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Research Article

The dual spaces of variable anisotropic Hardy–Lorentz spaces and continuity of a class of linear operators

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Abstract: In this paper, the authors obtain the continuity of a class of linear operators on variable anisotropic Hardy–Lorentz spaces. In addition, the authors also obtain that the dual space of variable anisotropic Hardy–Lorentz spaces is the anisotropic BMO-type spaces with variable exponents. This result is still new even when the exponent function $p(\cdot)$ is p.

Key words: Anisotropy, Hardy-Lorentz space, atom, Calderón-Zygmund operator, BMO space

1. Introduction

As is known to all, Hardy space on the Euclidean space \mathbb{R}^n is a good substitutes of Lebesgue space $L^p(\mathbb{R}^n)$ when $p \in (0, 1]$, and plays an important role in haronmic analysis and **PDE**s; see, for examples, [5, 11, 14, 20-22, 25]. Moreover, when studying the boundedness of some operators in the critical case, the weak Hardy space $w\mathcal{H}^p(\mathbb{R}^n)$ naturally appears and it is a good substitute of $\mathcal{H}^p(\mathbb{R}^n)$. $w\mathcal{H}^p(\mathbb{R}^n)$ with $p \in (0, 1)$ was first introduced by Fefferman and Soria [10] to find out the biggest space from which the Riesz transform is bounded to the weak Lebesgue space $wL^1(\mathbb{R}^n)$. In 2007, Abu-Shammala and Torchinsky [1] introduced the Hardy–Lorentz spaces $\mathcal{H}^{p,r}(\mathbb{R}^n)$ for the full range $p \in (0, 1]$ and $r \in (0, \infty]$, and obtained some real-variable characterizations of this space. In 2016, Liu et al. [16] introduced the anisotropic Hardy–Lorentz space of Abu-Shammala and Torchinsky.

As a generalization, variable exponent function spaces have their applications in fluid dynamics [2], image processing [4], **PDE**s and variational calculus [9, 25]. Let $p(\cdot) : \mathbb{R}^n \to (0, \infty)$ be a variable exponent function. Recently, Liu et al. [17] introduced the variable anisotropic Hardy–Lorentz space $\mathcal{H}_A^{p(\cdot), r}(\mathbb{R}^n)$, via the radial grand maximal function, and then established its some real-variable characterizations, respectively, in terms of atom, the radial and the nontangential maximal functions. For more information about variable function spaces, see [6–8, 13, 18, 19, 24, 27].

To complete the theory of the variable anisotropic Hardy–Lorentz space $\mathcal{H}_{A}^{p(\cdot), r}(\mathbb{R}^{n})$, in this article, we obtain the boundedness of a class of Calderón–Zygmund operators from $\mathcal{H}_{A}^{p(\cdot), r}(\mathbb{R}^{n})$ to variable Lorentz space

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 $L^{p(\cdot),r}(\mathbb{R}^n)$ and from $\mathcal{H}^{p(\cdot),r}_A(\mathbb{R}^n)$ to itself. In addition, we also obtain the dual space of $\mathcal{H}^{p(\cdot),r}_A(\mathbb{R}^n)$ is the anisotropic BMO-type space with variable exponents.

Precisely, this article is organized as follows.

In Section 2, we recall some notations and definitions concerning expansive dilations, the variable Lorentz space $L^{p(\cdot), r}(\mathbb{R}^n)$ and the variable anisotropic Hardy–Lorentz space $\mathcal{H}^{p(\cdot), r}_A(\mathbb{R}^n)$, via the radial grand maximal function.

Section 3 is devoted to establishing the boundedness of anisotropic convolutional δ -type Calderón–Zygmund operators from $\mathcal{H}^{p(\cdot), r}_{A}(\mathbb{R}^{n})$ to $L^{p(\cdot), r}(\mathbb{R}^{n})$ and from $\mathcal{H}^{p(\cdot), r}_{A}(\mathbb{R}^{n})$ to itself.

In Section 4, we prove that the dual space of $\mathcal{H}_{A}^{p(\cdot), r}(\mathbb{R}^{n})$ is the anisotropic BMO-type space with variable exponents (see Theorem 4.6). For this purpose, we first introduce a new kind of anisotropic BMO-type spaces with variable exponents $\mathcal{BMO}_{A}^{p(\cdot), q, s}(\mathbb{R}^{n})$ in Definition 4.1, which includes the space $\mathrm{BMO}(\mathbb{R}^{n})$ of John and Nirenberg [12]. It is worth pointing out that this result is also new, when $\mathcal{H}_{A}^{p(\cdot), r}(\mathbb{R}^{n})$ is reduced to $\mathcal{H}_{A}^{p, r}(\mathbb{R}^{n})$.

Finally, we make some conventions on notation. Let $\mathbb{N} := \{1, 2, ...\}$ and $\mathbb{Z}_+ := \{0\} \cup \mathbb{N}$. For any $\alpha := (\alpha_1, \ldots, \alpha_n) \in \mathbb{Z}_+^n := (\mathbb{Z}_+)^n$, let $|\alpha| := \alpha_1 + \cdots + \alpha_n$ and $\partial^{\alpha} := \left(\frac{\partial}{\partial x_1}\right)^{\alpha_1} \cdots \left(\frac{\partial}{\partial x_n}\right)^{\alpha_n}$. In this article, we denote by C a *positive constant* which is independent of the main parameters, but it may vary from line to line. For any $q \in [1, \infty]$, we denote by q' its conjugate index. For any $a \in \mathbb{R}$, $\lfloor a \rfloor$ denotes the *maximal integer* not larger than a. The symbol $D \leq F$ means that $D \leq CF$. If $D \leq F$ and $F \leq D$, we then write $D \sim F$. If a set $E \subset \mathbb{R}^n$, we denote by χ_E its characteristic function. If there are no special instructions, any space $\mathcal{X}(\mathbb{R}^n)$ is denoted simply by \mathcal{X} .

2. Preliminaries

Firstly, we recall the definitions of anisotropic dilations on \mathbb{R}^n ; see [3, p. 5]. A real $n \times n$ matrix A is called an *anisotropic dilation*, shortly a *dilation*, if $\min_{\lambda \in \sigma(A)} |\lambda| > 1$, where $\sigma(A)$ denotes the set of all *eigenvalues* of A. Let λ_- and λ_+ be two *positive numbers* such that

$$1 < \lambda_{-} < \min\{|\lambda| : \lambda \in \sigma(A)\} \le \max\{|\lambda| : \lambda \in \sigma(A)\} < \lambda_{+}$$

By [3, Lemma 2.2], we know that, for a given dilation A, there exist a number $r \in (1, \infty)$ and a set $\Delta := \{x \in \mathbb{R}^n : |Px| < 1\}$, where P is some nondegenerate $n \times n$ matrix, such that

$$\Delta \subset r\Delta \subset A\Delta,$$

and we can assume that $|\Delta| = 1$, where $|\Delta|$ denotes the *n*-dimensional Lebesgue measure of the set Δ . Let $B_k := A^k \Delta$ for $k \in \mathbb{Z}$. Then B_k is open,

$$B_k \subset rB_k \subset B_{k+1}$$
 and $|B_k| = b^k$

here and hereafter, $b := |\det A|$. An ellipsoid $x + B_k$ for some $x \in \mathbb{R}^n$ and $k \in \mathbb{Z}$ is called a *dilated ball*. Denote

$$\mathfrak{B} := \{ x + B_k : \ x \in \mathbb{R}^n, \ k \in \mathbb{Z} \}.$$

$$(2.1)$$

Throughout the whole paper, let σ be the *smallest integer* such that $2B_0 \subset A^{\sigma}B_0$ and, for any subset E of \mathbb{R}^n , let $E^{\complement} := \mathbb{R}^n \setminus E$. Then, for all $k, j \in \mathbb{Z}$ with $k \leq j$, it holds true that

$$B_k + B_j \subset B_{j+\sigma},\tag{2.2}$$

$$B_k + (B_{k+\sigma})^{\complement} \subset (B_k)^{\complement}, \tag{2.3}$$

where E + F denotes the algebraic sum $\{x + y : x \in E, y \in F\}$ of sets $E, F \subset \mathbb{R}^n$.

Recall a quasi-norm, associated with dilation A, is a Borel measurable mapping $\rho : \mathbb{R}^n \to [0,\infty)$, satisfying

- (i) $\rho(x) > 0$ for all $x \in \mathbb{R}^n \setminus \{\vec{0}_n\}$, here and hereafter, $\vec{0}_n$ denotes the origin of \mathbb{R}^n ;
- (ii) $\rho(Ax) = b\rho(x)$ for all $x \in \mathbb{R}^n$, where, as above, $b := |\det A|$;

(iii) $\rho(x+y) \leq H[\rho(x)+\rho(y)]$ for all $x, y \in \mathbb{R}^n$, where $H \in [1, \infty)$ is a constant independent of x and y.

