Turkish Journal of Mathematics<br>http://journals.tubitak.gov.tr/math/<br>Research Article

Turk J Math<br>(2021) 45: 1479 - 1494<br>© TÜBİTAK<br>doi:10.3906/mat-2101-115

# Extensions of the matrix-valued $q$-Sturm-Liouville operators 

Bilender PAŞAOĞLU ALLAHVERDİEV ${ }^{1}{ }^{(D)}$, Hüseyin TUNA ${ }^{2, *}$ ©<br>${ }^{1}$ Department of Mathematics, Faculty of Arts and Sciences, Süleyman Demirel University, Isparta, Turkey<br>${ }^{2}$ Department of Mathematics, Faculty of Arts and Sciences, Mehmet Akif Ersoy University, Burdur, Turkey

| Received: 27.01 .2021 | Accepted/Published Online: 02.05.2021 | Final Version: 20.05 .2021 |
| :--- | :--- | :--- | :--- | :--- |


#### Abstract

In this paper, we investigate the matrix-valued $q$-Sturm-Liouville problems. We establish an existence and uniqueness result. Later, we introduce the corresponding maximal and minimal operators for this system. Moreover, we give a criterion under which these operators are self-adjoint. Finally, we characterize extensions (maximal dissipative, maximal accumulative, and self-adjoint) of the minimal symmetric operator.


Key words: Boundary value space, boundary condition, dissipative extensions, accretive extensions, self-adjoint extensions

## 1. Introduction

This paper deals with the extension theory of symmetric operators. This topic is one of the main research areas of operator theory. This theory was developed originally by J. Von Neumann [44]. In [22], Calkin gave the description of self-adjoint extensions of a symmetric operator in terms of abstract boundary conditions. Extensions of a symmetric operator with aid of linear relations were given by Rofe-Beketov [39]. Later, in [21, 32], the notion of a space of boundary values was introduced. The readers may find some papers that are related to extension theory in [28, 38, 47]. In [34], the authors obtained a description of extensions of a second-order symmetric operator. In [27], the author obtained a description of self-adjoint extensions of SturmLiouville operators with an operator potential. In the case when the deficiency indices take indeterminate values, a description of extensions of differential operators was given in [6, 35-37].

On the other hand, the study of matrix-valued Sturm-Liouville equations has become an important area of research because such equations arise in a variety of physical problems (for example, see [15-17, 20, 23]). Although the matrix Sturm-Liouville equations is more difficult than the scalar Sturm-Liouville equations the matrix-valued Sturm-Liouville equations have intensively been investigated during the last two decades (see $[9,13,18,19,24,45,46]$ and references therein).

Recently, $q$-difference equations have attracted tremendous interest since they have a lot of applications in sciences, e.g., quantum theory, orthogonal polynomials, hypergeometric functions (see [25] for more details). Specially, $q$-Sturm-Liouville problems were studied in [1-5, 8, 10-12, 14, 26, 31, 41-43]. The goal of our paper is to study the matrix-valued $q$-Sturm-Liouville operators. In the analysis that follows, we will largely follow a development of the theory in $[6,29,33,40,47]$.

[^0]
## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

This paper is organized as follows. In Section 2, fundamental concepts of quantum analysis are given. In Section 3, an existence and uniqueness theorem is proved for the matrix-valued $q$-Sturm-Liouville equation. Later, the corresponding maximal and minimal operators for this equation are constructed. In Section 4, a criterion under which the matrix-valued $q$-Sturm-Liouville operators are self-adjoint is given. In Section 5, maximal dissipative, maximal accumulative, and self-adjoint extensions of the minimal operators are studied.

## 2. Preliminaries

In this section, we recall some basic concepts and useful results of quantum calculus. We refer to [7, 25, 30] and some references cited therein. Let $q$ be a positive number with $0<q<1$. A set $A \subset \mathbb{R}$ is called $q$-geometric if for every $x \in A, q x \in A$. Let $y$ be a complex-valued function on $A$. Then, the $q-$ difference operator $D_{q}$ is defined by

$$
D_{q} y(x)=[y(q x)-y(x)](q x-x)^{-1} \text { for all } x \in A
$$

The $q$-derivative at zero is defined by

$$
D_{q} y(0)=\lim _{n \rightarrow \infty}\left[y\left(x q^{n}\right)-y(0)\right] x^{-1} q^{-n} \quad(x \in A)
$$

if the limit exists and does not depend on $x$ (see [7]).
The Jackson $q$-integration is given by

$$
\int_{0}^{x} f(t) d_{q} t=(1-q) x \sum_{n=0}^{\infty} f\left(q^{n} x\right) q^{n} \quad(x \in A)
$$

provided that the series converges, and

$$
\int_{a}^{b} f(t) d_{q} t=\int_{0}^{b} f(t) d_{q} t-\int_{0}^{a} f(t) d_{q} t \quad(a, b \in A)
$$

A function $f$ which is defined on $A, 0 \in A$, is said to be $q$-regular at zero if

$$
\lim _{n \rightarrow \infty} f\left(x q^{n}\right)=f(0)
$$

for every $x \in A$. Through the remainder of the paper, we deal only with $q-$ regular functions at zero.

## 3. The matrix-valued $q$-Sturm-Liouville equation

Let us consider the following matrix-valued $q$-Sturm-Liouville equation

$$
\begin{equation*}
\Upsilon(y):=-\frac{1}{q} D_{q^{-1}}\left[P(x) D_{q} y(x)\right]+Q(x) y(x)=\lambda V(x) y(x), x \in(0, a), \tag{3.1}
\end{equation*}
$$

where $P(x), V(x)$, and $Q(x)$ are $n \times n$ complex Hermitian matrix-valued functions, defined on $\left[0, q^{-1} a\right]$, continuous at zero, $q$-integrable over $[0, a], \operatorname{det} P(x) \neq 0, P^{-1}(x)$ is $q$-integrable over $[0, a], V(x)$ is positive, and $\lambda$ is a complex parameter.

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

Now, we can transform equation (3.1) into the Hamiltonian system. Let

$$
\begin{aligned}
\mathcal{Y}(x) & =\binom{y(x)}{P(x) D_{q} y(x)}, \quad \mathcal{Y}^{[q]}(x)=\binom{D_{q} y(x)}{\frac{1}{q} D_{q^{-1}}\left(P(x) D_{q} y(x)\right)}, \\
W_{1}(x) & =\left(\begin{array}{cc}
V(x) & O_{n} \\
O_{n} & O_{n}
\end{array}\right), \quad J=\left(\begin{array}{cc}
O_{n} & -I_{n} \\
I_{n} & O_{n}
\end{array}\right), \\
W_{2}(x) & =\left(\begin{array}{cc}
-Q(x) & O_{n} \\
O_{n} & P^{-1}(x)
\end{array}\right),
\end{aligned}
$$

where $I_{n}$ is a unit matrix and $O_{n}$ is a zero matrix. Then, the equation (3.1) becomes

$$
\begin{equation*}
\tau(\mathcal{Y}):=J \mathcal{Y}^{[q]}(x)-W_{2}(x) \mathcal{Y}(x)=\lambda W_{1}(x) \mathcal{Y}(x), x \in(0, a) \tag{3.2}
\end{equation*}
$$

Let $L_{q, W_{1}}^{2}[(0, a) ; E]=\left\{\mathcal{Y}: \int_{0}^{a}\left(W_{1} \mathcal{Y}, \mathcal{Y}\right)_{E} d_{q} x=\int_{0}^{a} \mathcal{Y}^{*} W_{1} \mathcal{Y} d_{q} x<\infty\right\}$ with the inner product

$$
(\mathcal{X}, \mathcal{Y}):=\int_{0}^{a}\left(W_{1} \mathcal{X}, \mathcal{Y}\right)_{E} d_{q} x=\int_{0}^{a} \mathcal{Y}^{*} W_{1} \mathcal{Y} d_{q} x
$$

where $E=\mathbb{C}^{2 n}$ is the $2 n$-dimensional Euclidean space. For any function $\mathcal{Y} \in L_{q, W_{1}}^{2}[(0, a) ; E], \mathcal{Y}(0)$ can be defined as

$$
\begin{equation*}
\mathcal{Y}(0):=\lim _{n \rightarrow \infty} \mathcal{Y}\left(q^{n}\right) \tag{3.3}
\end{equation*}
$$

Since $\mathcal{Y}$ is $q$-regular at zero, the limit in (3.3) exists and is finite.
Let $C_{q}^{2}[(0, a) ; E]=\left\{\mathcal{Y}: \mathcal{Y}\right.$ and $D_{q} \mathcal{Y}$ are $q$ - regular at zero $\}$. It is clear that $C_{q}^{2}[(0, a) ; E] \subset L_{q, W_{1}}^{2}[(0, a) ; E]$.

