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# Mannheim curves with modified orthogonal frame in Euclidean 3-space 

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#### Abstract

In this paper, we investigate Mannheim pairs, Frenet-Mannheim curves, and weakened Mannheim curves with respect to the modified orthogonal frame in Euclidean 3-space $\left(E^{3}\right)$. We derive some characterization results of these curves.


Key words: Mannheim curves, Mannheim partner curve, Modified orthogonal frame

## 1. Introduction

In the study of the classical differential geometry of space curves, finding the corresponding relations between different space curves has been an important and interesting characterization problem of space curves. For example, if the normal vector of one space curve $\varphi$ is normal to another curve $\psi$, then $\psi$ is called the Bertrand mate of $\varphi$. Liu [4] characterized a similar type of curves, weakened Bertrand curves and Frenet-Bertrand curves under weakened conditions. There is another important class of space curves called Mannheim curves, where the normal vector of one curve is the binormal vector to some other curve, and such a pair of curves is called as Mannheim pair. Liu and Wang [5] derived the necessary and sufficient conditions for a curve to possess a Mannheim partner curve in Euclidean and Minkowski spaces. Öztekin and Ergüt [9] studied the null Mannheim curves in Minkowski space and derived some necessary and sufficient conditions. Recently, Tunçer et al. [11] obtained some characterization results of the nonnull weakened Mannheim curves in Minkowski 3-space. Moreover, Karacan [3] characterized the weakened Mannheim curves in Euclidean 3-space. Mostly, researchers have studied Mannheim curves with respect to the classical Frenet-Serret frame of a curve, where we are supposed to consider that the curvature $\kappa(s) \neq 0$. In this paper, we shall drop the condition of $\kappa(s) \neq 0$ and consider a general set of curves with a discrete set of zeros of $\kappa(s)$ to characterize the Mannheim curves according to the modified orthogonal frame in Euclidean 3-space. Bukcu and Karacan studied space curves [2] and spherical curves [1] with respect to the modified orthogonal frame in Minkowski and Euclidean space, respectively. Recently, Lone et al. [6] obtained some characterization results for helices and Bertrand curves with respect to the modified orthogonal frame.

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## 2. Preliminaries

Let $\varphi(s)$ be a $C^{3}$ space curve in Euclidean 3 -space $E^{3}$, parametrized by arc length $s$. We also assume that its curvature $\kappa(s) \neq 0$ anywhere. Then an orthonormal frame $\{t, n, b\}$ exists satisfying the Frenet-Serret equations

$$
\left[\begin{array}{c}
t^{\prime}(s)  \tag{2.1}\\
n^{\prime}(s) \\
b^{\prime}(s)
\end{array}\right]=\left[\begin{array}{ccc}
0 & \kappa & 0 \\
-\kappa & 0 & \tau \\
0 & -\tau & 0
\end{array}\right]\left[\begin{array}{c}
t(s) \\
n(s) \\
b(s)
\end{array}\right]
$$

where $t$ is the unit tangent, $n$ is the unit principal normal, $b$ is the unit binormal, and $\tau(s)$ is the torsion. For a given $C^{1}$ function $\kappa(s)$ and a continuous function $\tau(s)$, there exists a $C^{3}$ curve $\varphi$ that has an orthonormal frame $\{t, n, b\}$ satisfying the Frenet-Serret frame (2.1). Moreover, any other curve $\tilde{\varphi}$ satisfying the same conditions differs from $\varphi$ only by a rigid motion.

Now let $\varphi(t)$ be a general analytic curve, which can be reparametrized by its arc length. Assuming that the curvature function has discrete zero points or $\kappa(s)$ is not identically zero, we have an orthogonal frame $\{T, N, B\}$ defined as follows:

$$
T=\frac{d \varphi}{d s}, \quad N=\frac{d T}{d s}, \quad B=T \times N
$$

where $T \times N$ is the vector product of $T$ and $N$. The relations between $\{T, N, B\}$ and previous Frenet frame vectors at nonzero points of $\kappa$ are

$$
\begin{equation*}
T=t, N=\kappa n, B=\kappa b \tag{2.2}
\end{equation*}
$$

Thus, we see that $N\left(s_{0}\right)=B\left(s_{0}\right)=0$ when $\kappa\left(s_{0}\right)=0$ and squares of the length of $N$ and $B$ vary analytically in $s$. From Eq. (2.2), it is easy to calculate

$$
\left[\begin{array}{c}
T^{\prime}(s)  \tag{2.3}\\
N^{\prime}(s) \\
B^{\prime}(s)
\end{array}\right]=\left[\begin{array}{ccc}
0 & 1 & 0 \\
-\kappa^{2} & \frac{\kappa^{\prime}}{\kappa} & \tau \\
0 & -\tau & \frac{\kappa^{\prime}}{\kappa}
\end{array}\right]\left[\begin{array}{c}
T(s) \\
N(s) \\
B(s)
\end{array}\right]
$$

and

$$
\tau=\tau(s)=\frac{\operatorname{det}\left(\varphi^{\prime}, \varphi^{\prime \prime}, \varphi^{\prime \prime \prime}\right)}{\kappa^{2}}
$$

is the torsion of $\varphi$. From Frenet-Serret equations, we know that any point where $\kappa^{2}=0$ is a removable singularity of $\tau$. Let $\langle$,$\rangle be the standard inner product of E^{3}$, and then $\{T, N, B\}$ satisfies

$$
\begin{equation*}
\langle T, T\rangle=1,\langle N, N\rangle=\langle B, B\rangle=\kappa^{2},\langle T, N\rangle=\langle T, B\rangle=\langle N, B\rangle=0 \tag{2.4}
\end{equation*}
$$

The orthogonal frame defined in Eq. (2.3) satisfying Eq. (2.4) is called the modified orthogonal frame [10]. We see that for $\kappa=1$, the Frenet--Serret frame coincides with the modified orthogonal frame.

Definition 2.1 Let $\varphi$ and $\psi$ be two space curves in Euclidean 3-space such that, at the corresponding points of the curve, the principal normal vectors of $\varphi$ coincide with the binormal vectors of $\psi$ at the corresponding points. Then $\psi$ is called a Mannheim partner curve of Mannheim curve $\varphi$ and the pair $\{\varphi, \psi\}$ is called a Mannheim pair [5].

Definition 2.2 A Mannheim curve $\psi\left(s^{*}\right)$ is a $C^{\infty}$ regular curve with nonzero curvature for which there exists another $C^{\infty}$ regular curve $\varphi(s)$ with $\varphi^{\prime}(s) \neq 0$ and parameterized by arc length and it also has nonvanishing curvature, in bijection with it in such a way that the principal normal to $\psi\left(s^{*}\right)$ and the binormal to $\varphi(s)$ at each pair of corresponding points coincide with the line joining the corresponding points. The curve $\varphi(s)$ is said to be a Mannheim conjugate of $\psi\left(s^{*}\right)[3,8]$.

Definition 2.3 A Frenet-Mannheim (FM) curve $\psi\left(s^{*}\right)$ is a $C^{\infty}$ Frenet curve for which there exists another $C^{\infty}$ Frenet curve $\varphi(s)$ and $\varphi^{\prime}(s) \neq 0$, in bijection with it so that, by suitable choice of the Frenet frames, the principal normal vector $N_{\psi}\left(s^{*}\right)$ and the binormal vector $B_{\varphi}(s)$ at corresponding points on $\psi\left(s^{*}\right), \varphi(s)$, both lie on the line joining the corresponding points. The curve $\varphi(s)$ is called the FM conjugate of $\psi\left(s^{*}\right)$ [3, 8].