By [3, Lemma 2.4], we know that all homogeneous quasi-norms associated with a given dilation A are equivalent. Therefore, for a fixed dilation A, in what follows, for convenience, we always use the *step* homogeneous quasi-norm ρ defined by setting, for all $x \in \mathbb{R}^n$,

$$\rho(x) := \sum_{k \in \mathbb{Z}} b^k \chi_{B_{k+1} \setminus B_k}(x) \text{ if } x \neq \vec{0}_n, \text{ or else } \rho(\vec{0}_n) := 0.$$

By (2.2), we know that, for all $x, y \in \mathbb{R}^n$,

$$\rho(x+y) \le b^{\sigma}[\rho(x) + \rho(y)]$$

Moreover, (\mathbb{R}^n, ρ, dx) is a space of homogeneous type in the sense of Coifman and Weiss [5], where dx denotes the *n*-dimensional Lebesgue measure.

A measurable function $p(\cdot) : \mathbb{R}^n \to (0, \infty)$ is called a *variable exponent*. For any variable exponent $p(\cdot)$, let

$$p_{-} := \underset{x \in \mathbb{R}^{n}}{\operatorname{ess inf}} p(x) \quad \text{and} \quad p_{+} := \underset{x \in \mathbb{R}^{n}}{\operatorname{ess sup}} p(x).$$

$$(2.4)$$

Denote by \mathcal{P} the set of all variable exponents $p(\cdot)$ satisfying $0 < p_{-} \leq p_{+} < \infty$.

Let f be a measurable function on \mathbb{R}^n and $p(\cdot) \in \mathcal{P}$. Define

$$\|f\|_{L^{p(\cdot)}} := \inf \left\{ \lambda \in (0, \infty) : \varrho_{p(\cdot)}(f/\lambda) \le 1 \right\},\$$

where

$$\varrho_{p(\cdot)}(f) := \int_{\mathbb{R}^n} |f(x)|^{p(x)} \, dx.$$

Moreover, the variable Lebesgue space $L^{p(\cdot)}$ is defined to be the set of all measurable functions f satisfying that $\varrho_{p(\cdot)}(f) < \infty$, equipped with the quasi-norm $\|f\|_{L^{p(\cdot)}}$.

Remark 2.1 [17] Let $p(\cdot) \in \mathcal{P}$.

(i) For any $r \in (0, \infty)$ and $f \in L^{p(\cdot)}$, $|||f|^r||_{L^{p(\cdot)}} = ||f||_{L^{rp(\cdot)}}^r$. Moreover, for any $\mu \in \mathbb{C}$ and $f, g \in L^{p(\cdot)}$, $||\mu f||_{L^{p(\cdot)}} = |\mu| ||f||_{L^{p(\cdot)}}$ and $||f + g||_{L^{p(\cdot)}}^{\underline{p}} \le ||f||_{L^{p(\cdot)}}^{\underline{p}} + ||g||_{L^{p(\cdot)}}^{\underline{p}}$, where

$$\underline{p} := \min\{p_{-}, 1\} \tag{2.5}$$

with p_{-} as in (2.4).

(ii) For any function $f \in L^{p(\cdot)}$ with $||f||_{L^{p(\cdot)}} > 0$, $\varrho_{p(\cdot)}(f/||f||_{L^{p(\cdot)}}) = 1$ and, for $||f||_{L^{p(\cdot)}} \le 1$, then $\varrho_{p(\cdot)}(f) \le ||f||_{L^{p(\cdot)}}$.

Definition 2.2 Let $p(\cdot) \in \mathcal{P}$. The variable Lorentz space $L^{p(\cdot), r}$ is defined to be the set of all measurable functions f such that

$$\|f\|_{L^{p(\cdot),r}} := \begin{cases} \left[\int_0^\infty \lambda^r \left\| \chi_{\{x \in \mathbb{R}^n : |f(x)| > \lambda\}} \right\|_{L^{p(\cdot)}}^r \frac{d\lambda}{\lambda} \right]^{1/r}, & r \in (0,\infty), \\ \sup_{\lambda \in (0,\infty)} \left[\lambda \left\| \chi_{\{x \in \mathbb{R}^n : |f(x)| > \lambda\}} \right\|_{L^{p(\cdot)}} \right], & r = \infty \end{cases}$$

is finite.

We say that $p(\cdot) \in \mathcal{P}$ satisfy the globally log-Hölder continuous condition, denoted by $p(\cdot) \in C^{\log}$, if there exist two positive constants $C_{\log}(p)$ and C_{∞} , and $p_{\infty} \in \mathbb{R}$ such that, for any $x, y \in \mathbb{R}^n$,

$$|p(x) - p(y)| \le \frac{C_{\log}(p)}{\log(e + 1/\rho(x - y))}$$

and

$$|p(x) - p_{\infty}| \le \frac{C_{\infty}}{\log(e + \rho(x))}$$

A C^{∞} function φ is said to belong to the Schwartz class S if, for every integer $\ell \in \mathbb{Z}_+$ and multiindex α , $\|\varphi\|_{\alpha,\ell} := \sup_{x \in \mathbb{R}^n} [\rho(x)]^{\ell} |\partial^{\alpha} \varphi(x)| < \infty$. The dual space of S, namely, the space of all tempered distributions on \mathbb{R}^n equipped with the weak-* topology, is denoted by S'. For any $N \in \mathbb{Z}_+$, let

$$\mathcal{S}_N := \{ \varphi \in \mathcal{S} : \|\varphi\|_{\alpha,\ell} \le 1, \ |\alpha| \le N, \ \ell \le N \}.$$

In what follows, for $\varphi \in \mathcal{S}$, $k \in \mathbb{Z}$ and $x \in \mathbb{R}^n$, let $\varphi_k(x) := b^{-k} \varphi \left(A^{-k} x \right)$.

Definition 2.3 Let $\varphi \in S$ and $f \in S'$. For any given $N \in \mathbb{N}$, the radial grand maximal function $M_N(f)$ of $f \in S'$ is defined by setting, for any $x \in \mathbb{R}^n$,

$$M_N(f)(x) := \sup_{\varphi \in \mathcal{S}_N} \sup_{k \in \mathbb{Z}} |f * \varphi_k(x)|.$$

Definition 2.4 [17] Let $p(\cdot) \in C^{\log}$, $r \in (0, \infty)$, A be a dilation and $N \in [\lfloor (1/\underline{p} - 1) \ln b / \ln \lambda_{-} \rfloor + 2, \infty)$, where \underline{p} is as in (2.5). The variable anisotropic Hardy-Lorentz space $\mathcal{H}_{A}^{p(\cdot), r}$ is defined as

$$\mathcal{H}_A^{p(\cdot), r} := \left\{ f \in \mathcal{S}' : M_N(f) \in L^{p(\cdot), r} \right\}$$

and, for any $f \in \mathcal{H}_{A}^{p(\cdot), r}$, let $||f||_{\mathcal{H}_{A}^{p(\cdot), r}} := ||M_{N}(f)||_{L^{p(\cdot), r}}$.

Remark 2.5 Let $p(\cdot) \in C^{\log}$, $r \in (0, \infty)$.

- (i) When $p(\cdot) := p$, where $p \in (0, \infty)$, the space $\mathcal{H}_A^{p(\cdot), r}$ is reduced to the anisotropic Hardy-Lorentz space $\mathcal{H}_A^{p, r}$ studied in [16].
- (ii) When $A := 2I_{n \times n}$ and $p(\cdot) := p$, the space $\mathcal{H}_A^{p(\cdot), r}$ is reduced to the Hardy–Lorentz space $\mathcal{H}^{p, r}$ studied in [1].

Definition 2.6 [17] Let $p(\cdot) \in \mathcal{P}$, $q \in (1, \infty]$ and

$$s \in \left[\lfloor (1/p_{-} - 1) \ln b / \ln \lambda_{-} \rfloor, \infty \right) \cap \mathbb{Z}_{+}$$

with p_{-} as in (2.4). An anisotropic $(p(\cdot), q, s)$ -atom is a measurable function a on \mathbb{R}^{n} satisfying

- (i) (support) supp $a := \overline{\{x \in \mathbb{R}^n : a(x) \neq 0\}} \subset B$, where $B \in \mathfrak{B}$ and \mathfrak{B} is as in (2.1);
- (ii) (size) $||a||_{L^q} \le \frac{|B|^{1/q}}{||\chi_B||_{L^p(\cdot)}};$
- (iii) (vanishing moment) $\int_{\mathbb{R}^n} a(x) x^{\alpha} dx = 0$ for any $\alpha \in \mathbb{Z}^n_+$ with $|\alpha| \leq s$.