Theorem 3.1 For $K \in \mathbb{C}^{2 n}$, the equation (3.2) with initial condition

$$
\begin{equation*}
\mathcal{Y}(0, \lambda)=K \quad(\lambda \in \mathbb{C}) \tag{3.4}
\end{equation*}
$$

has a unique solution in $C_{q}^{2}[(0, a) ; E]$.
Proof An integration yields

$$
\begin{equation*}
\mathcal{Y}(x, \lambda)=K-q \int_{0}^{x} J\left[\lambda W_{1}(q t, \lambda)+W_{2}(q t, \lambda)\right] \mathcal{Y}(q t, \lambda) d_{q} t \tag{3.5}
\end{equation*}
$$

where $x \in(0, a)$. Conversely, every solution of equation (3.5) is also a solution of the equation (3.2).
Let us define the sequence $\left\{\mathcal{Y}_{n}\right\}_{n \in \mathbb{N}}$ of successive approximations by

$$
\begin{gather*}
\mathcal{Y}_{0}(x, \lambda)=K \\
\mathcal{Y}_{n+1}(x, \lambda)= \\
K-q \int_{0}^{x} J\left[\lambda W_{1}(q t, \lambda)+W_{2}(q t, \lambda)\right] \mathcal{Y}_{n}(q t, \lambda) d_{q} t \tag{3.6}
\end{gather*}
$$

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

where $n=0,1,2, \ldots$, and $x \in(0, a)$. Then, we shall prove that the sequence $\left\{\mathcal{Y}_{n}\right\}_{n \in \mathbb{N}}$ converges to a function $y$ uniformly on each compact subset of $(0, a)$. There exist positive numbers $\vartheta(\lambda)$ and $\varsigma(\lambda)$ such that

$$
\begin{gathered}
\left\|J\left[\lambda W_{1}(x, \lambda)+W_{2}(x, \lambda)\right]\right\| \leq \vartheta(\lambda) \\
\left\|\mathcal{Y}_{1}(x, \lambda)\right\| \leq \varsigma(\lambda), x \in(0, a)
\end{gathered}
$$

Using mathematical induction, we get

$$
\left\|\mathcal{Y}_{n+1}(x, \lambda)-\mathcal{Y}_{n}(x, \lambda)\right\| \leq \vartheta(\lambda) q^{\frac{n(n+1)}{2}} \frac{(\varsigma(\lambda) x(1-q))^{n}}{(q ; q)_{n}} \quad(n \in \mathbb{N})
$$

An application of the Weierstrass $M$-test implies that the sequence $\left\{\mathcal{Y}_{n}\right\}_{n \in \mathbb{N}}$ converges to a function $y$ uniformly on each compact subset of $(0, a)$. One can prove that $\mathcal{Y}$ and $D_{q} \mathcal{Y}$ are continuous on $[0, a]$. It is clear that the function $\mathcal{Y}$ satisfies the condition (3.4).

Now, we show that the equation (3.2) has a unique solution, assume $\mathcal{Z}$ is another one. Then $\mathcal{Z}$ is continuous. Therefore, there exists a positive number $\mathcal{M}$ such that $\|\mathcal{Y}-\mathcal{Z}\| \leq \mathcal{M}$. Proceeding as above we conclude that

$$
\|\mathcal{Y}(x, \lambda)-\mathcal{Z}(x, \lambda)\| \leq \mathcal{M} \vartheta(\lambda) q^{\frac{n(n+1)}{2}} \frac{(x(1-q))^{n}}{(q ; q)_{n}} \quad(n \in \mathbb{N})
$$

Since

$$
\lim _{n \rightarrow \infty} \mathcal{M} \vartheta(\lambda) q^{\frac{n(n+1)}{2}} \frac{(x(1-q))^{n}}{(q ; q)_{n}}=0
$$

we arrive at $\mathcal{Y}=\mathcal{Z}$ on $[0, a]$.
Now, we will give the definition of maximal and minimal operators for the matrix-valued $q$-SturmLiouville equations.

Denote

$$
\begin{gather*}
\mathcal{D}_{\max }:= \\
\left\{\begin{array}{c}
\mathcal{Y} \in L_{q, W_{1}}^{2}[(0, a) ; E]: J \mathcal{Y}^{[q]}(x)-W_{2}(x) \mathcal{Y}(x)=W_{1} F \text { exists in }(0, a) \\
\text { and } F \in L_{q, W_{1}}^{2}[(0, a) ; E]
\end{array}\right\}, \\
\mathcal{D}_{\min }:= \\
\left\{\begin{array}{c}
\mathcal{Y} \in L_{q, W_{1}}^{2}[(0, a) ; E]: \mathcal{Y} \text { and } P D_{q} \mathcal{Y} \text { are } q \text { - regular at zero, } \\
J \mathcal{Y}[q](x)-W_{2}(x) \mathcal{Y}(x)=W_{1} F \text { exists in }(0, a), \\
F \in L_{q, W_{1}}^{2}[(0, a) ; E], \\
\text { and } \widehat{\mathcal{Y}}(0)=\widehat{\mathcal{Y}}(a)=0 .
\end{array}\right\}, \tag{3.7}
\end{gather*}
$$

where $\widehat{\mathcal{Y}}(x)=\binom{y(x)}{P(x) D_{q^{-1}} y(x)}$.
The operator $T_{\min }$ defined by

$$
\begin{aligned}
T_{\min } & : \mathcal{D}_{\min } \rightarrow L_{q, W_{1}}^{2}[(0, a) ; E] \\
\mathcal{Y} & \rightarrow T_{\min } \mathcal{Y}=\tau(\mathcal{Y})
\end{aligned}
$$

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

is called the minimal operator generated by the matrix-valued $q$-Sturm-Liouville equation. Similarly, the operator $T_{\max }$ defined by

$$
\begin{aligned}
T_{\max } & : \mathcal{D}_{\max } \rightarrow L_{q, W_{1}}^{2}[(0, a) ; E] \\
\mathcal{Y} & \rightarrow T_{\max } \mathcal{Y}=\tau(\mathcal{Y})
\end{aligned}
$$

is called the maximal operator for the matrix-valued $q$-Sturm-Liouville equation.
Now, we establish the following Green's formula.
Theorem 3.2 Let $\mathcal{Y}, \mathcal{Z} \in \mathcal{D}_{\max }$. Then we have

$$
\begin{gathered}
\int_{0}^{a} \mathcal{Z}^{*}(x) J \mathcal{Y}^{[q]}(x) d_{q} x-\int_{0}^{a}\left\{J \mathcal{Z}^{[q]}(x)\right\}^{*} \mathcal{Y}(x) d_{q} x \\
=\widehat{\mathcal{Z}}^{*}(a) J \widehat{\mathcal{Y}}(a)-\widehat{\mathcal{Z}}^{*}(0) J \widehat{\mathcal{Y}}(0)
\end{gathered}
$$

where $x \in(0, a)$.