Definition 2.4 A weakened Mannheim (WM) curve $\psi\left(s^{*}\right)$, $s^{*} \in I^{*}$, is a $C^{\infty}$ regular curve for which there exists another $C^{\infty}$ regular curve $\varphi(s), s \in I$, and a homeomorphism $\rho: I \rightarrow I^{*}$, such that:
(i) There exist two disjoint closed subsets $M, N$ of $I$ with void interiors such that $\rho \in C^{\infty}$ on $I \backslash N$, $\frac{d s^{*}}{d s}=0$ on $M, \rho^{-1} \in C^{\infty}$ on $\rho(I \backslash M)$, and $\frac{d s}{d s^{*}}=0$ on $\rho(N)$.
(ii) The line joining the corresponding points $s, s^{*}$ of $\varphi(s)$ and $\psi\left(s^{*}\right)$ is orthogonal to $\varphi(s)$ and $\psi\left(s^{*}\right)$ at the points $s, s^{*}$, respectively, and is along the principal normal to $\varphi(s)$ or $\psi\left(s^{*}\right)$ at the points $s, s^{*}$ whenever it is well defined.

The curve $\varphi(s)$ is called a WM conjugate of $\psi\left(s^{*}\right)[3,8]$.
From classical differential geometry, we find that there is a rich literature available on the Bertrand pairs in comparison to Mannheim curves. Thus, in this paper, we study the Mannheim curves according to the modified orthogonal frame in Euclidean 3-space and obtain several conditions for Mannheim partner FM and WM curves.

## 3. Mannheim partner curves according to modified orthogonal frame in $E^{3}$

Theorem 3.1 Let $C: \varphi(s)$ be a Mannheim curve in $E^{3}$ parameterized by its arc length $s$ and let $C^{*}: \psi\left(s^{*}\right)$ be the Mannheim partner curve of $C$ with an arc length parameter $s^{*}$. The distance between corresponding points of the Mannheim partner curves in $E^{3}$ is $|c| \kappa_{\varphi}$, where $c$ is a nonzero constant and $\kappa_{\varphi}$ is the curvature of curve $\varphi$.

Proof From the definition of Mannheim pair $\left\{C, C^{*}\right\}$, we can write $\overrightarrow{\varphi(s) \psi\left(s^{*}\right)}=\mu(s) N_{\varphi}(s)$, or

$$
\begin{equation*}
\psi\left(s^{*}\right)=\varphi(s)+\mu(s) N_{\varphi}(s) \tag{3.1}
\end{equation*}
$$

for some function $\mu(s)$. Taking the derivative with respect to $s$ and using Eq. (2.3), we get

$$
\psi^{\prime}\left(s^{*}\right)=T_{\varphi}+\mu^{\prime} N_{\varphi}+\mu\left(-\kappa_{\varphi}^{2} T_{\varphi}+\frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}} N_{\varphi}+\tau_{\varphi} B_{\varphi}\right)
$$

or

$$
\begin{equation*}
T_{\psi} \frac{d s^{*}}{d s}=\left(1-\mu \kappa_{\varphi}^{2}\right) T_{\varphi}+\left(\mu^{\prime}+\mu \frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}}\right) N_{\varphi}+\mu \tau_{\varphi} B_{\varphi} \tag{3.2}
\end{equation*}
$$

Taking the inner product of Eq. (3.2) with $B_{\psi}$ and considering $\frac{N_{\varphi}}{\kappa_{\varphi}}=\epsilon \frac{B_{\psi}}{\kappa_{\psi}}(\epsilon= \pm 1)$, we get

$$
\begin{equation*}
\mu^{\prime}+\mu \frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}}=0 \text { or } \mu=\frac{c}{\kappa_{\varphi}} \text {. } \tag{3.3}
\end{equation*}
$$

This means that $\mu$ is not constant except $c=0$. On the other hand, from the distance function between the points of $\varphi(s)$ and $\psi\left(s^{*}\right)$, we have

$$
d\left(\varphi(s), \psi\left(s^{*}\right)\right)=|c| \kappa_{\varphi}(s) .
$$

Theorem 3.2 A space curve $\varphi(s)$ in $E^{3}$ with respect to the modified orthogonal frame is a Mannheim curve if and only if its curvature $\kappa_{\varphi}$ and torsion $\tau_{\varphi}$ satisfy

$$
\begin{equation*}
\kappa_{\varphi}=c\left(\kappa_{\varphi}^{2}+\tau_{\varphi}^{2}\right), \tag{3.4}
\end{equation*}
$$

where $c$ is a nonzero constant.
Proof Let $C: \varphi(s)$ be a Mannheim curve in $E^{3}$ with arc length parameter $s$ and $C^{*}: \psi\left(s^{*}\right)$ the Mannheim partner curve of $C$ with arc length parameter $s^{*}$. Inserting Eq. (3.3) into Eq. (3.1), we get

$$
\begin{equation*}
\psi\left(s^{*}\right)=\varphi(s)+\frac{c}{\kappa_{\varphi}(s)} N_{\varphi}(s) \tag{3.5}
\end{equation*}
$$

for some nonzero constant $c$. Differentiating Eq. (3.5) with respect to $s$ and applying the modified orthogonal frame formulas, we obtain

$$
\begin{equation*}
T_{\psi} \frac{d s^{*}}{d s}=\left(1-c \kappa_{\varphi}\right) T_{\varphi}+\frac{c \tau_{\varphi}}{\kappa_{\varphi}} B_{\varphi} . \tag{3.6}
\end{equation*}
$$

Again differentiating Eq. (3.6) with respect to $s$ and applying the modified orthogonal frame formulas, we get

$$
\begin{aligned}
N_{\psi}\left(\frac{d s^{*}}{d s}\right)^{2}+T_{\psi} \frac{d^{2} s^{*}}{d s^{2}}= & -c \kappa_{\varphi}^{\prime} T_{\varphi}+\left(1-c \kappa_{\varphi}\right) N_{\varphi}+\frac{c \tau_{\varphi}^{\prime} \kappa_{\varphi}-c \kappa_{\varphi}^{\prime} \tau_{\varphi}}{\kappa_{\varphi}^{2}} B_{\varphi} \\
& +\frac{c \tau_{\varphi}}{\kappa_{\varphi}}\left(-\tau_{\varphi} N_{\varphi}+\frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}} B_{\varphi}\right)
\end{aligned}
$$

or

$$
\begin{equation*}
N_{\psi}\left(\frac{d s^{*}}{d s}\right)^{2}+T_{\psi} \frac{d^{2} s^{*}}{d s^{2}}=-c \kappa_{\varphi}^{\prime} T+\frac{1}{\kappa_{\varphi}}\left(\kappa_{\varphi}-c \kappa_{\varphi}^{2}-c \tau_{\varphi}^{2}\right) N_{\varphi}+c \tau_{\varphi}^{\prime} \kappa_{\varphi} B_{\varphi} . \tag{3.7}
\end{equation*}
$$