Definition 2.7 [17] Let $p(\cdot) \in C^{\log}$, $r \in (0, \infty)$, $q \in (1, \infty]$, $s \in [\lfloor (1/p_{-} - 1)\ln b/\ln \lambda_{-} \rfloor, \infty) \cap \mathbb{Z}_{+}$ with p_{-} as in (2.4) and A be a dilation. The anisotropic variable atomic Hardy–Lorentz space $\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r}$ is defined to be the set of all distributions $f \in \mathcal{S}'$ satisfying that there exists a sequence of $(p(\cdot), q, s)$ -atoms, $\{a_i^k\}_{i \in \mathbb{N}, k \in \mathbb{Z}}$, supported, respectively, on $\{x_i^k + B_{\ell_i^k}\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathfrak{B}$ and a positive constant \widetilde{C} such that $\sum_{i \in \mathbb{N}} \chi_{x_i^k + A^{j_0} B_{\ell_i^k}}(x) \leq \widetilde{C}$ for any $x \in \mathbb{R}^n$ and $k \in \mathbb{Z}$, with some $j_0 \in \mathbb{Z} \setminus \mathbb{N}$, and

$$f = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k \quad \text{in} \quad \mathcal{S}',$$

where $\lambda_i^k \sim 2^k \|\chi_{x_i^k + B_{\ell_i^k}}\|_{L^{p(\cdot)}}$ for any $k \in \mathbb{Z}$ and $i \in \mathbb{N}$ with the equivalent positive constants independent of k and i.

Moreover, for any $f \in \mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r}$, define

$$\|f\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}}} := \inf\left[\sum_{k \in \mathbb{Z}} \left\| \left\{ \sum_{i \in \mathbb{N}} \left[\frac{\lambda_i^k \chi_{x_i^k + B_{\ell_i^k}}}{\left\| \chi_{x_i^k + B_{\ell_i^k}} \right\|_{L^{p(\cdot)}}} \right]^{\underline{p}} \right\}^{1/\underline{p}} \right\|_{L^{p(\cdot)}}^r \right]^{1/r}$$

3. The continuity of Calderón–Zygmund operators

In this section, we get the continuity of anisotropic convolutional δ -type Calderón–Zygmund operators from $\mathcal{H}_{A}^{p(\cdot), r}$ to $L^{p(\cdot), r}$ or from $\mathcal{H}_{A}^{p(\cdot), r}$ to itself.

Let $\delta \in (0, \frac{\ln \lambda_+}{\ln b})$. We call a linear operator T is an anisotropic convolutional δ -type Calderón–Zygmund operator, if T is bounded on L^2 with kernel $\mathcal{K} \in \mathcal{S}'$ coinciding with a locally integrable function on $\mathbb{R}^n \setminus \{\vec{0}_n\}$,

and satisfying that there exists a positive constant C such that, for any $x, y \in \mathbb{R}^n$ with $\rho(x) > b^{2\sigma}\rho(y)$,

$$|\mathcal{K}(x-y) - \mathcal{K}(x)| \le C \frac{[\rho(y)]^{\delta}}{[\rho(x)]^{1+\delta}}$$

For any $f \in L^2$, define $T(f)(x) := \text{p.v.} \mathcal{K} * f(x)$.

Theorem 3.1 Let $p(\cdot) \in C^{\log}$, $r \in (0, \infty)$ and $\delta \in (0, \frac{\ln \lambda_+}{\ln b})$. Assume that T is an anisotropic convolutional δ -type Calderón–Zygmund operator. If $p_- \in (\frac{1}{1+\delta}, 1)$ with p_- as in (2.4), then there exists a positive constant C such that, for any $\mathcal{H}^{p(\cdot), r}_A$,

- (i) $||T(f)||_{L^{p(\cdot),r}} \leq C ||f||_{\mathcal{H}^{p(\cdot),r}};$
- (ii) $||T(f)||_{\mathcal{H}^{p(\cdot), r}_{A}} \leq C ||f||_{\mathcal{H}^{p(\cdot), r}_{A}}.$

Remark 3.2 When $p(\cdot) := p$, Theorem 3.1 coincides with [16, Theorem 6.16].

To prove Theorem 3.1, we need some technical lemmas.

Lemma 3.3 [17, Theorem 4.8] Let $p(\cdot) \in C^{\log}$, $r \in (0, \infty)$, $q \in (\max\{p_+, 1\}, \infty]$ with p_+ as in (2.4) and $s \in [\lfloor (1/p_- - 1) \ln b / \ln \lambda_- \rfloor, \infty) \cap \mathbb{Z}_+$ with p_- as in (2.4). Then

$$\mathcal{H}_{A}^{p(\cdot), r} = \mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r}$$

with equivalent quasi-norms.

By the proof of [17, Theorem 4.8], we obtain the following conclusion, which plays an important role in this section.

Lemma 3.4 Let $p(\cdot) \in C^{\log}$, $r \in (0, \infty)$, $q \in (1, \infty)$ and $s \in [\lfloor (1/p_{-} - 1) \ln b / \ln \lambda_{-} \rfloor, \infty) \cap \mathbb{Z}_{+}$ with p_{-} as in (2.4). Then, for any $f \in \mathcal{H}_{A}^{p(\cdot), r} \cap L^{q}$, there exist $\{\lambda_{i}^{k}\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathbb{C}$, dilated balls $\{x_{i}^{k} + B_{\ell_{i}^{k}}\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathfrak{B}$ and $(p(\cdot), \infty, s)$ -atoms $\{a_{i}^{k}\}_{i \in \mathbb{N}, k \in \mathbb{Z}}$ such that

$$f = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k \text{ in } L^q \text{ and } \mathcal{H}_A^{p(\cdot), r},$$

where the series also converges almost everywhere.

Proof Let $f \in \mathcal{H}^{p(\cdot), r}_A \cap L^q$. For any $k \in \mathbb{Z}$, by the proof of [17, Theorem 4.8], we know that there exist

$$\left\{x_i^k\right\}_{i\in\mathbb{N}}\subset\Omega_k:=\left\{x\in\mathbb{R}^n:\ M_Nf(x)>2^k\right\},\ \left\{\ell_i^k\right\}_{i\in\mathbb{N},k\in\mathbb{Z}}\subset\mathbb{Z},$$

a sequence of $(p(\cdot), \infty, s)$ -atoms, $\{a_i^k\}_{k \in \mathbb{Z}, i \in \mathbb{N}}$, supported on $\{x_i^k + B_{\ell_i^k + 4\sigma}\}_{k \in \mathbb{Z}, i \in \mathbb{N}}$, respectively, and $\{\lambda_i^k\}_{k \in \mathbb{Z}, i \in \mathbb{N}} \subset \mathbb{C}$, such that

$$f = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k =: \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} b_i^k \quad \text{in} \quad \mathcal{S}',$$
(3.1)

and for any $k \in \mathbb{Z}$ and $i \in \mathbb{N}$, $\operatorname{supp} b_i^k \subset x_i^k + B_{\ell_i^k + 4\sigma} \subset \Omega_k$,

$$\left\|b_{i}^{k}\right\|_{L^{\infty}} \lesssim 2^{k} \text{ and } \sharp\left\{j \in \mathbb{N}: \left(x_{i}^{k} + B_{\ell_{i}^{k} + 4\sigma}\right) \cap \left(x_{j}^{k} + B_{\ell_{j}^{k} + 4\sigma}\right) \neq \emptyset\right\} \le R,$$

$$(3.2)$$

where R is as in [17, Lemma 4.7]. Moreover, by $f \in \mathcal{H}_A^{p(\cdot), r} \cap L^q$, we have, for almost every $x \in \Omega_k$, there exists a $k(x) \in \mathbb{Z}$ such that $2^{k(x)} < M_N f(x) \le 2^{k(x)+1}$. From this, $\operatorname{supp} b_i^k \subset \Omega_k$ and (3.2), we deduce that, for a.e. $x \in \mathbb{R}^n$,

$$\sum_{k\in\mathbb{Z}}\sum_{i\in\mathbb{N}} \left|b_i^k(x)\right| \sim \sum_{k\in\mathbb{Z}, k\in(-\infty, k(x)]}\sum_{i\in\mathbb{N}} \left|b_i^k(x)\right| \lesssim \sum_{k\in\mathbb{Z}, k\in(-\infty, k(x)]}\sum_{i\in\mathbb{N}} 2^k \chi_{x_i^k + B_{\ell_i^k + 4\sigma}}(x)$$

$$\sim \sum_{k\in(-\infty, k(x)]\cap\mathbb{Z}} 2^k \sim M_N f(x).$$
(3.3)

Therefore, there exists a subsequence of the series $\{\sum_{|k| < K} \sum_{i \in \mathbb{Z}} b_i^k\}_{K \in \mathbb{N}}$, denoted still by itself without loss of generality, which converges to some measurable function F almost everywhere in \mathbb{R}^n .