## Proof

$$
\begin{gathered}
\int_{0}^{a} \mathcal{Z}^{*}(x) J \mathcal{Y}^{[q]}(x) d_{q} x-\int_{0}^{a}\left\{J \mathcal{Z}^{[q]}(x)\right\}^{*} \mathcal{Y}(x) d_{q} x \\
=\int_{0}^{a}\binom{z(x)}{P(x) D_{q} z(x)}^{*}\left(\begin{array}{cc}
O_{n} & -I_{n} \\
I_{n} & O_{n}
\end{array}\right)\binom{D_{q} y(x)}{\frac{1}{q} D_{q^{-1}}\left(P(x) D_{q} y(x)\right)} d_{q} x \\
-\int_{0}^{a}\left\{\binom{D_{q} z(x)}{\frac{1}{q} D_{q^{-1}}\left(P(x) D_{q} z(x)\right)}\right\}^{*}\left(\begin{array}{cc}
O_{n} & -I_{n} \\
I_{n} & O_{n}
\end{array}\right)\left(\begin{array}{c} 
\\
P(x) D_{q} y(x)
\end{array}\right) d_{q} x \\
=\int_{0}^{a}\left[z^{*}(x)\left\{\frac{1}{q} D_{q^{-1}}\left(P(x) D_{q} y(x)\right)\right\}+\left(P(x) D_{q} z(x)\right)^{*} D_{q} y(x)\right] d_{q} x \\
-\int_{0}^{a}\left[\left\{-\frac{1}{q} D_{q^{-1}}\left(P(x) D_{q} z(x)\right)\right\}^{*} y(x)+\left(D_{q} z(x)\right)^{*} P(x) D_{q} y(x)\right] d_{q} x \\
\quad=\int_{0}^{a} D_{q}\left\{\left[\left(P D_{q} z\right)\left(q^{-1} x\right)\right]^{*} y(x)-z^{*}(x)\left(P D_{q} y\right)\left(q^{-1} x\right)\right\} d_{q} x \\
\quad=\widehat{\mathcal{Z}}^{*}(a) J \widehat{\mathcal{Y}}(a)-\widehat{\mathcal{Z}}^{*}(0) J \widehat{\mathcal{Y}}(0)
\end{gathered}
$$

## Theorem 3.3 (Green's formula)

$$
\begin{equation*}
\left(T_{\max } \mathcal{Y}, \mathcal{Z}\right)-\left(\mathcal{Y}, T_{\max } \mathcal{Z}\right)=[\mathcal{Y}, \mathcal{Z}]_{a}-[\mathcal{Y}, \mathcal{Z}]_{0} \tag{3.8}
\end{equation*}
$$

where $[\mathcal{Y}, \mathcal{Z}]_{x}:=\widehat{\mathcal{Z}}^{*}(a) J \widehat{\mathcal{Y}}(x), x \in[0, a]$.

Lemma 3.4 The operator $T_{\min }$ is Hermitian.
Proof For $\mathcal{Y}, \mathcal{Z} \in \mathcal{D}_{\text {min }}$, there exist $F, G \in L_{q, W_{1}}^{2}[(0, a) ; E]$ such that $\tau(\mathcal{Y})=W_{1} F$ and $\tau(\mathcal{Z})=W_{1} G$. From (3.7) and (3.8), we conclude that

$$
\begin{aligned}
& \left(T_{\min } \mathcal{Y}, \mathcal{Z}\right)-\left(\mathcal{Y}, T_{\min } \mathcal{Z}\right)=(F, \mathcal{Z})-(\mathcal{Y}, G) \\
& =\int_{0}^{a}\left[\mathcal{Z}^{*}(t) W_{1} F-G^{*}(t) W_{1} \mathcal{Y}(t)\right] d_{q} t \\
& =\int_{0}^{a}\left[\mathcal{Z}^{*}(t) \tau(\mathcal{Y})-\tau^{*}(\mathcal{Z}) \mathcal{Y}(t)\right] d_{q} t \\
& =[\mathcal{Y}, \mathcal{Z}]_{a}-[\mathcal{Y}, \mathcal{Z}]_{0}=0
\end{aligned}
$$

The following lemma has a similar proof to that of Lemma 3.4.

Lemma 3.5 For all $\mathcal{Y} \in \mathcal{D}_{\text {min }}$ and for all $\mathcal{Z} \in \mathcal{D}_{\max }$, we have

$$
\left(T_{\min } \mathcal{Y}, \mathcal{Z}\right)=\left(\mathcal{Y}, T_{\max } \mathcal{Z}\right)
$$

Lemma 3.6 Let the null space and the range of an operator $T$ be denoted by $\mathcal{N}(T)$ and $\mathcal{R}(T)$, respectively. Then we have

$$
\mathcal{R}\left(T_{\min }\right)=\mathcal{N}\left(T_{\max }\right)^{\perp}
$$

where the superscript $\perp$ denotes the orthogonal complement of a subspace.
Proof Given any $\xi \in \mathcal{R}\left(T_{\min }\right)$, there exists $\mathcal{Y} \in \mathcal{D}_{\min }$ such that $T_{\min } \mathcal{Y}=\xi$. It follows from Lemma 3.5 that

$$
(\xi, \mathcal{Z})=\left(T_{\min } \mathcal{Y}, \mathcal{Z}\right)=\left(\mathcal{Y}, T_{\max } \mathcal{Z}\right)=0
$$

for each $\mathcal{Z} \in \mathcal{N}\left(T_{\max }\right)$.
Now we prove that $\mathcal{N}\left(T_{\max }\right)^{\perp} \subset \mathcal{R}\left(T_{\min }\right)$. For any given $\xi \in \mathcal{N}\left(T_{\max }\right)^{\perp}$ and for all $\mathcal{Z} \in \mathcal{N}\left(T_{\max }\right)$, we have $(\xi, \mathcal{Z})=0$. Let us consider the following problem:

$$
\begin{gather*}
J \mathcal{Y}^{[q]}(x)-W_{2}(x) \mathcal{Y}(x)=W_{1}(x) \xi(x), x \in(0, a) \\
\mathcal{Y}(0, \lambda)=P(0) D_{q^{-1}} \mathcal{Y}(0, \lambda)=0 \tag{3.9}
\end{gather*}
$$

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

It follows from Theorem 3.1 that the problem (3.9) has a unique solution on $(0, a)$. Let $\Psi(x)=\left(\psi_{1}, \psi_{2}, \ldots, \psi_{n}\right)$ be the fundamental solution of the system

$$
J \mathcal{Y}^{[q]}(x)-W_{2}(x) \mathcal{Y}(x)=0, \Psi(a)=J, x \in(0, a)
$$

It is clear that $\psi_{i} \in \mathcal{N}\left(T_{\max }\right)$ for $1 \leq i \leq n$. By Theorem 3.3, for $1 \leq i \leq n$, we get

$$
\begin{aligned}
0 & =\left(\xi, \psi_{i}\right)=\int_{0}^{a} \psi_{i}^{*}(t) W_{1}(x) \xi(t) d_{q} t=\int_{0}^{a} \psi_{i}^{*}(t) \tau(\mathcal{Y})(t) d_{q} t \\
& =\int_{0}^{a} \psi_{i}^{*}(t) \tau(\mathcal{Y})(t) d_{q} t-\int_{0}^{a} \tau\left(\psi_{i}\right)^{*}(t) \mathcal{Y}(t) d_{q} t \\
& =\left[\mathcal{Y}, \psi_{i}\right]_{a}-\left[\mathcal{Y}, \psi_{i}\right]_{0}=\left[\mathcal{Y}, \psi_{i}\right]_{a}
\end{aligned}
$$

Thus, we have $\left[\mathcal{Y}, \psi_{i}\right]_{a}=\widehat{\Psi}^{*}(a) J \widehat{\mathcal{Y}}(a)=\widehat{\mathcal{Y}}(a)=0$, i.e., $\xi \in \mathcal{R}\left(T_{\text {min }}\right)$.

Theorem 3.7 The operator $T_{\min }$ is a symmetric operator and the operator $T_{\max }$ is a densely defined operator. Furthermore, $T_{\min }^{*}=T_{\max }$, where $T_{\min }^{*}$ denotes the adjoint operator of $T_{\min }$.

