the scalar curvature of $WS(g, w, \pi, \mu)$ is $\bar{R} = 6\kappa$. Thus, the scalar curvature of $WS(g, w, \pi, \mu)$ becomes the same as the scalar curvature of Riemannian space.

In the following section, sectional curvature is examined for $WS(g, w, \pi, \mu)$ in the sense of previous studies (see [9], p. 265, and [12]).

3.2. Sectional curvatures on Weyl manifolds with semisymmetric recurrent-metric connection

$WS(g, w, \pi, \mu)$

Let $X, Y \in T_p(WS)$ at a point $P \in WS(g, w, \pi, \mu)$. Let Π be the 2-plane spanned by X, Y . Then the sectional curvature K of $WS(g, w, \pi, \mu)$ at P with respect to plane Π is defined by [12]:

$$\begin{aligned} K(\Pi) &= K(X, Y) \\ &= \frac{\bar{R}(X, Y, X, Y)}{g(X, X)g(Y, Y) - g(X, Y)^2}. \end{aligned} \tag{3.66}$$

In local coordinates the equation (3.66) is

$$K(\Pi) = \frac{\bar{R}_{ijkl} X^i Y^j X^k Y^l}{(g_{ik}g_{jl} - g_{il}g_{jk}) X^i Y^j X^k Y^l}. \tag{3.67}$$

If the sectional curvature K of $WS(g, w, \pi, \mu)$ at a point P is the same for all 2-planes in $T_p(WS)$, then we say that $WS(g, w, \pi, \mu)$ is isotropic at P . If $WS(g, w, \pi, \mu)$ is isotropic at every point P on $WS(g, w, \pi, \mu)$, then $WS(g, w, \pi, \mu)$ is called isotropic (see [9], p. 265, and [12]).

On the other hand, if the sectional curvature K of $WS(g, w, \pi, \mu)$ has the same value at every point and for every section at that point, then we say that $WS(g, w, \pi, \mu)$ has constant curvature.

It is obvious that if $WS(g, w, \pi, \mu)$ has constant curvature, then $WS(g, w, \pi, \mu)$ is isotropic.

We recall that a Weyl manifold $WS(g, w, \pi, \mu)$ is said to be an Einstein manifold with respect to the semisymmetric recurrent-metric connection $EWS(g, w, \pi, \mu)$ if the symmetric part of the Ricci tensor is proportional to the metric; that is,

$$\bar{R}_{(ji)} = \theta g_{ij}, \tag{3.68}$$

for a scalar function θ defined on $WS(g, w, \pi, \mu)$ (see [14], Eq. 41).

It is shown that every 2-dimensional Weyl manifold is an Einstein manifold [12].

We now quote the following lemma (see [9], p. 265), which will be needed in the proof of the following theorem.

Lemma 3.7 *Suppose that S is any 4-covariant tensor, and that X and Y are two arbitrary linearly independent vectors. If for all X and Y*

$$S_{ijkl} X^i Y^j X^k Y^l = 0, \tag{3.69}$$

then we have

$$S_{ijkl} + S_{kl ij} + S_{il kj} + S_{kj il} = 0. \tag{3.70}$$

A sufficient condition for a Weyl manifold to be locally conformal to an Einstein manifold by using sectional curvature was given in [12]. By means of the notation used in [12], we state and prove the following theorem for $WS(g, w, \pi, \mu)$.

Theorem 3.8 *Any isotropic Weyl manifold with the semisymmetric recurrent-metric connection can be locally conformal to an Einstein manifold with the semisymmetric recurrent-metric connection, $EWS(g, w, \pi, \mu)$.*

Proof Assume that $WS(g, w, \pi, \mu)$ is an isotropic manifold. In Lemma 3.7, by taking

$$S_{ijkl} = \bar{R}_{ijkl} - K(x)(g_{ik}g_{jl} - g_{il}g_{jk}), \tag{3.71}$$

and using (3.70), we get

$$\bar{R}_{ijkl} + \bar{R}_{klij} + \bar{R}_{iljk} + \bar{R}_{kjil} = 4K g_{ik}g_{jl} - 2K(g_{lk}g_{ij} + g_{li}g_{kj}). \tag{3.72}$$