It follows from (3.3) that, for any $K \in \mathbb{N}$ and a.e. $x \in \mathbb{R}^n$,

.

$$\left| F(x) - \sum_{|k| < K} b_i^k(x) \right| \lesssim |F(x)| + \sum_{k \in \mathbb{Z}, k \in (-\infty, k(x)]} \sum_{i \in \mathbb{N}} \left| b_i^k(x) \right|$$
$$\lesssim |F(x)| + M_N f(x) \lesssim M_N f(x).$$

From this, the fact that $M_N(f) \in L^q$, and the dominated convergence theorem, we conclude that $F = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} b_i^k$ in L^q . By this and (3.3), we know $f = F \in L^q$ and hence

$$f = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} b_i^k$$
 in L^q and $\mathcal{H}_A^{p(\cdot), r}$

and also almost everywhere.

In what follows, we also need the definition of anisotropic Hardy-Littlewood maximal function $\mathcal{M}(f)$. For any $f \in L^1_{\text{loc}}$ and $x \in \mathbb{R}^n$,

$$\mathcal{M}(f)(x) := \sup_{x \in B \in \mathfrak{B}} \frac{1}{|B|} \int_{B} |f(z)| \, dz, \tag{3.4}$$

where \mathfrak{B} is as in (2.1).

Lemma 3.5 [17, Lemma 4.3] Let $q \in (1, \infty)$. Assume that $p(\cdot) \in C^{\log}$ satisfies $1 < p_{-} \leq p_{+} < \infty$, where p_{-} and p_{+} are as in (2.4). Then there exists a positive constant C such that, for any sequence $\{f_k\}_{k\in\mathbb{N}}$ of measurable functions,

$$\left\|\left\{\sum_{k\in\mathbb{N}}\left[\mathcal{M}(f_k)\right]^q\right\}^{1/q}\right\|_{L^{p(\cdot)}} \leq C \left\|\left(\sum_{k\in\mathbb{N}}|f_k|^q\right)^{1/q}\right\|_{L^{p(\cdot)}}$$

where \mathcal{M} denotes the Hardy–Littlewood maximal operator as in (3.4).

Lemma 3.6 [15, Lemma 4.5] Let $p(\cdot) \in C^{\log}$, $r \in (0, \infty)$ and $q \in (\max\{p_+, 1\}, \infty)$. Then $\mathcal{H}^{p(\cdot), r}_A \cap L^q$ is dense in $\mathcal{H}^{p(\cdot), r}_A$.

The following Lemma show that variable anisotropic Hardy–Lorentz space $\mathcal{H}_A^{p(\cdot), r}$ is complete. Its proof is similar to [26, Lemma 3.9], we only need to make some minor changes. To limit the length of this article, we omit the concrete details.

Lemma 3.7 Let $p(\cdot) \in C^{\log}$, $r \in (0, \infty)$. Then $\mathcal{H}_A^{p(\cdot), r}$ is complete.

Proof [Proof of Theorem 3.1]

By the density, we only prove that (i) holds true for any $f \in \mathcal{H}_A^{p(\cdot), r} \cap L^q$ with $q \in (1, \infty) \cap (p_+, \infty)$. For any $f \in \mathcal{H}_A^{p(\cdot), r} \cap L^q$, from Lemma 3.4, we know that there exist numbers $\{\lambda_i^k\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathbb{C}$ and a sequence of $(p(\cdot), q, s)$ -atom, $\{a_i^k\}_{i \in \mathbb{N}, k \in \mathbb{Z}}$, supported, respectively, on $\{x_i^k + B_{\ell_i^k}\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathfrak{B}$ such that

$$f = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k \text{ in } L^q$$

where $\lambda_i^k \sim 2^k \|\chi_{x_i^k + B_{\ell_i^k}}\|_{L^{p(\cdot)}}$ for any $k \in \mathbb{Z}$ and $i \in \mathbb{N}$, $\sum_{i \in \mathbb{N}} \chi_{x_i^k + A^{j_0} B_{\ell_i^k}}(x) \lesssim 1$ with some $j_0 \in \mathbb{Z} \setminus \mathbb{N}$ for any $x \in \mathbb{R}^n$ and $k \in \mathbb{Z}$, and,

$$\|f\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}}} \sim \left[\sum_{k \in \mathbb{Z}} 2^{kr} \left\| \sum_{i \in \mathbb{N}} \chi_{x_i^k + B_{\ell_i^k}} \right\|_{L^{p(\cdot)}}^r \right]^{1/r}.$$
(3.5)

By the fact that T is bounded on L^q , we have

$$T(f) = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k T(a_i^k) \text{ in } L^q$$

 Set

$$f = \sum_{k=-\infty}^{k_0-1} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k + \sum_{k=k_0}^{\infty} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k =: F_1 + F_2 \text{ in } L^q.$$

Then

$$\begin{split} \left\|\chi_{\{x\in\mathbb{R}^{n}:\,T(f)(x)>2^{k_{0}}\}}\right\|_{L^{p(\cdot)}} &\lesssim \left\|\chi_{\{x\in\mathbb{R}^{n}:\,T(F_{1})(x)>2^{k_{0}-1}\}}\right\|_{L^{p(\cdot)}} + \left\|\chi_{\{x\in E_{k_{0}}:\,T(F_{2})(x)>2^{k_{0}-1}\}}\right\|_{L^{p(\cdot)}} \\ &+ \left\|\chi_{\{x\in(E_{k_{0}})^{\complement}:\,T(f_{2})(x)>2^{k_{0}-1}\}}\right\|_{L^{p(\cdot)}} \\ &=: I_{1} + I_{2} + I_{3} \,, \end{split}$$
(3.6)

where

$$E_{k_0} := \bigcup_{k=k_0}^{\infty} \bigcup_{i \in \mathbb{N}} \left(x_i^k + A^{\sigma} B_{\ell_i^k} \right)$$

Therefore,

$$I_{1} \lesssim \left\| \chi_{\{x \in \mathbb{R}^{n}: \sum_{k=-\infty}^{k_{0}-1} \sum_{i \in \mathbb{N}} \lambda_{i}^{k} T(a_{i}^{k})(x) \chi_{x_{i}^{k}+A^{\sigma}B_{\ell_{i}^{k}}}(x) > 2^{k_{0}-2} \}} \right\|_{L^{p(\cdot)}} + \left\| \chi_{\{x \in \mathbb{R}^{n}: \sum_{k=-\infty}^{k_{0}-1} \sum_{i \in \mathbb{N}} \lambda_{i}^{k} T(a_{i}^{k})(x) \chi_{(x_{i}^{k}+A^{\sigma}B_{\ell_{i}^{k}})} \mathfrak{g}(x) > 2^{k_{0}-2} \}} \right\|_{L^{p(\cdot)}} =: I_{1,1} + I_{1,2}.$$

$$(3.7)$$

For the term $I_{1,1}$, from the fact that T is bounded on L^q , Remark 2.1, (3.5) and a similar proof of [17, (4.7)], we deduce that

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{I}_{1,1}\right)^{r}\right]^{1/r} \lesssim \left[\sum_{k\in\mathbb{Z}} 2^{kr} \left\|\sum_{i\in\mathbb{N}} \chi_{x_{i}^{k}+B_{\ell_{i}^{k}}}\right\|_{L^{p(\cdot)}}^{r}\right]^{1/r} \sim \|f\|_{\mathcal{H}^{p(\cdot),\,q,\,s,\,r}_{A,\,\mathrm{atom}}}.$$
(3.8)

For the term $I_{1,2}$, from the Hölder inequality and the size condition of $a_i^k(x)$, we conclude that, for any $x \in (x_i^k + A^{\sigma}B_{\ell_i^k})^{\complement}$,

$$\begin{split} \left| Ta_{i}^{k}(x) \right| &\leq \int_{x_{i}^{k}+B_{\ell_{i}^{k}+\sigma}} \left| \mathcal{K}(x-y) - \mathcal{K}(x-x_{i}^{k}) \right| \left| a_{i}^{k}(y) \right| dy \\ &\lesssim \int_{x_{i}^{k}+B_{\ell_{i}^{k}+\sigma}} \frac{\rho(y-x_{i}^{k})^{\delta}}{\rho(x-x_{i}^{k})^{1+\delta}} \left| a_{i}^{k}(y) \right| dy \lesssim \frac{\left| x_{i}^{k} + B_{\ell_{i}^{k}} \right|^{\delta}}{\rho(x-x_{i}^{k})^{1+\delta}} \left\| a_{i}^{k} \right\|_{L^{q}} \left| x_{i}^{k} + B_{\ell_{i}^{k}} \right|^{1/q'} \\ &\lesssim \frac{\left| x_{i}^{k} + B_{\ell_{i}^{k}} \right|^{1+\delta}}{\rho(x-x_{i}^{k})^{1+\delta}} \frac{1}{\left\| \chi_{x_{i}^{k}+B_{\ell_{i}^{k}}} \right\|_{L^{p(\cdot)}}} \lesssim \left[\mathcal{M}(\chi_{x_{i}^{k}+B_{\ell_{i}^{k}}})(x) \right]^{1+\delta} \frac{1}{\left\| \chi_{x_{i}^{k}+B_{\ell_{i}^{k}}} \right\|_{L^{p(\cdot)}}}. \end{split}$$