Taking the inner product of Eq. (3.7) with $B_{\psi}$, we get

$$
\begin{equation*}
\kappa_{\varphi}-c \kappa_{\varphi}^{2}-c \tau_{\varphi}^{2}=0 \text { or } \kappa_{\varphi}=c\left(\kappa_{\varphi}^{2}+\tau_{\varphi}^{2}\right) . \tag{3.8}
\end{equation*}
$$

This completes the proof.
Conversely, if the curvature $\kappa_{\varphi}$ and the torsion $\tau_{\varphi}$ of curve $C$ satisfy Eq. (3.4) for some nonzero constant $c$, then define a curve $C$ by Eq. (3.5), and we will prove that $C$ is Mannheim and $C^{*}$ is the partner curve of $C$. We have already found the equality below:

$$
T_{\psi} \frac{d s^{*}}{d s}=\left(1-c \kappa_{\varphi}\right) T_{\varphi}+\frac{c \tau_{\varphi}}{\kappa_{\varphi}} B_{\varphi} .
$$

Differentiating the last equality with respect to $s$ and with the help of Eq. (3.8), we get

$$
\begin{equation*}
N_{\psi}\left(\frac{d s^{*}}{d s}\right)^{2}+T_{\psi} \frac{d^{2} s^{*}}{d s^{2}}=-c \kappa_{\varphi}^{\prime} T_{\varphi}+\frac{c \tau_{\varphi}^{\prime}}{\kappa_{\varphi}} B_{\varphi} \tag{3.9}
\end{equation*}
$$

Taking the cross product of Eq. (3.6) with Eq. (3.9), we obtain

$$
\frac{d s}{d s^{*}} T_{\psi} \times\left[N_{\psi}\left(\frac{d s^{*}}{d s}\right)^{2}+T_{\psi} \frac{d^{2} s^{* 2}}{d s}\right]=\left[\left(1-c \kappa_{\varphi}\right) T_{\varphi}+\frac{c \tau_{\varphi}}{\kappa_{\varphi}} B_{\varphi}\right] \times\left(-c \kappa_{\varphi}^{\prime} T_{\varphi}+\frac{c \tau_{\varphi}^{\prime}}{\kappa_{\varphi}} B_{\varphi}\right)
$$

or

$$
\begin{equation*}
\left(\frac{d s^{*}}{d s}\right)^{3} B_{\psi}=c\left(-\tau_{\varphi}^{\prime}+c \tau_{\varphi}^{\prime} \kappa_{\varphi}-c \kappa_{\varphi}^{\prime} \tau_{\varphi}\right) \frac{N_{\varphi}}{\kappa_{\varphi}} \tag{3.10}
\end{equation*}
$$

Since both $\frac{N_{\varphi}}{\kappa_{\varphi}}$ and $\frac{B_{\psi}}{\kappa_{\psi}}$ have unit length, we get

$$
\begin{equation*}
\left(\frac{d s^{*}}{d s}\right)^{3}=\frac{c\left(-\tau_{\varphi}^{\prime}+c \tau_{\varphi}^{\prime} \kappa_{\varphi}-c \kappa_{\varphi}^{\prime} \tau_{\varphi}\right)}{\kappa_{\psi}} \tag{3.11}
\end{equation*}
$$

Thus, we have

$$
\frac{B_{\psi}}{\kappa_{\psi}}=\epsilon \frac{N_{\varphi}}{\kappa_{\varphi}}, \epsilon= \pm 1
$$

or $N_{\varphi}$ and $B_{\psi}$ are linearly dependent. This completes the proof.

Theorem 3.3 A pair of curves $\left(C, C^{*}\right)$ is a Mannheim pair if and only if the curvature $\kappa_{\psi}$ and the torsion $\tau_{\psi}$ of curve $C^{*}$ satisfy:

$$
\begin{equation*}
\tau_{\psi}^{\prime}=\frac{d \tau_{\psi}}{d s^{*}}=\frac{\kappa_{\psi}}{a}\left(1+a^{2} \tau_{\psi}^{2}\right) \tag{3.12}
\end{equation*}
$$

where $a$ is a nonzero constant.
Proof Suppose that $C: \varphi(s)$ is a Mannheim curve. By the definition of $\varphi(s)$, we may write

$$
\begin{equation*}
\varphi(s)=\psi\left(s^{*}\right)+\delta\left(s^{*}\right) B_{\psi}\left(s^{*}\right) \tag{3.13}
\end{equation*}
$$

for some function $\delta\left(s^{*}\right)$. Differentiating Eq. (3.13) with respect to $s^{*}$, we get

$$
\begin{equation*}
T_{\varphi} \frac{d s}{d s^{*}}=T_{\psi}-\delta \tau_{\psi} N_{\psi}+\left(\delta^{\prime}+\delta \frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}}\right) B_{\psi} \tag{3.14}
\end{equation*}
$$

Since $N_{\varphi}$ and $B_{\psi}$ are linearly dependent, we get

$$
\begin{equation*}
\delta^{\prime}+\delta \frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}}=0 \text { or } \delta\left(s^{*}\right)=\frac{a}{\kappa_{\psi}} . \tag{3.15}
\end{equation*}
$$

This means that $\delta\left(s^{*}\right)$ is not a constant for each $s^{*}$ except $a=0$. Thus, with the help of Eq. (3.15), we can rewrite Eq. (3.14) as follows:

$$
\begin{equation*}
T_{\varphi} \frac{d s}{d s^{*}}=T_{\psi}-\frac{a \tau_{\psi}}{\kappa_{\psi}} N_{\psi} \tag{3.16}
\end{equation*}
$$

Let $\theta$ be the angle between $T_{\varphi}$ and $T_{\psi}$ at the corresponding points of $C$ and $C^{*}$ in Eq. (3.13). Then, taking the inner product of Eq. (3.16) with $T_{\psi}$ and considering the equality $\cos ^{2} \theta+\sin ^{2} \theta=1$, we get

$$
\begin{equation*}
\cos \theta=\frac{d s^{*}}{d s} \tag{3.17}
\end{equation*}
$$

and

$$
\begin{equation*}
\frac{d s^{*}}{d s}=-a \tau_{\psi} \sin \theta \tag{3.18}
\end{equation*}
$$

From Eq. (3.17) and Eq. (3.18), we find

$$
\begin{equation*}
\frac{d s}{d s^{*}}=\frac{1}{\cos \theta}=-\frac{a \tau_{\psi}}{\sin \theta} \tag{3.19}
\end{equation*}
$$

and

$$
\begin{equation*}
\tan \theta=-a \tau_{\psi} . \tag{3.20}
\end{equation*}
$$

Thus, we can write Eq. (3.16) as follows:

$$
\begin{equation*}
T_{\varphi}=(\cos \theta) T_{\psi}+\frac{\sin \theta}{\kappa_{\psi}} N_{\psi} \tag{3.21}
\end{equation*}
$$

Differentiating Eq. (3.21) with respect to $s^{*}$, we get

$$
\begin{equation*}
N_{\varphi} \frac{d s}{d s^{*}}=-\sin \theta\left(\kappa_{\psi}+\theta^{\prime}\right) T_{\psi}+\frac{\cos \theta}{\kappa_{\psi}}\left(\kappa_{\psi}+\theta^{\prime}\right) N_{\psi}+\left(\frac{\tau_{\psi}}{\kappa_{\psi}} \sin \theta\right) B_{\psi} \tag{3.22}
\end{equation*}
$$