Proof Firstly, we prove that $\mathcal{D}_{\text {min }}^{\perp}=\{0\}$. Assume that $\xi \in \mathcal{D}_{\min }^{\perp}$. Then, for all $\mathcal{Z} \in \mathcal{D}_{\text {min }}$, we have $(\xi, \mathcal{Z})=0$. Set $T_{\min } \mathcal{Z}(x)=\varphi(x)$. Let $y($.$) be any solution of the system$

$$
J \mathcal{Y}^{[q]}(x)-W_{2}(x) \mathcal{Y}(x)=W_{1}(x) \xi(x), x \in(0, a)
$$

It follows from Theorem 3.3 that

$$
\begin{aligned}
(\mathcal{Y}, \varphi)-(\xi, \mathcal{Z}) & =\int_{0}^{a} \varphi^{*}(t) W_{1}(t) \mathcal{Y}(t) d_{q} t-\int_{0}^{a} \mathcal{Z}^{*}(t) W_{1}(t) \xi(t) d_{q} t \\
& =\int_{0}^{a} \tau(\mathcal{Z})^{*}(t) \mathcal{Y}(t) d_{q} t-\int_{0}^{a} \mathcal{Z}^{*}(t) \tau(\mathcal{Y})(t) d_{q} t \\
& =-[\mathcal{Y}, \mathcal{Z}]_{a}+[\mathcal{Y}, \mathcal{Z}]_{0}=0
\end{aligned}
$$

i.e. $(\mathcal{Y}, \varphi)=(\xi, \mathcal{Z})=0$. From Lemma 3.6, we see that $\mathcal{Y} \in \mathcal{R}\left(T_{\text {min }}\right)=\mathcal{N}\left(T_{\max }\right)^{\perp}$. Thus, $\xi=0$.

We will denote by $\mathcal{D}_{\min }^{*}$ the domain of the operator $T_{\min }^{*}$. Now, we prove that $\mathcal{D}_{\min }^{*}=\mathcal{D}_{\text {max }}$, and $T_{\min }^{*} \mathcal{Y}=T_{\max } \mathcal{Y}$ for all $\mathcal{Y} \in \mathcal{D}_{\min }^{*}$. It follows from Lemma 3.5 that, for any given $\mathcal{Y} \in \mathcal{D}_{\max }$,

$$
\left(\mathcal{Y}, T_{\min } \mathcal{Z}\right)=\left(T_{\max } \mathcal{Y}, \mathcal{Z}\right) \text { for all } \mathcal{Z} \in \mathcal{D}_{\min }
$$

Consequently, the functional $\left(\mathcal{Y}, T_{\min }().\right)$ is continuous on $\mathcal{D}_{\min }$ and $\mathcal{Y} \in \mathcal{D}_{\min }^{*}$, i.e. $\mathcal{D}_{\max } \subset \mathcal{D}_{\min }^{*}$.
We prove the reverse conclusion, i.e. $\mathcal{D}_{\min }^{*} \subset \mathcal{D}_{\max }$. If $\mathcal{Y} \in \mathcal{D}_{\min }^{*}$, then $\mathcal{Y}, \varphi \in L_{q, W_{1}}^{2}[(0, a) ; E]$, where $\varphi:=T_{\min }^{*} \mathcal{Y}$. Assume that $\mathcal{U}$ is a solution of the equation

$$
\begin{equation*}
J \mathcal{U}^{[q]}(x)-W_{2}(x) \mathcal{U}(x)=W_{1}(x) \varphi(x) . \tag{3.10}
\end{equation*}
$$

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

By Lemma 3.5, we see that

$$
(\varphi, \mathcal{Z})=\left(T_{\max } \mathcal{U}, \mathcal{Z}\right)=\left(\mathcal{U}, T_{\min } \mathcal{Z}\right)
$$

Thus, we get

$$
\begin{aligned}
\left(\mathcal{Y}-\mathcal{U}, T_{\min } \mathcal{Z}\right) & =\left(\mathcal{Y}, T_{\min } \mathcal{Z}\right)-\left(\mathcal{U}, T_{\min } \mathcal{Z}\right) \\
& =\left(T_{\min }^{*} \mathcal{Y}, \mathcal{Z}\right)-(\varphi, \mathcal{Z})=0
\end{aligned}
$$

i.e. $\mathcal{Y}-\mathcal{U} \in \mathcal{R}\left(T_{\min }\right)^{\perp}$. It follows from Lemma 3.6 that $\mathcal{Y}-\mathcal{U} \in \mathcal{N}\left(T_{\max }\right)$.

Using (3.10), we get

$$
\begin{aligned}
& J \mathcal{Y}^{[q]}(x)-W_{2}(x) \mathcal{Y}(x) \\
& =J \mathcal{U}^{[q]}(x)-W_{2}(x) \mathcal{U}(x)=W_{1}(x) \varphi(x), x \in(0, a)
\end{aligned}
$$

Since $\mathcal{Y}, \varphi \in L_{q, W_{1}}^{2}[(0, a) ; E]$, we have that $\mathcal{Y} \in \mathcal{D}_{\max }$ and

$$
T_{\max } \mathcal{Y}=\varphi=T_{\min }^{*} \mathcal{Y}
$$

This completes the proof.

## 4. A criterion of the self-adjoint matrix-valued $q$-Sturm-Liouville operators

In this section, we give a criterion under which the matrix-valued $q$-Sturm-Liouville operators are self-adjoint.
Let $\Sigma$ and $\Lambda$ be $m \times 2 n$ matrices such that $\operatorname{rank}(\Sigma: \Lambda)=m$. Then, we define the operator $T$ by

$$
\begin{equation*}
T: \mathcal{D} \rightarrow L_{q, W_{1}}^{2}[(0, a) ; E] \tag{4.1}
\end{equation*}
$$

where

$$
\mathcal{D}:=\left\{\begin{array}{c}
\mathcal{Y} \in L_{q, V}^{2}(0, a): \mathcal{Y} \text { and } P D_{q} \mathcal{Y} \text { are } q \text { - regular at zero }  \tag{4.2}\\
J \mathcal{Y}^{[q]}(x)-W_{2}(x) \mathcal{Y}(x)=W_{1}(x) F(x) \text { exists in }(0, a) \\
F \in L_{q, V}^{2}(0, a), \text { and } \\
\Sigma \widehat{\mathcal{Y}}(0)+\Lambda \widehat{\mathcal{Y}}(a)=0
\end{array}\right\}
$$

Let $\Omega$ and $\Gamma$ be $(4 n-m) \times 2 n$ matrices, chosen so that $\operatorname{rank}(\Omega: \Gamma)=4 n-m$ and $\left(\begin{array}{ll}\Sigma & \Lambda \\ \Omega & \Gamma\end{array}\right)$ is nonsingular. Let $\left(\begin{array}{cc}\widetilde{\Sigma} & \widetilde{\Lambda} \\ \widetilde{\Omega} & \widetilde{\Gamma}\end{array}\right)$ be chosen so that

$$
\left(\begin{array}{cc}
\widetilde{\Sigma} & \widetilde{\Lambda}  \tag{4.3}\\
\widetilde{\Omega} & \widetilde{\Gamma}
\end{array}\right)^{*}\left(\begin{array}{ll}
\Sigma & \Lambda \\
\Omega & \Gamma
\end{array}\right)=\left(\begin{array}{cc}
-J & 0 \\
0 & J
\end{array}\right)
$$

Theorem 4.1 For $\mathcal{Y}, \mathcal{Z} \in \mathcal{D}_{\max }$, we have

$$
\begin{aligned}
\left(T_{\max } \mathcal{Y}, \mathcal{Z}\right)-\left(\mathcal{Y}, T_{\max } \mathcal{Z}\right) & =[\widetilde{\Sigma} \widehat{\mathcal{Z}}(0)+\widetilde{\Lambda} \widehat{\mathcal{Z}}(a)]^{*}[\Sigma \widehat{\mathcal{Y}}(0)+\Lambda \widehat{\mathcal{Y}}(a)] \\
& +[\widetilde{\Omega} \widehat{\mathcal{Z}}(0)+\widetilde{\Gamma} \widehat{\mathcal{Z}}(a)]^{*}[\Omega \widehat{\mathcal{Y}}(0)+\Gamma \widehat{\mathcal{Y}}(a)]
\end{aligned}
$$