By this and a similar estimate of [17, p. 374], we obtain

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{I}_{1,2}\right)^r\right]^{1/r} \lesssim \left[\sum_{k\in\mathbb{Z}} 2^{kr} \left\|\sum_{i\in\mathbb{N}} \chi_{x_i^k + B_{\ell_i^k}}\right\|_{L^{p(\cdot)}}^r\right]^{1/r} \sim \|f\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}}}.$$

Therefore, it follows from (3.7) and (3.8) that

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{I}_{1}\right)^{r}\right]^{1/r} \lesssim \left\|f\right\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \text{ atom}}}.$$
(3.9)

For I_2 and I_3 , by a proof similar to those of [17, (4.12) and (4.13)], we obtain

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{I}_{2}\right)^{r}\right]^{1/r} \lesssim \left\|f\right\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}}} \text{ and } \left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{I}_{3}\right)^{r}\right]^{1/r} \lesssim \left\|f\right\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}}}.$$

$$(3.10)$$

Combining the estimate of (3.6), (3.9) and (3.10), we obtain

$$\begin{aligned} \|T(f)\|_{L^{p(\cdot),r}} &\sim \left[\sum_{k\in\mathbb{Z}} 2^{kr} \|\chi_{\{x\in\mathbb{R}^{n}: |Tf(x)|>2^{k}\}}\|_{L^{p(\cdot)}}^{r}\right]^{1/r} \\ &\lesssim \left[\sum_{k\in\mathbb{Z}} 2^{kr} (\mathbf{I}_{1})^{r}\right]^{1/r} + \left[\sum_{k\in\mathbb{Z}} 2^{kr} (\mathbf{I}_{2})^{r}\right]^{1/r} + \left[\sum_{k\in\mathbb{Z}} 2^{kr} (\mathbf{I}_{3})^{r}\right]^{1/r} \\ &\lesssim \|f\|_{\mathcal{H}^{p(\cdot),q,s,r}_{A,\operatorname{atom}}} \sim \|f\|_{\mathcal{H}^{p(\cdot),r}_{A}}, \end{aligned}$$

which implies that $T(f) \in L^{p(\cdot), r}$. This finishes the proof of Theorem 3.1(i).

Now we show Theorem 3.1(ii). By Lemma 3.6, we only need to prove that (ii) holds true for any $f \in \mathcal{H}_A^{p(\cdot), r} \cap L^q$ with $q \in (1, \infty) \cap (p_+, \infty)$. Let $f \in \mathcal{H}_A^{p(\cdot), r} \cap L^q$. From Lemma 3.4, we know that there exist numbers $\{\lambda_i^k\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathbb{C}$ and a sequence of $(p(\cdot), q, s)$ -atom, $\{a_i^k\}_{i \in \mathbb{N}, k \in \mathbb{Z}}$, supported, respectively, on $\{x_i^k + B_{\ell_i^k}\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathfrak{B}$ such that

$$f = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k \text{ in } L^q$$

where $\lambda_i^k \sim 2^k \|\chi_{x_i^k + B_{\ell_i^k}}\|_{L^{p(\cdot)}}$ for any $k \in \mathbb{Z}$ and $i \in \mathbb{N}$, $\sum_{i \in \mathbb{N}} \chi_{x_i^k + A^{j_0} B_{\ell_i^k}}(x) \lesssim 1$ with some $j_0 \in \mathbb{Z} \setminus \mathbb{N}$ for any $x \in \mathbb{R}^n$ and $k \in \mathbb{Z}$, and,

$$\|f\|_{\mathcal{H}^{p(\cdot),q,s,r}_{A,\operatorname{atom}}} \sim \left[\sum_{k\in\mathbb{Z}} 2^{kr} \left\|\sum_{i\in\mathbb{N}} \chi_{x_i^k + B_{\ell_i^k}}\right\|_{L^{p(\cdot)}}^r\right]^{1/r}.$$
(3.11)

By the fact that T is bounded on L^q , we have

$$T(f) = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k T(a_i^k) \text{ in } L^q.$$

 Set

$$f = \sum_{k=-\infty}^{k_0 - 1} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k + \sum_{k=k_0}^{\infty} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k =: F_1 + F_2 \text{ in } L^q.$$

Then

$$\begin{aligned} \left\| \chi_{\left\{ x \in \mathbb{R}^{n} : M_{N}(T(f))(x) > 2^{k_{0}} \right\}} \right\|_{L^{p(\cdot)}} \tag{3.12} \\ \lesssim \left\| \chi_{\left\{ x \in \mathbb{R}^{n} : M_{N}(T(F_{1}))(x) > 2^{k_{0}-1} \right\}} \right\|_{L^{p(\cdot)}} + \left\| \chi_{\left\{ x \in G_{k_{0}} : M_{N}(T(F_{2}))(x) > 2^{k_{0}-1} \right\}} \right\|_{L^{p(\cdot)}} \\ &+ \left\| \chi_{\left\{ x \in (G_{k_{0}})^{\mathfrak{c}} : M_{N}(T(f_{2}))(x) > 2^{k_{0}-1} \right\}} \right\|_{L^{p(\cdot)}} \\ &= : J_{1} + J_{2} + J_{3} \,, \end{aligned}$$

where M_N is as in Definition 2.3 and

$$G_{k_0} := \bigcup_{k=k_0}^{\infty} \bigcup_{i \in \mathbb{N}} \left(x_i^k + A^{\sigma} B_{\ell_i^k} \right).$$

Therefore,

$$\begin{aligned} \mathbf{J}_{1} \lesssim \left\| \chi_{\{x \in \mathbb{R}^{n} : \sum_{k=-\infty}^{k_{0}-1} \sum_{i \in \mathbb{N}} \lambda_{i}^{k} M_{N}(T(a_{i}^{k}))(x) \chi_{x_{i}^{k}+A^{\sigma}B_{\ell_{i}^{k}}}(x) > 2^{k_{0}-2} \}} \right\|_{L^{p(\cdot)}} \\ &+ \left\| \chi_{\{x \in \mathbb{R}^{n} : \sum_{k=-\infty}^{k_{0}-1} \sum_{i \in \mathbb{N}} \lambda_{i}^{k} M_{N}(T(a_{i}^{k}))(x) \chi_{(x_{i}^{k}+A^{\sigma}B_{\ell_{i}^{k}})}^{(s)}(x) > 2^{k_{0}-2} } \right\|_{L^{p(\cdot)}} \\ &=: \mathbf{J}_{1,1} + \mathbf{J}_{1,2} \,. \end{aligned}$$

$$(3.13)$$

For the term $J_{1,1}$, from the fact that M_N and T are bounded on L^q , Remark 2.1, (3.11) and a similar proof of [17, (4.7)], we deduce that

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{J}_{1,1}\right)^{r}\right]^{1/r} \lesssim \left[\sum_{k\in\mathbb{Z}} 2^{kr} \left\|\sum_{i\in\mathbb{N}} \chi_{x_{i}^{k}+B_{\ell_{i}^{k}}}\right\|_{L^{p(\cdot)}}^{r}\right]^{1/r} \sim \|f\|_{\mathcal{H}^{p(\cdot),\,q,\,s,\,r}_{A,\,\mathrm{atom}}}.$$
(3.14)