From this equation and the fact that the direction of $\frac{N_{\varphi}}{\kappa_{\varphi}}$ is coincident with $\frac{B_{\psi}}{\kappa_{\psi}}$, we get

$$
\left\{\begin{array}{c}
-\sin \theta\left(\kappa_{\psi}+\theta^{\prime}\right)=0  \tag{3.23}\\
\frac{\cos \theta}{\kappa_{\psi}}\left(\kappa_{\psi}+\theta^{\prime}\right)=0
\end{array}\right.
$$

or

$$
\begin{equation*}
\theta^{\prime}=-\kappa_{\psi} . \tag{3.24}
\end{equation*}
$$

Differentiating Eq. (3.20) with respect to $s^{*}$ and applying Eq. (3.24), we get

$$
\kappa_{\psi}+a^{2} \kappa_{\psi} \tau_{\psi}^{2}-a \tau_{\psi}^{\prime}=0
$$

or

$$
\tau_{\psi}^{\prime}=\frac{\kappa_{\psi}}{a}\left(1+a^{2} \tau_{\psi}^{2}\right)
$$

Conversely, if the curvature $\kappa_{\psi}$ and the torsion $\tau_{\psi}$ of $C^{*}$ satisfy Eq. (3.12) for some nonzero constant $a$, then define a curve $C$ by Eq. (3.13) and we will prove that $C$ is a Mannheim and $C^{*}$ is the partner curve of $C$. We can easily reduce Eq. (3.13) in the following expression:

$$
T_{\varphi} \frac{d s}{d s^{*}}=T_{\psi}-\frac{a \tau_{\psi}}{\kappa_{\psi}} N_{\psi}
$$

Differentiating the above equality with respect to $s^{*}$ and with the help of Eq. (2.3), we get

$$
N_{\varphi}\left(\frac{d s}{d s^{*}}\right)^{2}+T_{\varphi} \frac{d^{2} s}{d s^{* 2}}=N_{\psi}+\frac{a \kappa_{\psi}^{\prime} \tau_{\psi}-a \kappa_{\psi} \tau_{\psi}^{\prime}}{\kappa_{\psi}^{2}} N_{\psi}-\frac{a \tau_{\psi}}{\kappa_{\psi}}\left(-\kappa_{\psi}^{2} T_{\psi}+\frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}} N_{\psi}+\tau_{\psi} B_{\psi}\right)
$$

or noticing Eq. (3.12), we get

$$
\begin{equation*}
N_{\varphi}\left(\frac{d s}{d s^{*}}\right)^{2}+T_{\varphi} \frac{d^{2} s}{d s^{* 2}}=a \tau_{\psi} \kappa_{\psi} T_{\psi}-a^{2} \tau_{\psi}^{2} N_{\psi}+\frac{c \tau_{\psi}^{2}}{\kappa_{\psi}} B_{\psi} \tag{3.25}
\end{equation*}
$$

Taking the cross product of Eq. (3.16) with Eq. (3.25), we have

$$
\frac{d s}{d s^{*}} T_{\varphi} \times\left[\left(\frac{d s}{d s^{*}}\right)^{2} N_{\varphi}+T_{\varphi} \frac{d^{2} s}{d s^{* 2}}\right]=\left(T_{\psi}-\frac{a \tau_{\psi}}{\kappa_{\psi}} N_{\psi}\right) \times\left(a \tau_{\psi} \kappa_{\psi} T_{\psi}-a^{2} \tau_{\psi}^{2} N_{\psi}+\frac{a \tau_{\psi}^{2}}{\kappa_{\psi}} B_{\psi}\right)
$$

or

$$
\begin{equation*}
\left(\frac{d s}{d s^{*}}\right)^{3} B_{\varphi}=-\frac{a \tau_{\psi}^{2}}{\kappa_{\psi}}\left(\frac{a \tau_{\psi}}{\kappa_{\psi}} T_{\psi}+N_{\psi}\right) \tag{3.26}
\end{equation*}
$$

Again taking the cross product of Eq. (3.16) with Eq. (3.26), we obtain

$$
\left(\frac{d s}{d s^{*}}\right)^{4} N_{\varphi}=\frac{a \tau_{\psi}^{2}}{\kappa_{\psi}}\left(\kappa_{\psi}^{2}+a^{2} \tau_{\psi}^{2}\right) B_{\psi}
$$

or

$$
N_{\varphi}=\frac{\kappa_{\varphi}}{\kappa_{\psi}} B_{\psi}
$$

This means that the principal normal direction $\frac{N_{\varphi}}{\kappa_{\varphi}}$ of $C: \varphi(s)$ coincides with the binormal direction $\frac{B_{\psi}}{\kappa_{\psi}}$ of $C^{*}: \psi\left(s^{*}\right)$. Hence, $C: \varphi(s)$ is a Mannheim curve and $C^{*}: \psi\left(s^{*}\right)$ is its Mannheim partner curve. Therefore, for each Mannheim curve, there is a unique Mannheim partner curve.

Proposition 3.4 A simple parametric transformation reduces the condition

$$
\tau_{\psi}^{\prime}=\frac{\kappa_{\psi}}{a}\left(1+a^{2} \tau_{\psi}^{2}\right)
$$

to

$$
\tau_{\psi}=\frac{1}{a} \tan \left(\int \kappa_{\psi} d s+c_{0}\right)
$$

Thus, the existence of a Mannheim partner curve to a Mannheim curve is unique.

Proposition 3.5 Let $\left\{\varphi(s), \psi\left(s^{*}\right)\right\}$ be a Mannheim pair, where both $\varphi$ and $\psi$ are parameterized by arc length $s$ and $s^{*}$, respectively. If $\varphi(s)$ is a generalized helix according to the modified frame in $E^{3}$, then $\psi\left(s^{*}\right)$ is a straight line.

Proof Let $T_{\varphi}, N_{\varphi}$, and $B_{\varphi}$ be the tangent, the principal normal, and the binormal vectors of $\varphi(s)$, respectively. From the definition of the Mannheim curve and properties of generalized helices, we have

$$
\left\langle N_{\varphi}, u\right\rangle=\left\langle B_{\psi}, u\right\rangle=0
$$

where $u$ is some constant vector. Differentiating the last equality, we obtain

$$
\left\langle-\tau_{\psi} N_{\psi}+\frac{\kappa^{\prime}}{\kappa} B_{\psi}, u\right\rangle=-\tau_{\psi}\left\langle N_{\psi}, u\right\rangle+\frac{\kappa^{\prime}}{\kappa}\left\langle B_{\psi}, u\right\rangle=-\tau_{\psi}\left\langle N_{\psi}, u\right\rangle=0
$$

Since $\left\langle N_{\psi}, u\right\rangle \neq 0$, we get

$$
\tau_{\psi}=0
$$

Thus, using the last equality in Eq. (3.12), we easily obtain

$$
\kappa_{\psi}=0
$$

## 4. Frenet-Mannheim curves according to modified orthogonal frame in $E^{3}$

In this section, we characterize $F M$ curves. For that, we begin with a lemma.