Transvecting (3.72) by g^{lh} ,

$$\bar{R}_{ijk}{}^h + \bar{R}_{kji}{}^h + (\bar{R}_{kji}l + \bar{R}_{kjil})g^{lh} = [4K g_{ik}g_{jl} - 2K(g_{lk}g_{ij} + g_{li}g_{kj})]g^{lh}, \tag{3.73}$$

and using symmetry properties (3.45) and (3.46) and the first Bianchi identity for $EWS(g, w, \pi, \mu)$,

$$\bar{R}_{kji}{}^l + \bar{R}_{jik}{}^l + \bar{R}_{ikj}{}^l = 2(\delta_j^l \nabla_{[k} \pi_{i]} + \delta_i^l \nabla_{[j} \pi_{k]} + \delta_k^l \nabla_{[i} \pi_{j]}), \tag{3.74}$$

we find that

$$\begin{aligned} \bar{R}_{ijk}{}^h + \bar{R}_{kji}{}^h + \bar{R}_{ikj}{}^h + 2\bar{R}_{ilkj}g^{lh} &= 2K(2g_{ik}\delta_j^h - g_{kj}\delta_i^h - g_{ij}\delta_k^h) \\ &\quad - g^{lh}A_{lkji} + 4\delta_j^h(\nabla_{[k} w_{i]} + \nabla_{[k} \mu_{i]}), \end{aligned} \tag{3.75}$$

where

$$A_{lkji} = 2(g_{ij}\nabla_{[l} \pi_{k]} + g_{kj}\nabla_{[i} \pi_{l]} + g_{lj}\nabla_{[k} \pi_{i]}). \tag{3.76}$$

Contracting (3.75) with h and i and using equation (3.74),

$$\begin{aligned} \bar{R}_{jk} + \bar{R}_{kji}{}^i + \bar{R}_{kj} + 2\bar{R}_{ilkj}g^{il} &= 2K(2g_{ik}\delta_j^i - g_{kj}\delta_i^i - g_{ij}\delta_k^i) - g^{li}A_{lkji} \\ &\quad + 4\delta_j^i(\nabla_{[k} w_{i]} + \nabla_{[k} \mu_{i]}). \end{aligned} \tag{3.77}$$

For (3.77), by using (3.27), (3.28), and (2.15), let us calculate $\bar{R}_{kji}{}^i$ and $g^{li}A_{lkji}$:

$$\bar{R}_{kji}{}^i = Q_{kji}{}^i + W_{kji}{}^i, \tag{3.78}$$

where

$$\begin{aligned} Q_{kji}{}^i &= \delta_j^i \lambda_{ki} - \delta_k^i \lambda_{ji} + \lambda_{jl}g^{li}g_{ki} - \lambda_{km}g^{mi}g_{ji} + \delta_i^i(\partial_j \mu_k - \partial_k \mu_j) \\ &= 2n \nabla_{[j} \mu_{k]}. \end{aligned} \tag{3.79}$$

Thus,

$$Q_{kji}{}^i = 2n \nabla_{[j} \mu_{k]},$$

$$W_{kji}{}^i = n(w_{jk} - w_{kj}) = 2n \nabla_{[j} w_{k]}, \tag{3.80}$$

and from (3.76),

$$A_{lkji}g^{li} = 2g^{li}g_{kj}\nabla_{[i}\pi_{l]} . \tag{3.81}$$

Hence, we reach

$$\begin{aligned} \bar{R}_{jk} + \bar{R}_{kj} + 2n(\nabla_{[j}\mu_{k]} + \nabla_{[j}w_{k]}) &= 2K(2g_{ik}\delta_i^j - ng_{kj} - g_{kj}) \\ &+ 4(\nabla_{[k}w_{j]} + \nabla_{[k}\mu_{j]}) - 2g^{li}g_{kj}\nabla_{[i}\pi_{l]} . \end{aligned} \tag{3.82}$$

From (3.82), we observe that the symmetric part of the Ricci tensor is

$$\bar{R}_{(jk)} = (1 - n)Kg_{kj} - (n + 2)(\nabla_{[j}w_{k]} + \nabla_{[j}\mu_{k]}) . \tag{3.83}$$

Since $\bar{R}_{(jk)}$ is symmetric, the second term of (3.83) must satisfy the following relation:

$$\nabla_{[j}w_{k]} + \nabla_{[j}\mu_{k]} = 0 . \tag{3.84}$$

Thus, the symmetric Ricci tensor of $EWS(g, w, \pi, \mu)$ is

$$\bar{R}_{(jk)} = (1 - n)Kg_{kj} , \tag{3.85}$$

and equation (3.84) implies that w_k and μ_k are gradient. This completes the proof. \square

Acknowledgment

The authors would like to thank the referees for their valuable comments and suggestions, which helped to improve the paper.

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