For the term $J_{1,2}$, from the Hölder inequality, the size condition of $a_i^k(x)$, and a similar proof of [20, p.117, Lemma], we conclude that, for any $x \in (x_i^k + A^{\sigma}B_{\ell_i^k})^{\complement}$,

$$\begin{split} M_{N}(Ta_{i}^{k})(x) &= \sup_{\varphi \in \mathcal{S}_{N}} \sup_{j \in \mathbb{Z}} \left| (\varphi_{j} * Ta_{i}^{k})(x) \right| = \sup_{\varphi \in \mathcal{S}_{N}} \sup_{j \in \mathbb{Z}} \left| (\varphi_{j} * \mathcal{K} * a_{i}^{k})(x) \right| \\ &\leq \sup_{\varphi \in \mathcal{S}_{N}} \sup_{j \in \mathbb{Z}} \int_{x_{i}^{k} + B_{\ell_{i}^{k} + \sigma}} \left| (\varphi_{j} * \mathcal{K})(x - y) - (\varphi_{j} * \mathcal{K})(x - x_{i}^{k}) \right| \left| a_{i}^{k}(y) \right| \, dy \\ &\lesssim \int_{x_{i}^{k} + B_{\ell_{i}^{k} + \sigma}} \frac{\rho(y - x_{i}^{k})^{\delta}}{\rho(x - x_{i}^{k})^{1 + \delta}} \left| a_{i}^{k}(y) \right| \, dy \lesssim \frac{\left| x_{i}^{k} + B_{\ell_{i}^{k}} \right|^{\delta}}{\rho(x - x_{i}^{k})^{1 + \delta}} \left\| a_{i}^{k} \right\|_{L^{q}} \left| x_{i}^{k} + B_{\ell_{i}^{k}} \right|^{1/q'} \\ &\lesssim \frac{\left| x_{i}^{k} + B_{\ell_{i}^{k}} \right|^{1 + \delta}}{\rho(x - x_{i}^{k})^{1 + \delta}} \frac{1}{\left\| \chi_{x_{i}^{k} + B_{\ell_{i}^{k}}} \right\|_{L^{p(\cdot)}}} \lesssim \left[\mathcal{M}(\chi_{x_{i}^{k} + B_{\ell_{i}^{k}}})(x) \right]^{1 + \delta} \frac{1}{\left\| \chi_{x_{i}^{k} + B_{\ell_{i}^{k}}} \right\|_{L^{p(\cdot)}}}. \end{split}$$

From this and a similar estimate of [17, p. 374], we deduce that

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{J}_{1,2}\right)^{r}\right]^{1/r} \lesssim \left[\sum_{k\in\mathbb{Z}} 2^{kr} \left\|\sum_{i\in\mathbb{N}} \chi_{x_{i}^{k}+B_{\ell_{i}^{k}}}\right\|_{L^{p(\cdot)}}^{r}\right]^{1/r} \sim \|f\|_{\mathcal{H}_{A,\operatorname{atom}}^{p(\cdot),\,q,\,s,\,r}}.$$

Therefore, it follows from (3.13) and (3.14) that

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{J}_{1}\right)^{r}\right]^{1/r} \lesssim \|f\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}}}.$$
(3.15)

For J_2 and J_3 , by a proof similar to those of [17, (4.12) and (4.13)], we obtain

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{J}_{2}\right)^{r}\right]^{1/r} \lesssim \left\|f\right\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}}} \text{ and } \left[\sum_{k\in\mathbb{Z}} 2^{kr} \left(\mathbf{J}_{3}\right)^{r}\right]^{1/r} \lesssim \left\|f\right\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}}}.$$

$$(3.16)$$

It follows from the estimates of (3.12), (3.15) and (3.16) that

$$\begin{aligned} \|T(f)\|_{\mathcal{H}^{p(\cdot),r}_{A}} &= \|M_{N}(T(f))\|_{L^{p(\cdot),r}} \\ &\sim \left[\sum_{k\in\mathbb{Z}} 2^{kr} \|\chi_{\{x\in\mathbb{R}^{n}: |M_{N}(Tf)(x)|>2^{k}\}}\|_{L^{p(\cdot)}}^{r}\right]^{1/r} \\ &\lesssim \left[\sum_{k\in\mathbb{Z}} 2^{kr} (J_{1})^{r}\right]^{1/r} + \left[\sum_{k\in\mathbb{Z}} 2^{kr} (J_{2})^{r}\right]^{1/r} + \left[\sum_{k\in\mathbb{Z}} 2^{kr} (J_{3})^{r}\right]^{1/r} \\ &\lesssim \|f\|_{\mathcal{H}^{p(\cdot),q,s,r}_{A,\operatorname{atom}}} \sim \|f\|_{\mathcal{H}^{p(\cdot),r}_{A}}, \end{aligned}$$

which implies that $T(f) \in \mathcal{H}_A^{p(\cdot), r}$. Therefore, we complete the proof of Theorem 3.1.

4. The dual space of variable anisotropic Hardy–Lorentz space $\mathcal{H}_{A}^{p(\cdot), r}$

In this section, we establish the dual space of $\mathcal{H}_{A}^{p(\cdot), r}$. More precisely, we prove that the dual space of $\mathcal{H}_{A}^{p(\cdot), r}$ is the variable anisotropic BMO-type space $\mathcal{BMO}_{A}^{p(\cdot), q, s}$.

Now, we define two new variable anisotropic BMO-type space as follows. In this article, for any $m \in \mathbb{Z}_+$, we use P_m to denote the set of polynomials on \mathbb{R}^n with order not more than m. For any $B \in \mathfrak{B}$ and any locally integrable function g on \mathbb{R}^n , we use $P_B^m(g)$ to denote the minimizing polynomial of g with degree not greater than m, which means that $P_B^m(g)$ is the unique polynomial $f \in P_m$ such that, for any $h \in P_m$,

$$\int_B h(x)(g(x) - f(x)) \, dx = 0.$$

Definition 4.1 Let A be a given dilation, $p(\cdot) \in \mathcal{P}$, s be a nonnegative integer and $q \in [1, \infty)$. Then the variable anisotropic BMO-type space $\mathcal{BMO}_A^{p(\cdot), q, s}$ is defined to be the set of all $f \in L^q_{loc}$ such that

$$\|f\|_{\mathcal{BMO}_{A}^{p(\cdot), q, s}} := \sup_{B \in \mathfrak{B}} \inf_{P \in P_{s}} \frac{|B|^{1-1/q}}{\|\chi_{B}\|_{L^{p(\cdot)}}} \left[\int_{B} |f(x) - P(x)|^{q} \, dx \right]^{1/q} < \infty$$

where \mathfrak{B} is as in (2.1).

Definition 4.2 Let A be a given dilation, $p(\cdot) \in \mathcal{P}$, s be a nonnegative integer and $q \in [1, \infty)$. Then the variable anisotropic BMO-type space $\widetilde{\mathcal{BMO}}_A^{p(\cdot), q, s}$ is defined to be the set of all $f \in L^q_{loc}$ such that

$$\|f\|_{\widetilde{\mathcal{BMO}}_{A}^{p(\cdot),\,q,\,s}} := \sup_{B \in \mathfrak{B}} \frac{|B|^{1-1/q}}{\|\chi_B\|_{L^{p(\cdot)}}} \left[\int_B |f(x) - P_B^s(f)(x)|^q \, dx \right]^{1/q} < \infty$$

where \mathfrak{B} is as in (2.1).

Lemma 4.3 [3, (8.9)] Let $q \in [1, \infty]$, A be a given dilation, $f \in L^q_{loc}$ and s be a nonnegative integer and $B \in \mathfrak{B}$. Then there exists a positive constant C, independent of f and B, such that

$$\sup_{x \in B} |P_B^s(f)(x)| \le C \frac{\int_B |f(x)| \, dx}{|B|}$$

Lemma 4.4 Let A be a given dilation, $p(\cdot) \in \mathcal{P}$, s be a nonnegative integer and $q \in [1, \infty)$. Then

$$\mathcal{BMO}_A^{p(\cdot), q, s} = \widetilde{\mathcal{BMO}}_A^{p(\cdot), q, s}$$

with equivalent quasi-norms.

Proof By the above definition, it is easy to see that

$$\mathcal{BMO}_A^{p(\cdot), q, s} \supseteq \widetilde{\mathcal{BMO}}_A^{p(\cdot), q, s}$$

Conversely, from Lemma 4.3 and the Hölder inequality, we conclude that, for any $B \in \mathfrak{B}$, $Q \in P_s$,

$$\begin{split} \left[\frac{1}{|B|} \int_{B} |P_{B}^{s}(Q-f)(x)|^{q} \, dx\right]^{1/q} &\lesssim \frac{1}{|B|} \int_{B} |Q(x) - f(x)| \, dx \\ &\lesssim \left[\frac{1}{|B|} \int_{B} |Q(x) - f(x)|^{q} \, dx\right]^{1/q} \end{split}$$

Therefore, by the Minkowski inequality, we obtain

$$\begin{aligned} &\frac{|B|}{\|\chi_B\|_{L^{p(\cdot)}}} \left[\frac{1}{|B|} \int_B |P_B^s(f)(x) - f(x)|^q \, dx\right]^{1/q} \\ &= \frac{|B|}{\|\chi_B\|_{L^{p(\cdot)}}} \left[\frac{1}{|B|} \int_B |P_B^s(Q - f)(x) + f(x) - Q(x)|^q \, dx\right]^{1/q} \\ &\lesssim \frac{|B|}{\|\chi_B\|_{L^{p(\cdot)}}} \left[\frac{1}{|B|} \int_B |Q(x) - f(x)|^q \, dx\right]^{1/q}, \end{aligned}$$

which implies that

$$\mathcal{BMO}_A^{p(\cdot), q, s} \subseteq \widetilde{\mathcal{BMO}}_A^{p(\cdot), q, s}.$$

This completes the proof of Lemma 4.4.