Lemma 4.1 Suppose $\psi\left(s^{\star}\right)$, $s^{\star} \in I^{\star}$, is a FM curve with FM conjugate $\varphi(s)$. We mark all the quantities of $\psi\left(s^{\star}\right)$ with an asterisk and suppose

$$
\begin{equation*}
\psi\left(s^{\star}\right)=\varphi(s)+\delta(s) B_{\varphi}(s) \tag{4.1}
\end{equation*}
$$

Then the distance $(|\delta|)$ between corresponding points of $\varphi(s), \psi\left(s^{\star}\right)$ is not constant, i.e. $\delta=c \kappa_{\varphi}, c \in R$, and $\left\langle T_{\varphi}, T_{\psi}\right\rangle=\cos \theta$, where $\theta$ is a constant angle and
(i) $\sin \theta=-a \tau_{\varphi} \cos \theta$,
(ii) $\left(1+\epsilon a \kappa_{\psi}\right) \sin \theta=a \tau_{\psi} \cos \theta$,
(iii) $\cos ^{2} \theta=1+\epsilon a \kappa_{\psi}$,
(iv) $\sin ^{2} \theta=a^{2} \tau_{\varphi} \tau_{\psi}$.

Proof From Eq. (4.1), we have

$$
\delta(s)=\left\langle\psi\left(s^{\star}\right)-\varphi(s), B_{\varphi}(s)\right\rangle
$$

where $\delta(s)$ is of class $C^{\infty}$. Differentiating Eq. (4.1) with respect to $s$, we get

$$
T_{\psi} \frac{d s^{\star}}{d s}=T_{\varphi}+\delta^{\prime} B_{\varphi}+\delta\left(-\tau_{\varphi} N_{\varphi}+\frac{\kappa^{\prime}}{\kappa} B_{\varphi}\right)
$$

or

$$
\begin{equation*}
T_{\psi} \frac{d s^{\star}}{d s}=T_{\varphi}-\tau_{\varphi} \delta N_{\varphi}+\left(\delta^{\prime}+\delta \frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}}\right) B_{\varphi} \tag{4.2}
\end{equation*}
$$

By the given conditions, we have $B_{\varphi}=\epsilon N_{\psi}$ with $\epsilon= \pm 1$. Taking the scalar multiplication of Eq. (4.2) with $B_{\varphi}$, we obtain

$$
\frac{\delta^{\prime}}{\delta}=-\frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}} \Rightarrow \delta=\frac{a}{\kappa_{\varphi}}, a \in R
$$

Therefore,

$$
\begin{equation*}
T_{\psi} \frac{d s^{\star}}{d s}=T_{\varphi}-\frac{a \tau_{\varphi}}{\kappa_{\varphi}} N_{\varphi} \tag{4.3}
\end{equation*}
$$

Now, by the definition of the $F M$ curve, we have $\frac{d s^{\star}}{d s} \neq 0$, so $T_{\psi}$ is a function of $s$ of class $C^{\infty}$. Hence,

$$
\left\langle T_{\varphi}, T_{\psi}\right\rangle^{\prime}=\left\langle N_{\varphi}, T_{\psi}\right\rangle+\left\langle T_{\varphi}, N_{\psi}\right\rangle \frac{d s^{\star}}{d s}=0
$$

This implies that $\left\langle T_{\varphi}, T_{\psi}\right\rangle$ is constant and thus there exists a constant angle $\theta$, such that

$$
\begin{equation*}
T_{\psi}=T_{\varphi} \cos \theta+N_{\varphi} \frac{\sin \theta}{\kappa_{\varphi}} \tag{4.4}
\end{equation*}
$$

From Eq. (4.3) and Eq. (4.4), we get

$$
\left(\frac{d s}{d s^{\star}}-\cos \theta\right) T_{\psi}-\left(\frac{d s}{d s^{\star}} \frac{a \cdot \tau_{\varphi}}{\kappa_{\varphi}}+\frac{\sin \theta}{\kappa_{\varphi}}\right) N_{\psi}=0 .
$$

Since $T_{\psi}$ and $N_{\psi}$ are linearly independent vectors, we have

$$
\begin{equation*}
\frac{d s}{d s^{\star}}=\cos \theta \tag{4.5}
\end{equation*}
$$

and

$$
\frac{d s^{\star}}{d s} \sin \theta=-a \tau_{\varphi} .
$$

Using Eq. (4.5) in the last equality, we get

$$
\begin{equation*}
\sin \theta=-a \tau_{\varphi} \cos \theta \tag{4.6}
\end{equation*}
$$

which is $(i)$. Now write

$$
\varphi(s)=\psi\left(s^{\star}\right)-\epsilon \delta(s) N_{\psi}(s)
$$

The above equation implies that

$$
\begin{aligned}
& T_{\varphi}=\frac{d s^{\star}}{d s}\left[T_{\psi}-\epsilon \delta^{\prime} N_{\psi}-\epsilon \delta\left(-\kappa_{\psi}^{2} T_{\psi}+\frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}} N_{\psi}+\tau_{\psi} B_{\psi}\right)\right] \\
& \text { or } \\
& T_{\varphi}=\frac{d s^{\star}}{d s}\left[\left(1+\epsilon \delta \kappa_{\psi}^{2}\right) T_{\psi}-\epsilon\left(\delta^{\prime}+\delta \frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}}\right) N_{\psi}-\epsilon \delta \tau_{\psi} B_{\psi}\right]
\end{aligned}
$$

Using $\frac{\delta^{\prime}}{\delta}=-\frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}}$, it follows that

$$
\begin{equation*}
T_{\varphi}=\frac{d s^{\star}}{d s}\left[\left(1+\epsilon a \kappa_{\psi}\right) T_{\psi}-\epsilon a \frac{\tau_{\psi}}{\kappa_{\psi}} B_{\psi}\right] \tag{4.7}
\end{equation*}
$$

On the other hand, Eq. (4.4) gives

$$
\begin{aligned}
T_{\psi} \wedge N_{\psi} & =\left[T_{\varphi} \cos \theta+N_{\varphi} \frac{\sin \theta}{\kappa_{\varphi}}\right] \wedge N_{\psi}=\left[T_{\varphi} \cos \theta+N_{\varphi} \frac{\sin \theta}{\kappa_{\varphi}}\right] \wedge\left(\epsilon B_{\varphi}\right) \\
& =\epsilon\left(T_{\varphi} \wedge B_{\varphi}\right) \cos \theta+\epsilon\left(N_{\varphi} \wedge B_{\varphi}\right) \frac{\sin \theta}{\kappa_{\varphi}}=-\epsilon N_{\varphi} \cos \theta+\epsilon T_{\varphi} \kappa_{\varphi} \sin \theta \\
& =-\epsilon N_{\varphi} \cos \theta+\epsilon T_{\varphi} \kappa_{\varphi}^{2} \frac{\sin \theta}{\kappa_{\varphi}}=\epsilon T_{\varphi} \sin \theta-\epsilon N_{\varphi} \kappa_{\varphi} \cos \theta \\
\Rightarrow B_{\psi} & =\epsilon T_{\varphi} \sin \theta-\epsilon N_{\varphi} \cos \theta
\end{aligned}
$$

Using Eq. (4.5) again, we get

$$
\begin{equation*}
T_{\varphi}=T_{\psi} \cos \theta-\epsilon B_{\psi} \frac{\sin \theta}{\kappa_{\psi}} \tag{4.8}
\end{equation*}
$$

Taking the vector product of Eq. (4.5) and Eq. (4.6), we obtain

$$
\left(1+\epsilon a \kappa_{\psi}\right) \sin \theta=a \tau_{\psi} \cos \theta
$$

which is (ii). On the other hand, from Eq. (4.7) and Eq. (4.8), it follows that

$$
\begin{align*}
& \frac{d s^{\star}}{d s}\left(1+a \epsilon \kappa_{\psi}\right)=\cos \theta  \tag{4.9}\\
& \frac{d s^{\star}}{d s}\left(a \tau_{\psi}\right)=\sin \theta \tag{4.10}
\end{align*}
$$

Thus, inserting Eq. (4.5) into Eq. (4.9) and using Eq. (4.6) in Eq. (4.10), we get (iii) and (iv), respectively.