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

Proof By virtue of (3.8) and (4.3), we get

$$
\begin{aligned}
& \left(T_{\max } \mathcal{Y}, \mathcal{Z}\right)-\left(\mathcal{Y}, T_{\max } \mathcal{Z}\right)=[\mathcal{Y}, \mathcal{Z}]_{a}-[\mathcal{Y}, \mathcal{Z}]_{0} \\
& =\left(\begin{array}{cc}
\widehat{\mathcal{Z}}^{*}(0) & \widehat{\mathcal{Z}}^{*}(a)
\end{array}\right)\left(\begin{array}{cc}
-J & 0 \\
0 & J
\end{array}\right)\binom{\widehat{\mathcal{Y}}(0)}{\widehat{\mathcal{Y}}(a)} \\
& =\left(\begin{array}{cc}
\widehat{\mathcal{Z}}^{*}(0) & \widehat{\mathcal{Z}}^{*}(a)
\end{array}\right)\left(\begin{array}{cc}
\widetilde{\Sigma} & \widetilde{\Lambda} \\
\widetilde{\Omega} & \widetilde{\Gamma}
\end{array}\right)^{*}\left(\begin{array}{ll}
\Sigma & \Lambda \\
\Omega & \Gamma
\end{array}\right)\binom{\widehat{\mathcal{Y}}(0)}{\widehat{\mathcal{Y}}(a)} \\
& =\left[\left(\begin{array}{cc}
\widetilde{\Sigma} & \widetilde{\Lambda} \\
\widetilde{\Omega} & \widetilde{\Gamma}
\end{array}\right)\binom{\widehat{\mathcal{Z}}(0)}{\widehat{\mathcal{Z}}(a)}\right]^{*}\left[\left(\begin{array}{cc}
\Sigma & \Lambda \\
\Omega & \Gamma
\end{array}\right)\binom{\widehat{\mathcal{Y}}(0)}{\widehat{\mathcal{Y}}(a)}\right] \\
& =\binom{\widetilde{\Sigma} \widehat{\mathcal{Z}}(0)+\widetilde{\Lambda} \widehat{\mathcal{Z}}(a)}{\widetilde{\Omega} \widehat{\mathcal{Z}}(0)+\widetilde{\Gamma} \widehat{\mathcal{Z}}(a)}^{*}\binom{\Sigma \widehat{\mathcal{Y}}(0)+\Lambda \widehat{\mathcal{Y}}(a)}{\Omega \widehat{\mathcal{Y}}(0)+\Gamma \widehat{\mathcal{Y}}(a)} .
\end{aligned}
$$

Now, we describe the adjoint of the operator $T$. Denote

$$
\begin{gathered}
\mathcal{D}^{*}:= \\
\left\{\begin{array}{c}
\mathcal{Z} \in L_{q, W_{1}}^{2}[(0, a) ; E]: J \mathcal{Z}^{[q]}-W_{2}(x) \mathcal{Z}(x)=W_{1}(x) F_{1}(x) \text { exists in }(0, a) \\
F_{1} \in L_{q, W_{1}}^{2}[(0, a) ; E] \text { and } \widetilde{\Omega} \widehat{\mathcal{Z}}(0)+\widetilde{\Gamma} \widehat{\mathcal{Z}}(a)=0
\end{array}\right\},
\end{gathered}
$$

Theorem 4.2 For $\mathcal{Z} \in \mathcal{D}^{*}, T^{*} \mathcal{Z}=\widehat{F_{1}}$ if and only if

$$
J \mathcal{Z}^{[q]}-W_{2}(x) \mathcal{Z}(x)=W_{1}(x) F_{1}(x)
$$

Proof Since $T_{\min } \subset T \subset T_{\max }$, we have $T_{\min } \subset T^{*} \subset T_{\max }$. Let $\mathcal{Y} \in \mathcal{D}$ and $\mathcal{Z} \in \mathcal{D}^{*}$. It follows from Theorem 4.1 that

$$
\begin{aligned}
(T \mathcal{Y}, \mathcal{Z})-\left(\mathcal{Y}, T^{*} \mathcal{Z}\right) & =[\widetilde{\Sigma} \widehat{\mathcal{Z}}(0)+\widetilde{\Lambda} \widehat{\mathcal{Z}}(a)]^{*}[\Sigma \widehat{\mathcal{Y}}(0)+\Lambda \widehat{\mathcal{Y}}(a)] \\
& +[\widetilde{\Omega} \widehat{\mathcal{Z}}(0)+\widetilde{\Gamma} \widehat{\mathcal{Z}}(a)]^{*}[\Omega \widehat{\mathcal{Y}}(0)+\Gamma \widehat{\mathcal{Y}}(a)]
\end{aligned}
$$

Then we get

$$
0=[\widetilde{\Omega} \widehat{\mathcal{Z}}(0)+\widetilde{\Gamma} \widehat{\mathcal{Z}}(a)]^{*}[\Omega \widehat{\mathcal{Y}}(0)+\Gamma \widehat{\mathcal{Y}}(a)]
$$

Since $\Omega \widehat{\mathcal{Y}}(0)+\Gamma \widehat{\mathcal{Y}}(a)$ is arbitrary, we have $\widetilde{\Omega} \widehat{\mathcal{Z}}(0)+\widetilde{\Gamma} \widehat{\mathcal{Z}}(a)=0$.
Conversely, if $\mathcal{Z}$ satisfies the criteria listed above then $\mathcal{Z} \in \mathcal{D}^{*}$.

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

Now, we find parametric boundary conditions for $\mathcal{D}$ and $\mathcal{D}^{*}$. Recall that

$$
\begin{equation*}
\Sigma \widehat{\mathcal{Y}}(0)+\Lambda \widehat{\mathcal{Y}}(a)=0, \Omega \widehat{\mathcal{Y}}(0)+\Gamma \widehat{\mathcal{Y}}(a)=F_{2} \tag{4.4}
\end{equation*}
$$

where $F_{2}$ is arbitrary. Hence, we get

$$
\left(\begin{array}{cc}
\Sigma & \Lambda  \tag{4.5}\\
\Omega & \Gamma
\end{array}\right)\binom{\widehat{\mathcal{Y}}(0)}{\widehat{\mathcal{Y}}(a)}=\binom{0}{F_{2}}
$$

If we multiply both sides of (4.5) by

$$
\left(\begin{array}{cc}
-J & 0 \\
0 & J
\end{array}\right)\left(\begin{array}{ll}
\widetilde{\Sigma} & \widetilde{\Lambda} \\
\widetilde{\Omega} & \widetilde{\Gamma}
\end{array}\right)^{*}
$$

we conclude that

$$
\begin{equation*}
\binom{\widehat{\mathcal{Y}}(0)}{\widehat{\mathcal{Y}}(a)}=\binom{J \widetilde{\Omega}^{*} F_{2}}{-J \widetilde{\Gamma}^{*} F_{2}} . \tag{4.6}
\end{equation*}
$$

Similarly, we can find parametric boundary conditions for $\mathcal{D}^{*}$. Since

$$
\widetilde{\Sigma} \widehat{\mathcal{Z}}(0)+\widetilde{\Lambda} \widehat{\mathcal{Z}}(a)=F_{3}, \widetilde{\Omega} \widehat{\mathcal{Z}}(0)+\widetilde{\Gamma} \widehat{\mathcal{Z}}(a)=0
$$

where $F_{3}$ is arbitrary, we obtain

$$
\left(\begin{array}{cc}
\widehat{\mathcal{Z}}^{*}(0) & \widehat{\mathcal{Z}}^{*}(a)
\end{array}\right)\left(\begin{array}{cc}
\widetilde{\Sigma} & \widetilde{\Lambda}  \tag{4.7}\\
\widetilde{\Omega} & \widetilde{\Gamma}
\end{array}\right)^{*}=\left(\begin{array}{ll}
F_{3}^{*} & 0
\end{array}\right) .
$$

Multiplying both sides of (4.7) by

$$
\left(\begin{array}{ll}
\Sigma & \Lambda \\
\Omega & \Gamma
\end{array}\right)\left(\begin{array}{cc}
-J & 0 \\
0 & J
\end{array}\right)
$$

it follows that

$$
\begin{equation*}
\widehat{\mathcal{Z}}(0)=-J \Sigma^{*} F_{3}, \quad \widehat{\mathcal{Z}}(a)=J \Lambda^{*} F_{3} . \tag{4.8}
\end{equation*}
$$

Now, we have the following theorem.