Lemma 4.5 Let A be a given dilation, $p(\cdot) \in C^{\log}$, $r \in (0, 1]$, s be a nonnegative integer and $q \in [1, \infty)$. Then, for any continuous linear functional \mathcal{L} on $\mathcal{H}_A^{p(\cdot), r} = \mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r}$,

$$\begin{split} \|\mathcal{L}\|_{(\mathcal{H}^{p(\cdot),\,q,\,s,\,r}_{A,\,\mathrm{atom}})^{*}} := &\sup\left\{|\mathcal{L}(f)| : \|f\|_{\mathcal{H}^{p(\cdot),\,q,\,s,\,r}_{A,\,\mathrm{atom}}} \leq 1\right\} \\ &= &\sup\{|\mathcal{L}(a)| : a \text{ is } (p(\cdot),\,q,\,s) - \mathrm{atom}\}, \end{split}$$

where $(\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r})^*$ denotes the dual space of $\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r}$.

Proof Let a be a $(p(\cdot), q, s)$ -atom. Then we have that $||a||_{\mathcal{H}^{p(\cdot), q, s, r}_{A, \text{atom}}} \leq 1$. Therefore,

$$\sup\{|\mathcal{L}(a)|: a \text{ is } (p(\cdot), q, s) - \operatorname{atom}\} \le \sup\left\{|\mathcal{L}(f)|: \|f\|_{\mathcal{H}^{p(\cdot), q, s, r}_{A}} \le 1\right\}$$

Moreover, let $f \in \mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r}$ and $\|f\|_{\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r}} \leq 1$. Then, for any $\varepsilon > 0$, we know that there exist $\{\lambda_i^k\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathbb{C}$ and a sequence of $(p(\cdot), q, s)$ -atoms, $\{a_i^k\}_{i \in \mathbb{N}, k \in \mathbb{Z}}$, supported, respectively, on $\{x_i^k + B_{\ell_i^k}\}_{i \in \mathbb{N}, k \in \mathbb{Z}} \subset \mathfrak{B}$ such that

$$f = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k$$
 in \mathcal{S}' and a. e.

and

$$\left[\sum_{k\in\mathbb{Z}} 2^{kr} \left\| \sum_{i\in\mathbb{N}} \chi_{x_i^k + B_{\ell_i^k}} \right\|_{L^{p(\cdot)}}^r \right]^{1/r} \le 1 + \varepsilon.$$

Therefore, from the boundedness of \mathcal{L} , $\lambda_i^k \sim 2^k \|\chi_{x_i^k + B_{\ell_i^k}}\|_{L^{p(\cdot)}}$ and $r \in (0, 1]$, we further conclude that

$$\begin{split} \mathcal{L}(g)| &\leq \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \left| \lambda_i^k \right| |\mathcal{L}(a_i^k)| \leq \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \left| \lambda_i^k \right| \sup\{|\mathcal{L}(a)| : a \text{ is } (p(\cdot), q, s) - \operatorname{atom}\} \\ &\lesssim \left[\sum_{k \in \mathbb{Z}} 2^{kr} \left\| \sum_{i \in \mathbb{N}} \chi_{x_i^k + B_{\ell_i^k}} \right\|_{L^{p(\cdot)}}^r \right]^{1/r} \sup\{|\mathcal{L}(a)| : a \text{ is } (p(\cdot), q, s) - \operatorname{atom}\} \\ &\lesssim (1 + \varepsilon) \sup\{|\mathcal{L}(a)| : a \text{ is } (p(\cdot), q, s) - \operatorname{atom}\}. \end{split}$$

Combined with the arbitrariness of ε and hence finishes the proof of Lemma 4.5.

For any $q \in [1, \infty]$ and $s \in \mathbb{Z}_+$. Denote by L^q_{comp} the set of all functions $f \in L^q$ with compact support and

$$L^{q,\,s}_{\text{comp}} := \left\{ f \in L^q_{\text{comp}} : \int_{\mathbb{R}^n} f(x) x^\alpha \, dx = 0, \, |\alpha| \le s \right\}$$

The main result of this section is as follows.

].

Theorem 4.6 Let A be a given dilation, $r \in (0, 1]$, $p(\cdot) \in C^{\log}$, $p_+ \in (0, 1]$, $q \in (\max\{p_+, 1\}, \infty)$ and $s \in [\lfloor (1/p_- - 1) \ln b / \ln \lambda_- \rfloor, \infty) \cap \mathbb{Z}_+$ with p_- as in (2.4). Then

$$(\mathcal{H}_{A}^{p(\cdot), r})^{*} = (\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r})^{*} = \mathcal{BMO}_{A}^{p(\cdot), q', s} = \widetilde{\mathcal{BMO}}_{A}^{p(\cdot), q', s}$$

in the following sense: for any $\psi \in \mathcal{BMO}_A^{p(\cdot), q', s}$ or $\widetilde{\mathcal{BMO}}_A^{p(\cdot), q', s}$, the linear functional

$$\mathcal{L}_{\psi}(g) := \int_{\mathbb{R}^n} \psi(x) g(x) \, dx, \tag{4.1}$$

initial defined for all $g \in L^{q, s}_{\text{comp}}$, has a bounded extension to $\mathcal{H}^{p(\cdot), q, s, r}_{A, \text{atom}} = \mathcal{H}^{p(\cdot), r}_{A}$.

Conversely, if \mathcal{L} is a bounded linear functional on $\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r} = \mathcal{H}_{A}^{p(\cdot), r}$, then \mathcal{L} has the form as in (4.1) with a unique $\psi \in \mathcal{BMO}_{A}^{p(\cdot), q', s}$ or $\widetilde{\mathcal{BMO}}_{A}^{p(\cdot), q', s}$. Moreover,

$$\|\psi\|_{\widetilde{\mathcal{BMO}}_{A}^{p(\cdot),\,q',\,s}} \sim \|\psi\|_{\mathcal{BMO}_{A}^{p(\cdot),\,q',\,s}} \sim \|\mathcal{L}_{\psi}\|_{(\mathcal{H}_{A,\,\mathrm{atom}}^{p(\cdot),\,q,\,s,\,r})^{*}},$$

where the implicit positive constants are independent of ψ .

Remark 4.7 We should point that, when $p(\cdot) := p \in (0, 1]$, this result is also new.

Proof [Proof of Theorem 4.6] By Lemmas 3.3 and 4.4, we only need to show

$$\mathcal{BMO}_A^{p(\cdot),\,q',\,s} = (\mathcal{H}_{A,\,\mathrm{atom}}^{p(\cdot),\,q,\,s,\,r})^*.$$

Firstly, we prove that

$$\mathcal{BMO}_A^{p(\cdot),\,q',\,s} \subset (\mathcal{H}_{A,\,\mathrm{atom}}^{p(\cdot),\,q,\,s,\,r})^*.$$

Let $\psi \in \mathcal{BMO}_A^{p(\cdot), q', s}$ and a be a $(p(\cdot), q, s)$ -atom with $\operatorname{supp} a \subset B \in \mathfrak{B}$. Then, by the vanishing moment condition of a, Hölder's inequality and the size condition of a, we obtain

$$\begin{aligned} \left| \int_{\mathbb{R}^{n}} \psi(x) a(x) \, dx \right| &= \inf_{P \in P_{s}} \left| \int_{B} (\psi(x) - P(x)) a(x) \, dx \right| \\ &\leq \|a\|_{L^{q}} \inf_{P \in P_{s}} \left[\int_{B} |\psi(x) - P(x)|^{q'} \, dx \right]^{1/q'} \\ &\leq \frac{|B|^{1/q}}{\|\chi_{B}\|_{L^{p(\cdot)}}} \inf_{P \in P_{s}} \left[\int_{B} |\psi(x) - P(x)|^{q'} \, dx \right]^{1/q'} \\ &\leq \|\psi\|_{\mathcal{BMO}_{A}^{p(\cdot), q', s}}. \end{aligned}$$
(4.2)

Therefore, for $\{\lambda_i^k\}_{i\in\mathbb{N},k\in\mathbb{Z}} \subset \mathbb{C}$ and a sequence $\{a_i^k\}_{i\in\mathbb{N},k\in\mathbb{Z}}$ of $(p(\cdot), q, s)$ -atoms supported, respectively, on $\{x_i + B_{\ell_i^k}\}_{i\in\mathbb{N},k\in\mathbb{Z}} \subset \mathfrak{B}$ and

$$g = \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \lambda_i^k a_i^k \in \mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r},$$

from (4.2), we deduce that

$$\begin{aligned} |\mathcal{L}_{\psi}(g)| &= \left| \int_{\mathbb{R}^{n}} \psi(x)g(x) \, dx \right| \leq \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \left| \lambda_{i}^{k} \right| \left| \int_{B} \left| \psi(x) - P(x) \right| \left| a_{i}^{k}(x) \right| \, dx \\ &\leq \sum_{k \in \mathbb{Z}} \sum_{i \in \mathbb{N}} \left| \lambda_{i}^{k} \right| \left\| \psi \right\|_{\mathcal{BMO}_{A}^{p(\cdot), q', s}} \\ &\lesssim \left[\sum_{k \in \mathbb{Z}} 2^{kr} \left\| \sum_{i \in \mathbb{N}} \chi_{x_{i}^{k} + B_{\ell_{i}^{k}}} \right\|_{L^{p(\cdot)}}^{r} \right]^{1/r} \left\| \psi \right\|_{\mathcal{BMO}_{A}^{p(\cdot), q', s}} \\ &\lesssim \left\| g \right\|_{\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r}} \left\| \psi \right\|_{\mathcal{BMO}_{A}^{p(\cdot), q', s}}. \end{aligned}$$

This implies that $\mathcal{BMO}_A^{p(\cdot), q', s} \subset (\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r})^*$.