Theorem 4.2 Let $\psi\left(s^{\star}\right) \in C^{\infty}$, $s^{\star} \in I^{\star}$, be a Frenet curve with $\tau_{\psi}$ nowhere vanishing and satisfying

$$
\begin{equation*}
\left(1+a \epsilon \kappa_{\psi}\right) \sin \theta=a \tau_{\psi} \cos \theta \tag{4.11}
\end{equation*}
$$

for some constant $a \neq 0$. Then $\psi\left(s^{\star}\right)$ is a FM curve, which is nonplanar.
Proof Define the position vector of curve $\psi\left(s^{\star}\right)$ as follows:

$$
\psi\left(s^{\star}\right)=\varphi(s)+\frac{a}{\kappa(s)} B_{\varphi}(s)
$$

Let us denote differentiation with respect to $s$ by a dash. We have

$$
\psi^{\prime}\left(s^{*}\right)=T_{\varphi}-\frac{a \tau_{\varphi}}{\kappa_{\varphi}} N_{\varphi}
$$

Since $\tau_{\varphi} \neq 0$, we see that $\psi\left(s^{\star}\right)$ is a $C^{\infty}$ regular curve. Suppose all the quantities of $\psi\left(s^{\star}\right)$ are marked by an asterisk. Then

$$
T_{\psi} \frac{d s^{\star}}{d s}=T_{\varphi}-\frac{a \tau_{\varphi}}{\kappa_{\varphi}} N_{\varphi}
$$

Hence, we have

$$
\frac{d s^{\star}}{d s}=\sqrt{1-a^{2} \tau_{\varphi}^{2}}
$$

Using Eq. (4.11), we get

$$
T_{\psi}=T_{\varphi} \cos \theta+N_{\varphi} \frac{\sin \theta}{\kappa_{\varphi}}
$$

and notice that from Eq. (4.11), we have $\sin \theta \neq 0$. Therefore,

$$
\frac{d T_{\psi}}{d s^{\star}} \frac{d s^{\star}}{d s}=N_{\varphi} \cos \theta+\frac{\sin \theta}{\kappa_{\varphi}}\left(-\kappa_{\varphi}^{2} T_{\varphi}+\frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}} N_{\varphi}+\tau_{\varphi} B_{\varphi}\right)-\frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}^{2}} \sin \theta N_{\varphi}
$$

or

$$
\begin{equation*}
\kappa_{\psi} \frac{d s^{\star}}{d s} \frac{N_{\psi}}{\kappa_{\psi}}=-\left(\kappa_{\varphi} \sin \theta\right) T_{\varphi}+N_{\varphi} \cos \theta+\left(\tau_{\varphi} \sin \theta\right) \frac{B_{\varphi}}{\kappa_{\varphi}} \tag{4.12}
\end{equation*}
$$

Now define $\frac{N_{\psi}}{\kappa_{\psi}}=\epsilon \frac{B_{\varphi}}{\kappa_{\varphi}}$. Then, from Eq. (4.12), we get

$$
\kappa_{\psi}=\epsilon \frac{d s}{d s^{\star}} \tau_{\varphi} \sin \theta
$$

These are $C^{\infty}$ functions of $s$ (and hence of $s^{\star}$ ), and

$$
\frac{d T_{\psi}}{d s^{\star}}=N_{\psi}
$$

Again, define $B_{\psi}=T_{\psi} \wedge B_{\varphi}$ and

$$
\begin{aligned}
& \left\langle\frac{d B_{\psi}}{d s^{\star}}, N_{\psi}\right\rangle=\left\langle-\tau_{\psi} N_{\psi}+\frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}} B_{\psi}, N_{\psi}\right\rangle=-\tau_{\psi}\left\langle N_{\psi}, N_{\psi}\right\rangle=-\tau_{\psi} \kappa_{\psi}^{2} \\
& \text { or } \\
& \tau_{\psi}=-\frac{\left\langle\frac{d B_{\psi}}{d s^{\star}}, N_{\psi}\right\rangle}{\kappa_{\psi}^{2}}
\end{aligned}
$$

These are also $C^{\infty}$ functions on $I^{\star}$. It is then easy to verify that with the modified frame $\left\{T_{\psi}, N_{\psi}, B_{\psi}\right\}$ and the functions $\kappa_{\psi}, \tau_{\psi}$, the curve $\psi\left(s^{\star}\right)$ becomes a $C^{\infty}$ Frenet curve. However, $B_{\varphi}$ and $N_{\psi}$ lie on the line joining the corresponding points of $\varphi(s)$ and $\psi\left(s^{\star}\right)$. Thus, $\psi\left(s^{\star}\right)$ is a FM curve and $\varphi(s)$ is a $F M$ conjugate of $\psi\left(s^{\star}\right)$.

Lemma 4.3 A necessary and sufficient condition for a regular curve $\psi \in C^{\infty}$ to be a FM curve with a FM conjugate is that $\psi$ is either a line or a nonplanar circular helix.

Proof $\Rightarrow$ : Suppose a line $\varphi$ is a $F M$ conjugate of $\psi$. This implies $\kappa_{\varphi}=0$. Using Lemma 4.1, (iii) and (i), (ii), we have

$$
\begin{equation*}
\cos ^{2} \theta=1+a \epsilon \kappa_{\psi} \tag{4.13}
\end{equation*}
$$

and then

$$
\begin{equation*}
\cos ^{2} \theta \sin \theta=a \tau_{\psi} \cos \theta \tag{4.14}
\end{equation*}
$$

$$
\begin{equation*}
\cos \theta=-a \tau_{\varphi} \sin \theta \tag{4.15}
\end{equation*}
$$

From Eq. (4.15), it follows that $\cos \theta \neq 0$. Hence, Eq. (4.14) is proportional to

$$
\begin{equation*}
a \tau_{\psi}=\cos \theta \sin \theta \tag{4.16}
\end{equation*}
$$

Case 1. $\sin \theta=0$. Then $\cos \theta= \pm 1$, so (4.13) implies $\kappa_{\psi}=0$, and $\psi$ is a line. We also note that Eq. (4.16) implies that $\tau_{\psi}=0$.