Theorem 4.3 The operator $T$ is self-adjoint if and only if $m=2 n$ and $\Sigma J \Sigma^{*}=\Lambda J \Lambda^{*}$.
Proof Let $T$ be a self-adjoint operator. Then $\mathcal{Z}$ satisfies the boundary conditions for $\mathcal{D}$, that is

$$
\Sigma \widehat{\mathcal{Z}}(0)+\Lambda \widehat{\mathcal{Z}}(a)=0
$$

It follows from (4.8) that

$$
\begin{aligned}
\Sigma\left(-J \Sigma^{*} F_{3}\right)+\Lambda\left(J \Lambda^{*} F_{3}\right) & =0 \\
{\left[\Sigma J \Sigma^{*}-\Lambda J \Lambda^{*}\right] F_{3} } & =0
\end{aligned}
$$

Since $F_{3}$ is arbitrary, we see that

$$
\Sigma J \Sigma^{*}=\Lambda J \Lambda^{*}
$$

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

Conversely, if $\Sigma J \Sigma^{*}=\Lambda J \Lambda^{*}$, then we get

$$
\left(\begin{array}{cc}
-\Sigma J & \Lambda J
\end{array}\right)\binom{\Sigma^{*}}{\Lambda^{*}}=0
$$

i.e. the columns of $\binom{\Sigma^{*}}{\Lambda^{*}}$ for $2 n$ independent solutions to the equation

$$
\left(\begin{array}{cc}
-\Sigma J & \Lambda J
\end{array}\right) X=0
$$

By virtue of (4.4) and (4.6), we deduce that

$$
\left(\begin{array}{cc}
-\Sigma J & \Lambda J
\end{array}\right)\binom{\widetilde{\Omega}^{*}}{\widetilde{\Gamma}^{*}}=0
$$

Thus, there must be a constant, nonsingular matrix $K$ such that

$$
\binom{\widetilde{\Omega}^{*}}{\widetilde{\Gamma}^{*}} K^{*}=\binom{\Sigma^{*}}{\Lambda^{*}}
$$

or

$$
\left(\begin{array}{cc}
\Sigma & \Lambda
\end{array}\right)=K\left(\begin{array}{ll}
\widetilde{\Omega} & \widetilde{\Gamma}
\end{array}\right)
$$

Clearly, the conditions $\Sigma \widehat{\mathcal{Y}}(0)+\Lambda \widehat{\mathcal{Y}}(a)=0$ and $\Omega \widehat{\mathcal{Y}}(0)+\Gamma \widehat{\mathcal{Y}}(a)=0$ are equivalent. Since the forms of $T$ and $T^{*}$ are the same, this gives $T=T^{*}$.

## 5. Extensions of the matrix-valued $q$-Sturm-Liouville operators

In this section, we shall describe all the self-adjoint, dissipative, and accumulative extensions of the corresponding minimal operator $T_{\text {min }}$.

We begin this section with a definition (see [21, 29, 32]).

Definition 5.1 Let $\mathbb{H}$ be a Hilbert space; let $\Pi_{1}$ and $\Pi_{2}$ be linear mappings of $\mathcal{D}\left(\mathcal{B}^{*}\right)$ into $\mathbb{H}$, where $\mathcal{B}$ is a closed symmetric operator acting in a Hilbert space $\mathcal{H}$ with equal (finite or infinite) deficiency indices. Then the triplet $\left(\mathbb{H}, \Pi_{1}, \Pi_{2}\right)$ is called a space of boundary values of the operator $\mathcal{B}$ if
(i) $\left(\mathcal{B}^{*} h, g\right)_{\mathcal{H}}-\left(h, \mathcal{B}^{*} g\right)_{\mathcal{H}}=\left(\Pi_{1} h, \Pi_{2} g\right)_{\mathbb{H}}-\left(\Pi_{2} h, \Pi_{1} g\right)_{\mathbb{H}}, \forall h, g \in \mathcal{D}\left(\left(\mathcal{B}^{*}\right)\right.$, and
(ii) for every $G_{1}, G_{2} \in \mathbb{H}$, there exists a vector $g \in \mathcal{D}\left(\mathcal{B}^{*}\right)$ such that $\Pi_{1} g=G_{1}$ and $\Pi_{2} g=G_{2}$.

In the next results, we use the following notation:
$\Pi_{1}, \Pi_{2}: \mathcal{D}_{\max } \rightarrow E \oplus E$, where $E=\mathbb{C}^{n}$,

$$
\begin{equation*}
\Pi_{1} \mathcal{Y}=\binom{-\mathcal{Y}(0)}{\mathcal{Y}(a)}, \Pi_{2} \mathcal{Y}=\binom{\left(P D_{q^{-1}} \mathcal{Y}\right)(0)}{\left(P D_{q^{-1}} \mathcal{Y}\right)(a)} \tag{5.1}
\end{equation*}
$$

and $y \in \mathcal{D}_{\text {max }}$.

Theorem 5.2 The triplet $\left(E \oplus E, \Pi_{1}, \Pi_{2}\right)$ defined by (5.1) is a space of boundary values of the symmetric operator $T_{\text {min }}$.

Proof From (5.1) and (3.8), we conclude that

$$
\begin{gathered}
\left(\Pi_{1} \mathcal{Y}, \Pi_{2} \mathcal{Z}\right)_{E \oplus E}-\left(\Pi_{2} \mathcal{Y}, \Pi_{1} \mathcal{Z}\right)_{E \oplus E}=-\left(\mathcal{Y}(0),\left(P D_{q^{-1}} \mathcal{Z}\right)(0)\right)_{E} \\
+\left(\mathcal{Y}(a),\left(P D_{q^{-1}} \mathcal{Z}\right)(a)\right)_{E}+\left(\left(P D_{q^{-1}} \mathcal{Y}\right)(0), \mathcal{Z}(0)\right)_{E} \\
-\left(\left(P D_{q^{-1}} \mathcal{Y}\right)(a), \mathcal{Z}(a)\right)_{E}=[\mathcal{Y}, \mathcal{Z}](a)-[\mathcal{Y}, \mathcal{Z}](0) \\
=\left(T_{\max } \mathcal{Y}, \mathcal{Z}\right)-\left(\mathcal{Y}, T_{\max } \mathcal{Z}\right)
\end{gathered}
$$

where $\mathcal{Y}, \mathcal{Z} \in \mathcal{D}_{\text {max }}$.
Now, we show the second assumption of the definition of space of boundary values.
Let $\Lambda=\binom{\Lambda_{1}}{\Lambda_{2}}, \Gamma=\binom{\Gamma_{1}}{\Gamma_{2}} \in E \oplus E$. Then the vector-valued function

$$
\mathcal{Y}(t)=\alpha_{1}(t) o \Lambda_{1}+\alpha_{2}(t) o \Gamma_{1}+\beta_{1}(t) o \Lambda_{2}+\beta_{2}(t) o \Gamma_{2}
$$

where o is a symbol of the Hadamard product of vectors and

$$
\begin{aligned}
& \alpha_{1}(t)=\left(\begin{array}{c}
\alpha_{11}(t) \\
\ldots \\
\alpha_{1 n}(t)
\end{array}\right), \alpha_{2}(t)=\left(\begin{array}{c}
\alpha_{21}(t) \\
\ldots \\
\alpha_{2 n}(t)
\end{array}\right) \in E, \\
& \beta_{1}(t)=\left(\begin{array}{c}
\beta_{11}(t) \\
\ldots \\
\beta_{1 n}(t)
\end{array}\right), \beta_{2}(t)=\left(\begin{array}{c}
\beta_{21}(t) \\
\ldots \\
\beta_{2 n}(t)
\end{array}\right) \in E,
\end{aligned}
$$

satisfy the conditions

$$
\begin{aligned}
& \alpha_{1}(0)=\left(\begin{array}{c}
\alpha_{11}(0) \\
\ldots \\
\alpha_{1 n}(0)
\end{array}\right)=\left(\begin{array}{c}
-1 \\
\ldots \\
-1
\end{array}\right) \\
& \alpha_{1}(a)=\left(\begin{array}{c}
\alpha_{11}(a) \\
\ldots \\
\alpha_{1 n}(a)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right) \\
& D_{q^{-1}} \alpha_{1}(0)=\left(\begin{array}{c}
D_{q^{-1}} \alpha_{11}(0) \\
\ldots \\
D_{q^{-1}} \alpha_{1 n}(0)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right), \\
& D_{q^{-1}} \alpha_{1}(a)=\left(\begin{array}{c}
D_{q^{-1}} \alpha_{11}(a) \\
\ldots \\
D_{q^{-1}} \alpha_{1 n}(a)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right),
\end{aligned}
$$