Next we show that $(\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r})^* \subset \mathcal{BMO}_A^{p(\cdot), q', s}$. For any $B \in \mathfrak{B}$, let

$$S_B: L^1(B) \to P_s$$

be the natural projection satisfying, for any $g \in L^1$ and $Q \in P_s$,

$$\int_{B} S_B(g)(x)Q(x) \, dx = \int_{B} g(x)Q(x) \, dx.$$

By a similar proof of [3, (8.9)], we obtain that, for any $B \in \mathfrak{B}$ and $g \in L^1(B)$,

$$\sup_{x \in B} |S_B(g)(x)| \lesssim \frac{\int_B |g(z)| \, dz}{|B|}$$

Define

$$L_0^q(B) := \{g \in L^q(B) : S_B(g)(x) = 0 \text{ and } g \text{ is not zero almost everywhere} \},\$$

where $L^q(B) := \{ f \in L^q : \operatorname{supp} f \subset B \}$ with $q \in (1, \infty]$ and $B \in \mathfrak{B}$, For any $g \in L^q_0(B)$, set

$$a(x) := \frac{|B|^{1/q}}{\|\chi_B\|_{L^{p(\cdot)}}} \|g\|_{L^q(B)}^{-1} g(x)\chi_B(x)$$

Then a is a $(p(\cdot), q, s)$ -atom. By this and Lemma 4.5, we obtain, for any $\mathcal{L} \in (\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r})^*$ and $g \in L_0^q(B)$,

$$|\mathcal{L}(g)| \leq \frac{\|\chi_B\|_{L^{p(\cdot)}}}{|B|^{1/q}} \|g\|_{L^q(B)} \|\mathcal{L}\|_{(\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}})^*}.$$
(4.3)

Thus, by the Hahn-Banach theorem, it can be extended to a bounded linear functional on $L^q(B)$ with the same norm.

If $q \in (1, \infty]$, by the duality of $L^q(B)$ is $L^{q'}(B)$, we see that there exists a $\Phi \in L^{q'}(B)$ such that, for any $f \in L^q_0(B)$, $\mathcal{L}(f) = \int_B f(x)\Phi(x) dx$. In what follows, for any $B \in \mathfrak{B}$, let $P_s(B)$ denote all the P_s elements vanishing outside B. Now we prove that, if there exists another function $\Phi' \in L^{q'}(B)$ such that, for any $f \in L^q_0(B)$ and $\mathcal{L}(f) = \int_B f(x)\Phi'(x) dx$, then $\Phi' - \Phi \in P_s(B)$. For this, we only need to show that, if $\Phi, \Phi' \in L^1(B)$ such that, for any $f \in L^\infty_0(B)$, $\int_B f(x)\Phi'(x) dx = \int_B f(x)\Phi(x) dx$, then $\Phi - \Phi' \in P_s(B)$. In fact, for any $f \in L^\infty_0(B)$, we have

$$0 = \int_{B} [f(x) - S_{B}(f)(x)] [\Phi'(x) - \Phi(x)] dx$$

= $\int_{B} f(x) [\Phi'(x) - \Phi(x)] dx - \int_{B} f(x) S_{B}(\Phi'(x) - \Phi(x)) dx$
= $\int_{B} f(x) [\Phi'(x) - \Phi(x) - S_{B}(\Phi' - \Phi)(x)] dx.$

Therefore, for a.e. $x \in B \in \mathfrak{B}$, we have

$$\Phi'(x) - \Phi(x) = S_B(\Phi' - \Phi)(x).$$

Hence $\Phi' - \Phi \in P_s(B)$. From this, we see that, for any $q \in (1, \infty]$ and $f \in L_0^q(B)$, there exists a unique $\Phi \in L^{q'}(B)/P_s(B)$ such that $\mathcal{L}(f) = \int_B f(x)\Phi(x) dx$.

For any $j \in \mathbb{N}$ and $g \in L_0^q(B_j)$ with $q \in (1, \infty)$, let $f_j \in L^{q'}(B_j)/P_s(B_j)$ be a unique function such that $\mathcal{L}(g) = \int_{B_j} f_j(x)g(x) dx$. Then, for any $i, j \in \mathbb{N}$ with $i < j, f_j|_{B_i} = f_i$. From this and the fact that, for any $g \in (\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r})^*$, there exists a number $j_0 \in \mathbb{N}$ such that $g \in L_0^q(B_{j_0})$, we conclude that, for any

 $g \in (\mathcal{H}_{A, \operatorname{atom}}^{p(\cdot), q, s, r})^*$, we have

$$\mathcal{L}(g) = \int_{B} \psi(x)g(x) \, dx, \tag{4.4}$$

where $\psi(x) := f_j(x)$ with $x \in B_j$.

Next we show that $\psi \in \mathcal{BMO}_A^{p(\cdot), q', s}$. By [3, (8.12)], (4.3) and (4.4), we have that, for any $q \in (1, \infty)$, $B \in \mathfrak{B}$,

$$\inf_{P \in P_s} \|\psi - P\|_{L^{q'}(B)} = \|\psi\|_{(L^q_0(B))^*} \le \frac{\|\chi_B\|_{L^{p(\cdot)}}}{|B|^{1/q}} \|\mathcal{L}\|_{(\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}})^*}.$$

Therefore, we have that, for any $q \in (1, \infty)$,

$$\begin{aligned} \|\psi\|_{\mathcal{BMO}^{p(\cdot), q', s}} &= \sup_{B \in \mathfrak{B}} \frac{|B|^{1/q}}{\|\chi_B\|_{L^{p(\cdot)}}} \inf_{P \in P_s} \|\psi - P\|_{L^{q'}(B)} = \sup_{B \in \mathfrak{B}} \frac{|B|^{1/q}}{\|\chi_B\|_{L^{p(\cdot)}}} \|\psi\|_{(L^q_0(B))^*} \\ &\leq \|\mathcal{L}\|_{(\mathcal{H}^{p(\cdot), q, s, r}_{A, \operatorname{atom}})^*}, \end{aligned}$$

which implies $\psi \in \mathcal{BMO}_A^{p(\cdot), q', s}$. This finishes the proof of Theorem 4.6.

From Theorem 4.6, we easily obtain the following two conclusions. Moreover, the proof of Corollary 4.8 is similar to [28, Lemma 2.21], we omit the details.

Corollary 4.8 Let A be a given dilation, $p(\cdot) \in C^{\log}$ and $s \in [\lfloor (1/p_- - 1) \ln b / \ln \lambda_- \rfloor, \infty) \cap \mathbb{Z}_+$ with p_- as in (2.4). Assume that $f \in \mathcal{BMO}_A^{p(\cdot), 1, s}$ and $p_+ \in (0, 1]$. Then there exist two positive constants c_1 and c_2 , such that, for any $B \in \mathfrak{B}$ and $\lambda \in (0, \infty)$,

$$|x \in B : |f(x) - P_B^s(f)(x)| > \lambda| \le c_1 \exp\left\{\frac{c_2\lambda|B|}{\|f\|_{\mathcal{BMO}_A^{p(\cdot), 1, s}} \|\chi_B\|_{L^{p(\cdot)}}}\right\}$$

Corollary 4.9 Let A be a given dilation, $p(\cdot) \in C^{\log}$, $p_+ \in (0, 1]$, $q \in (1, \infty)$ and $s \in [\lfloor (1/p_- - 1) \ln b / \ln \lambda_- \rfloor, \infty) \cap \mathbb{Z}_+$ with p_+ , p_- as in (2.4). Then

$$\mathcal{BMO}_A^{p(\cdot),\,1,\,s} = \mathcal{BMO}_A^{p(\cdot),\,q,\,s}$$

with equivalent quasi-norms.

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