Case 2. $\sin \theta \neq 0$. Then $\cos \theta \neq \pm 1$, and Eqn. (4.13) and (4.16) imply that $\kappa_{\psi}, \tau_{\psi}$ are nonvanishing constants, and $\psi$ is a nonplanar circular helix.
$\Leftarrow$ : Suppose that $\psi$ is a nonplanar circular helix given by

$$
\psi=(r \cos t, r \sin t, b t),\left(a, b \in R^{+}\right)
$$

where $t=\frac{s}{\sqrt{r^{2}+b^{2}}}$.
We may write

$$
\psi^{\prime}(s)=T_{\psi}=\frac{r}{\sqrt{r^{2}+b^{2}}}(-\sin t, \cos t, k)
$$

and

$$
\begin{gathered}
\psi^{\prime \prime}(s)=N_{\psi}=\frac{-r}{r^{2}+b^{2}}(\cos t, \sin t, 0) \\
\Rightarrow \kappa_{\psi}=\frac{r}{r^{2}+b^{2}} \text { and }-r \frac{N_{\psi}}{\kappa_{\psi}}=r(\cos t, \sin t, 0)
\end{gathered}
$$

Then the curve $\psi$ with

$$
\begin{aligned}
\psi & =(r \cos t, r \sin t, b t)=r(\cos t, \sin t, 0)+(0,0, b t) \\
& =-\frac{r}{\kappa_{\psi}} N_{\psi}+\varphi=\varphi-\frac{r}{\kappa_{\psi}}\left(\frac{\kappa_{\psi}}{\kappa_{\varphi}} B_{\varphi}\right)=\varphi-\frac{r}{\kappa_{\varphi}} B_{\varphi}
\end{aligned}
$$

or putting $\delta=-\frac{r}{\kappa_{\varphi}}$,

$$
\psi=\varphi+\delta B_{\varphi}
$$

will be a line along the $z$-axis and can be turned into a $F M$ conjugate of $\psi$ by defining $\frac{N_{\psi}}{\kappa_{\psi}}$ to be equal to $\frac{B_{\varphi}}{K_{\varphi}}$.

Theorem 4.4 Let $\psi\left(s^{\star}\right) \in C^{\infty}$ be a plane Frenet curve with zero torsion and the curvature being bounded above or bounded below. Then $\psi$ is a FM curve, with FM conjugates that are curves in the plane.

Proof Suppose $\psi$ is a curve satisfying the given conditions. Then $\kappa_{\psi}<-\frac{1}{a}$ or $\kappa_{\psi}>-\frac{1}{a}$ on $I$ for some constant $a \neq 0$. For such $a$, let $\varphi$ be the curve with position vector

$$
\varphi=\psi-\varepsilon \delta N_{\psi}
$$

Differentiating the last equality with respect to $s^{*}$ and considering equality $\delta^{\prime}+\delta \frac{\kappa_{\varphi}^{\prime}}{\kappa_{\varphi}}=0$ and $\delta=\frac{a}{\kappa_{\varphi}}(a \in R)$, we get

$$
T_{\varphi} \frac{d s}{d s^{*}}=T_{\psi}-\varepsilon \delta^{\prime} N_{\psi}-\varepsilon \delta\left(-\kappa_{\psi}^{2} T_{\psi}+\frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}} N_{\psi}\right)=\left(1+\varepsilon \delta \kappa_{\psi}^{2}\right) T_{\psi}-\varepsilon\left(\delta^{\prime}+\delta \frac{\kappa_{\psi}^{\prime}}{\kappa_{\psi}}\right) N_{\psi}
$$

or

$$
T_{\varphi}=\frac{d s^{*}}{d s}\left(1+\varepsilon a \kappa_{\psi}\right) T_{\psi}
$$

Since $\frac{d s^{*}}{d s}\left(1+\varepsilon a \kappa_{\psi}\right) \neq 0, \varphi$ is a $C^{\infty}$ regular curve, and $T_{\varphi}=T_{\psi}$. It is then easy to verify that $\varphi$ is a $F M$ conjugate of $\psi$.

## 5. Weakened Mannheim curves according to modified orthogonal frame in $E^{3}$

Definition 5.1 Let $\mathcal{D}$ be a subset of a topological space $\mathcal{X}$. If a function $\mathcal{X} \mapsto \mathcal{Y}$ is constant for each component of $\mathcal{D}$, we say that the function is $\mathcal{D}$-piecewise constant [4].

Lemma 5.2 Suppose $\mathcal{D}$ is an open subset of a proper interval $\mathcal{X}$ of the real line. A necessary and sufficient condition for every $\mathcal{D}$-piecewise constant real continuous function on $\mathcal{X}$ to be constant is that there is an empty dense-in-itself kernel of $\mathcal{X} \backslash \mathcal{D}$ [4].

We remark, however, that if $\mathcal{D}$ is dense in $\mathcal{X}$, any $\mathcal{D}$-piecewise constant, $C^{1}$ real function on $\mathcal{X}$ is constant, even if $\mathcal{D}$ has a nonempty dense-in-itself kernel.

Theorem 5.3 A WM curve for which $M$ and $N$ (defined in def. 2.4) have void dense-in-itself kernels is a FM curve.

Proof Let $\varphi(s), s \in I$, be a $W M$ conjugate of $W M$ curve $\psi\left(s^{\star}\right), s^{\star} \in I^{\star}$. From the definition of $\varphi(s)$ and $\psi\left(s^{\star}\right)$, it follows that both the curves have a $C^{\infty}$ family of the tangent vectors $T_{\psi}\left(s^{\star}\right), T_{\varphi}(s)$. Let

$$
\begin{equation*}
\psi(s)=\psi(\rho(s))=\varphi(s)+\delta(s) B_{\varphi}(s) \tag{5.1}
\end{equation*}
$$

where $B_{\varphi}(s)$ is a vector function and $\delta(s) \geq 0, \forall s \in I$. Let $\mathcal{D}=I \backslash N, \mathcal{D}^{\star}=I^{\star} \backslash \rho(M)$. Then $s^{\star}(s) \in C^{\infty}$ on $D^{\star}$.

Step 1. Prove $\delta=\frac{a}{\kappa_{\varphi}(s)}$. Since $\delta(=\|\psi(s)-\varphi(s)\|)$ is $C^{\infty}$ on $I$ and nowhere zero on every interval of $\mathcal{D}$, let $\mathcal{X}$ be any component of $\mathcal{P}:=\{s \in I: \delta(s) \neq 0\}$. Hence, both $\mathcal{P}$ and $\mathcal{X}$ are open in $I$. Consider $\mathcal{L}$ as a component interval of $\mathcal{X} \cap \mathcal{D}$. Then $\delta(s)$ and $B_{\varphi}(s)$ are of class $C^{\infty}$ on $\mathcal{L}$, and from Eq. (5.1), we have

$$
\psi^{\prime}(s)=T_{\varphi}+\delta^{\prime}(s) B_{\varphi}(s)+\delta(s) B_{\varphi}^{\prime}(s)
$$

By the definition of a $W M$ curve, we know that $\left\langle T_{\varphi}, B_{\varphi}(s)\right\rangle=0=\left\langle\psi^{\prime}\left(s^{\star}\right), B_{\varphi}(s)\right\rangle$. Hence, using $\left\langle B_{\varphi}^{\prime}(s), B_{\varphi}(s)\right\rangle=$ 0 , we get

$$
0=\left(\delta^{\prime}(s)+\delta(s) \frac{\kappa_{\varphi}^{\prime}(s)}{\kappa_{\varphi}(s)}\right)\left\langle B_{\varphi}(s), B_{\varphi}(s)\right\rangle
$$

Therefore, $\delta=\frac{a}{\kappa_{\varphi}}$ on $\mathcal{L}$, where $a$ is constant. Thus, $\delta$ is not constant on each interval of the set $\mathcal{X} \cap \mathcal{D}$, but by the given conditions $\mathcal{X} \backslash \mathcal{D}$ has a void dense-in-itself kernel. It follows from Lemma 5.2 that $\delta$ is not constant (and nonzero) on $\mathcal{X}$. As $\delta$ is continuous on $I, \mathcal{X}$ should be closed in $I$, but in $I, \mathcal{X}$ is also open. Hence, by connectedness, $\mathcal{X}=I$, i.e. $\delta$ is not constant on $I$.