$$
\begin{aligned}
& \alpha_{2}(0)=\left(\begin{array}{c}
\alpha_{21}(0) \\
\ldots \\
\alpha_{2 n}(0)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right), \\
& \alpha_{2}(a)=\left(\begin{array}{c}
\alpha_{21}(a) \\
\ldots \\
\alpha_{2 n}(a)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right) \text {, } \\
& D_{q^{-1}} \alpha_{2}(0)=\left(\begin{array}{c}
D_{q^{-1}} \alpha_{21}(0) \\
\ldots \\
D_{q^{-1}} \alpha_{2 n}(0)
\end{array}\right)=\left(\begin{array}{c}
1 \\
\ldots \\
1
\end{array}\right), \\
& D_{q^{-1}} \alpha_{2}(a)=\left(\begin{array}{c}
D_{q^{-1}} \alpha_{21}(a) \\
\ldots \\
D_{q^{-1}} \alpha_{2 n}(a)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right), \\
& \beta_{1}(0)=\left(\begin{array}{c}
\beta_{11}(0) \\
\cdots \\
\beta_{1 n}(0)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\cdots \\
0
\end{array}\right), \\
& \beta_{1}(a)=\left(\begin{array}{c}
\beta_{11}(a) \\
\ldots \\
\beta_{1 n}(a)
\end{array}\right)=\left(\begin{array}{c}
1 \\
\ldots \\
1
\end{array}\right), \\
& D_{q^{-1}} \beta_{1}(0)=\left(\begin{array}{c}
D_{q^{-1}} \beta_{11}(0) \\
\ldots \\
D_{q^{-1}} \beta_{1 n}(0)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right), \\
& D_{q^{-1}} \beta_{1}(a)=\left(\begin{array}{c}
D_{q^{-1}} \beta_{11}(a) \\
\ldots \\
D_{q^{-1}} \beta_{1 n}(a)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right), \\
& \beta_{2}(0)=\left(\begin{array}{c}
\beta_{21}(0) \\
\cdots \\
\beta_{2 n}(0)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\cdots \\
0
\end{array}\right), \\
& \beta_{2}(a)=\left(\begin{array}{c}
\beta_{21}(a) \\
\ldots \\
\beta_{2 n}(a)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right), \\
& D_{q^{-1}} \beta_{2}(0)=\left(\begin{array}{c}
D_{q^{-1}} \beta_{21}(0) \\
\ldots \\
D_{q^{-1}} \beta_{2 n}(0)
\end{array}\right)=\left(\begin{array}{c}
0 \\
\ldots \\
0
\end{array}\right), \\
& D_{q^{-1}} \beta_{2}(a)=\left(\begin{array}{c}
D_{q^{-1}} \beta_{21}(a) \\
\ldots \\
D_{q^{-1}} \beta_{2 n}(a)
\end{array}\right)=\left(\begin{array}{c}
1 \\
\ldots \\
1
\end{array}\right) \text {. }
\end{aligned}
$$

belongs to the set $\mathcal{D}_{\max }$ and $\Pi_{1} \mathcal{Y}=\Lambda, \Pi_{2} \mathcal{Y}=\Gamma$. This completes the proof.
Now, we give the following definition.
Definition 5.3 ([29]) Let $\mathcal{L}$ be a linear operator with dense domain $\mathcal{D}(\mathcal{L})$ acting on some Hilbert space $\mathcal{H}$. The operator $\mathcal{L}$ is called dissipative if

$$
\operatorname{Im}(\mathcal{L} f, f) \geq 0
$$

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

for all $f \in \mathcal{D}(\mathcal{L})$ and is called maximal dissipative if it does not have a proper dissipative extension. Similarly, The operator $\mathcal{L}$ is called accumulative

$$
\operatorname{Im}(\mathcal{L} f, f) \leq 0
$$

for all $f \in \mathcal{D}(\mathcal{L})$ and is called maximal accumulative if it does not have a proper accumulative extension.
Let

$$
\begin{align*}
& D_{1}=\left\{\xi \in \mathcal{D}_{\max }:(M-I) \Pi_{1} \xi+i(M+I) \Pi_{2} \xi=0\right\},  \tag{5.2}\\
& D_{2}=\left\{\xi \in \mathcal{D}_{\max }:(M-I) \Pi_{1} \xi+i(M+I) \Pi_{2} \xi=0\right\}, \tag{5.3}
\end{align*}
$$

where $M$ is a contraction operator in $E \oplus E$.
Then by Theorem 5.2, the following theorem is obtained [29].

Theorem 5.4 The restriction of the operator $T_{\max }$ to the set $D_{1}$ is a maximal dissipative extension of the symmetric operator $T_{\min }$. Conversely, any maximal dissipative extensions of $T_{\min }$ is the restriction of $T_{\max }$ to a set $D_{1}$. Similarly, the restriction of the operator $T_{\max }$ to the set $D_{2}$ is a maximal accumulative extension of the symmetric operator $T_{\min }$. Conversely, any maximal accumulative extensions of $T_{\min }$ is the restriction of $T_{\max }$ to a set $D_{2}$. Here, the contraction $M$ is uniquely determined by the extension. If the operator $M$ is unitary, these conditions define a self-adjoint extension of $T_{\min }$.

## References

[1] Adıvar M, Bohner M. Spectral analysis of $q$-difference equations with spectral singularities. Mathematical and Computer Modelling 2006; 43 (7-8): 69 5-703. doi: $10.1016 / \mathrm{j} . \mathrm{mcm}$.2005.04.014
[2] Allahverdiev BP, Tuna H. An expansion theorem for $q$ - Sturm-Liouville operators on the whole line. Turkish Journal of Mathematics 2018; 42: 1060-1071. doi: 10.3906/mat-1705-22
[3] Allahverdiev BP, Tuna H. Limit-point criteria for $q$-Sturm-Liouville equations. Quaestiones Mathematicae 2019; 42 (10): 1291-1299. doi: 10.2989/16073606.2018.1514541
[4] Allahverdiev BP, Tuna H. Qualitative spectral analysis of singular $q$-Sturm-Liouville operators. Bulletin of the Malaysian Mathematical Sciences Society 2020; 43: 1391. doi: 10.1007/s40840-019-00747-3
[5] Allahverdiev BP, Tuna H. Eigenfunction expansion in the singular case for $q$-Sturm-Liouville operators. Caspian Journal of Mathematical Sciences 2019; 8 (2): 91-102. doi: 10.22080/CJMS.2018.13943.1339
[6] Allakhverdiev BP. On extensions of symmetric Schrödinger operators with a matrix potential, Izvestiya Rossiĭskoĭ Akademii Nauk. Seriya Matematicheskaya 1995; 59: 19-54; English transl. Izvestiya: Mathematics 1995; 59: 45-62. doi: 10.1070/IM1995v059n01ABEH000002
[7] Annaby MH, Mansour ZS. q- Fractional Calculus and Equations. Lecture Notes in Mathematics. vol. 2056, Berlin, Germany: Springer-Verlag, 2012.
[8] Annaby MH, Mansour ZS. Basic Sturm-Liouville problems. Journal of Physics. A. Mathematical and Theoretical 2005; 38: 3775-3797. doi: 10.1088/0305-4470/38/17/005
[9] Aygar Y, Bairamov E. Jost solution and the spectral properties of the matrix-valued difference operators. Applied Mathematics and Computation 2012; 218 (3): 9676-9681. doi: 10.1016/j.amc.2012.02.081

## PAŞAOĞLU ALLAHVERDİEV and TUNA/Turk J Math

[10] Aygar Y, Bohner M. Spectral analysis of a matrix-valued quantum-difference operator. Dynamic Systems and Applications 2016; 25: 1-9.
[11] Aygar Y, Bohner M. A Polynomial-Type Jost Solution and spectral properties of a self-adjoint quantum-difference operator. Complex Analysis and Operator Theory 2016; 10 (6): 1171-1180. doi: 10.1007/s11785-015-0463-x
[12] Aygar Y. A research on spectral analysis of a matrix quantum difference equations with spectral singularities. Quaestiones Mathematicae 2017; 40 (2): 245-249. doi: 10.2989/16073606.2017.1284911
[13] Bairamov E, Cebesoy Ş. Spectral singularities of the matrix Schrödinger equations. Hacettepe Journal of Mathematics and Statistics 2016; 45 (4): 1007-1014. doi: 10.15672/HJMS. 20164514275
[14] Bairamov E, Aygar Y, Cebesoy S. Spectral analysis of a self-adjoint matrix-valued discrete operator on the whole axis. Journal of Nonlinear Sciences and Applications 2016; 9 (6): 4257-4262.
[15] Bastard G, Brum JA. Electronic states in semi conductor heterostructures. IEEE Journal of Quantum Electronics 1986; 22: 1625-1644. doi: 10.1109/JQE.1986.1073186
[16] Bastard G. Wave mechanics applied to semi conductor hetero structures. Paris, Éditions de Physique: 1989.
[17] Beals R, Henkin GM, Novikova NN. The inverse boundary problem for the Rayleigh system. Journal of Mathematical Physics 1965; 36 (12): 6688-6708. doi: 10.1063/1.531182
[18] Bondarenko N. Spectral analysis for the matrix Sturm-Liouville operator on a finite interval. Tamkang Journal of Mathematics 2011; 42 (3): 305-327. doi: 10.5556/j.tkjm.42.2011.305-327
[19] Bondarenko N. Matrix Sturm-Liouville equation with a Bessel-type singularity on a finite interval. Analysis and Mathematical Physics 2017; 7 (1): 77-92. doi: 10.1007/s13324-016-0131-y
[20] Boutet de Monvel A, Shepelsky D. Inverse scattering problem for anisotropic media. Journal of Mathematical Physics 1995; 36 (7): 3443-3453. doi: 10.1063/1.530971
[21] Bruk VM. On a class of boundary -value problemswith a spectral parameter in the boundary conditions, Matematicheskiĭ Sbornik 1976; 100: 210-216. doi: 10.1070/SM1976v029n02ABEH003662
[22] Calkin JW. Abstract symmetric boundary conditions. Transactions of the American Mathematical Society 1939; 45 (3): 369-442.
[23] Chabanov VM. Recovering the M-channel Sturm-Liouville operator from M +1 spectra. Journal of Mathematical Physics 2004; 45 (11): 4255-4260. doi: 10.1063/1.1794844
[24] Coskun C, Olgun M. Principal functions of non-selfadjoint matrix Sturm-Liouville equations. Journal of Computational and Applied Mathematics 2011; 235 (16): 4834-4838. doi: 10.1016/j.cam.2010.12.004
[25] Ernst T. The History of $q$ - Calculus and a New Method. U. U. D. M. Report (2000): 16, ISSN1101-3591, Department of Mathematics, Uppsala University, 2000.
[26] Eryılmaz A, Tuna H. Spectral theory of dissipative $q$-Sturm-Liouville problems. Studia Scientiarum Mathematicarum Hungarica 2014; 51 (3): 366-383. doi: 10.1556/SScMath.51.2014.3.1289
[27] Gorbachuk ML. On spectral functions of a second order differential operator with operator coefficients. Ukrains'kyi Matematychnyi Zhurnal 1966; 18 (2): 3-21; English transl. American Mathematical Society Translations: Series 2 1968; 72: 177-202.
[28] Gorbachuk ML, Gorbachuk VI, Kochubei AN. The theory of extensions of symmetric operators and boundaryvalue problems for differential equations. Ukrains'kyi Matematychnyi Zhurnal 1989; 41: 1299-1312; English transl. in Ukrainian Mathematical Journal 1989; 41: 1117-1129. doi: 10.1007/BF01057246
[29] Gorbachuk ML, Gorbachuk VI. Boundary Value Problems for Operator Differential Equations, Naukova Dumka, Kiev, 1984; English transl. Birkhauser Verlag, 1991.
[30] Kac V, Cheung P. Quantum Calculus. Berlin, Germany: Springer-Verlag, 2002.
[31] Karahan D, Mamedov KhR. Sampling theory associated with $q$-Sturm-Liouville operator with discontinuity conditions. Journal of Contemporary Applied Mathematics 2020; 10 (2): 1-9.
[32] Kochubei AN. Extensions of symmetric operators and symmetric binary relations. Matematicheskie Zametki 1975; 17: 41-48; English transl. in Mathematical Notes 1975; 17: 25-28. doi: 10.1007/BF01093837
[33] Krall AM. Hilbert Space, Boundary Value Problems and Orthogonal Polynomials. Berlin, Germany: Birkhäuser Verlag, 2002.
[34] Krein MG. On the indeterminate case of the Sturm-Liouville boundaryvalue problem in the interval ( $0, \infty$ ), Izvestiya Akademii Nauk SSSR. Seriya Matematicheskaya 1952; 16: 292-324.
[35] Maksudov FG, Allahverdiev BP. On the extensions of Schrödinger operators with a matrix potentials, Doklady Akademii Nauk 1993; 332(1):18-20;English transl. Russian Academy of Sciences. Doklady. Mathematics 1994; 48 (2):240-243. doi: 10.1070/IM1995v059n01ABEH000002
[36] Malamud MM, Mogilevskiy VI. On extensions of dual pairs of operators. Dopovidi Natsional'noï Akademiï Nauk Ukraïny 1997; 1: 30-37.
[37] Mogilevskiy VI. On proper extensions of a singular differential operator in a space of vector functions. Dopovidi Natsional'noï Akademiï Nauk Ukraïny 1994; 9: 29-33 (in Russian with an English abstract) . doi: 10.1007/BF01059050
[38] Naimark MA. Linear Differential Operators. 2nd ed. Moscow, USSR: Nauka, 1969 (in Russian).
[39] Rofe-Beketov FS. Self-adjoint extensions of differential operators in a space of vector valued functions, Doklady Akademii Nauk SSSR 1969; 184: 1034-1037; English transl. in Soviet Mathematics. Doklady 1969; 10:188-192.
[40] Shi YM. Spectral theory of discrete linear Hamiltonian systems. Journal of Mathematical Analysis and Applications 2004; 289 (2): 554-570. doi: 10.1016/j.jmaa.2003.08.039
[41] Tuna H, Eryılmaz A. Completeness of the system of root functions of $q$-Sturm-Liouville operators. Mathematical Communications 2014; 19 (1): 65-73.
[42] Tuna H, Eryılmaz A. On $q$-Sturm-Liouville operators with eigenvalue parameter contained in the boundary conditions. Dynamic Systems and Applications 2015; 24 (4): 491-501.
[43] Tuna H, Eryılmaz A. Livšic's theorem for $q-$ Sturm-Liouville operators. Studia Scientiarum Mathematicarum Hungarica 2016; 53 (4): 512-524. doi: 10.1556/012.2016.53.4.1348
[44] von Neumann J. Allgemeine Eigenwertheorie Hermitischer Functionaloperatoren. Mathematische Annalen 1929; 102: 49-131 (in German). doi: 10.1007/BF01782338
[45] Yardimci S. A note on the spectral singularities of non-selfadjoint matrix-valued difference operators. Journal of Computational and Applied Mathematics 2010; 234: 3039-3042. doi: 10.1016/j.cam.2010.04.017
[46] Yurko V. Inverse problems for the matrix Sturm-Liouville equation on a finite interval. Inverse Problems 2006; 22: 1139-1149. doi: 10.1088/0266-5611/22/4/002
[47] Zettl A. Sturm-Liouville Theory. Mathematical Surveys and Monographs. vol. 121. Providence, Rhode Island, USA: American Mathematical Society, 2005.


[^0]:    *Correspondence: hustuna@gmail.com
    2010 AMS Mathematics Subject Classification: 34B20, 34B24, 47B25