Step 2. Take the existence of frames

$$
\left\{T_{\varphi}(s), N_{\varphi}(s), N_{\varphi}(s)\right\},\left\{T_{\psi}\left(s^{\star}\right), N_{\psi}\left(s^{\star}\right), B_{\psi}\left(s^{\star}\right)\right\},
$$

which are the modified orthogonal frames for $\varphi(s)$ on $\mathcal{D}$ and $\psi\left(s^{\star}\right)$ on $\mathcal{D}^{\star}$, respectively. Since $\delta=\frac{a}{\kappa_{\varphi}}$ is a nonzero function, it follows from Eq. (5.1) that $B_{\varphi}(s)$ is continuous on $I$ and is of class $C^{\infty}$ on $\mathcal{D}$, and orthogonal to $T_{\varphi}(s)$. Define $B_{\varphi}(s)=T_{\varphi}(s) \wedge N_{\varphi}(s)$. Then $\left\{T_{\varphi}(s), N_{\varphi}(s), B_{\varphi}(s)\right\}$ forms a right-handed modified orthonormal frame for $\varphi(s)$, which is of class $C^{\infty}$ on $\mathcal{D}$ and continuous on $I$.

From the definition of a $W M$ curve, it follows that there exists a scalar function $\kappa_{\psi}\left(s^{\star}\right)$ such that $T_{\psi}^{\prime}\left(s^{\star}\right)=\kappa_{\psi}\left(s^{\star}\right) N_{\psi}\left(s^{\star}\right)$ on $I^{\star}$. Thus, $\left\langle T_{\psi}^{\prime}\left(s^{\star}\right), N_{\psi}\left(s^{\star}\right)\right\rangle=\kappa_{\psi}^{2}\left(s^{\star}\right)$ is continuous on $I^{\star}$ and of class $C^{\infty}$ on $\mathcal{D}^{\star}$. Hence, the first Frenet formula holds on $\mathcal{D}^{\star}$. It is now easy to show that there exists a function $\tau_{\varphi}(s)$ of class $C^{\infty}$ on $\mathcal{D}$ such that the Frenet formulas hold. Thus, $\left\{T_{\varphi}(s), N_{\varphi}(s), B_{\varphi}(s)\right\}$ is a modified orthogonal frame for $\varphi(s)$ on $\mathcal{D}$.

Similarly, there exists a modified orthogonal frame $\left\{T_{\psi}\left(s^{\star}\right), N_{\psi}\left(s^{\star}\right), B_{\psi}\left(s^{\star}\right)\right\}$ for $\psi\left(s^{\star}\right)$, which is continuous on $I^{\star}$ and is a Frenet frame for $\psi\left(s^{\star}\right)$ on $\mathcal{D}^{\star}$. Moreover, we may choose

$$
B_{\varphi}(s)=N_{\psi}(\rho(s))
$$

Step 3. Show that $N=\varnothing, M=\varnothing$.
Noticing that on $\mathcal{D}$, we have

$$
\left\langle T_{\psi}, T_{\varphi}\right\rangle^{\prime}=\left\langle N_{\psi} \frac{d s^{\star}}{d s}, T_{\varphi}\right\rangle+\left\langle T_{\psi}, \kappa_{\varphi} N_{\varphi}\right\rangle=0
$$

on each component of $\mathcal{D},\left\langle T_{\psi}, T_{\varphi}\right\rangle$ is constant and hence on $I$ by Lemma 5.2. Thus, there exists an angle $\theta$ such that

$$
T_{\psi}=T_{\varphi} \cos \theta+N_{\varphi} \sin \theta
$$

Furthermore,

$$
B_{\varphi}(s)=N_{\psi}(\rho(s))
$$

and so

$$
B_{\psi}\left(s^{\star}\right)=-T_{\varphi} \sin \theta+\frac{N_{\varphi}}{\kappa_{\varphi}} \cos \theta
$$

Hence, $\left\{T_{\psi}\left(s^{\star}\right), N_{\psi}\left(s^{\star}\right), B_{\varphi}(s)\right\}$ are also of class $C^{\infty}$ on $\mathcal{D}$. On the other hand, with respect to $s^{\star}$ on $\mathcal{D}^{\star}$, $\left\{T_{\psi}\left(s^{\star}\right), N_{\psi}\left(s^{\star}\right), B_{\psi}\left(s^{\star}\right)\right\}$ are of class $C^{\infty}$. Writing Eq. (5.1) in the form

$$
\varphi=\psi-\frac{a}{\kappa_{\psi}} N_{\psi} \text { or } \varphi=\psi-\delta N_{\psi}
$$

and differentiating with respect to $s$ on $\mathcal{D} \cap \rho^{-1}\left(\mathcal{D}^{\star}\right)$, we have

$$
\begin{equation*}
T_{\varphi}=\frac{d s^{\star}}{d s}\left[\left(1+a \kappa_{\psi}\right) T_{\psi}-\frac{a \tau_{\psi}}{\kappa_{\psi}} B_{\psi}\right] \tag{5.2}
\end{equation*}
$$

However,

$$
T_{\varphi}=T_{\psi} \cos \theta-B_{\psi} \sin \theta
$$

Hence,

$$
\begin{equation*}
\frac{d s^{\star}}{d s}\left(1+a \kappa_{\psi}\right)=\cos \theta \text { and } \frac{a \tau_{\psi}}{\kappa_{\psi}}=-\sin \theta \tag{5.3}
\end{equation*}
$$

Since $\kappa_{\psi}^{2}\left(s^{\star}\right)=\left\langle T_{\psi}^{\prime}, N_{\psi}\right\rangle$ is continuous on $I^{\star}$ and $\rho^{-1}\left(\mathcal{D}^{\star}\right)$ is dense, it follows by continuity that Eq. (5.3) holds throughout $\mathcal{D}$.

Case 1. $\cos \theta \neq 0$. Then Eq. (5.3) implies $\frac{d s^{\star}}{d s} \neq 0$ on $\mathcal{D}$. Hence, $M=\varnothing$. Similarly, $N=\varnothing$.
Case 2. $\cos \theta=0$. Then

$$
\begin{equation*}
T_{\psi}= \pm \frac{N_{\varphi}}{\kappa_{\varphi}} \tag{5.4}
\end{equation*}
$$

Taking the derivative of Eq. (5.1) with respect to $s$ in $\mathcal{D}$, we get

$$
\begin{equation*}
T_{\psi} \frac{d s^{\star}}{d s}= \pm \frac{a \tau_{\varphi}}{\kappa_{\varphi}} N_{\varphi} \tag{5.5}
\end{equation*}
$$

Hence, using Eq. (5.4) in Eq. (5.5), we have

$$
\frac{d s^{\star}}{d s}= \pm a \tau_{\varphi}
$$

Therefore, we get

$$
\tau_{\varphi}= \pm \frac{1}{a} \frac{d s^{\star}}{d s}
$$

and so also on $I$, by Lemma 5.2. It follows that $\tau_{\varphi}$ is nowhere zero on $I$. Consequently, $\psi\left(s^{\star}\right)=\varphi(s)+\delta(s) B_{\varphi}(s)$ is of class $C^{\infty}$ on $I^{\star}$. Hence, $N=\varnothing$. Similarly, $M=\varnothing$.

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