

## On the BMO spaces associated with the Laplace-Bessel differential operator

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**Abstract:** In this paper, the characteristic properties of the space of functions of bounded mean oscillation called the  $B$ - $BMO$  associated with the Laplace-Bessel differential operator are obtained. The John-Nirenberg type inequality on the  $B$ - $BMO$  space and a relation between the  $B$ -Poisson integral and the  $B$ - $BMO$  functions are proved.

**Key words:** Laplace-Bessel differential operator, Generalized translation operator,  $BMO$  spaces,  $B$ - $BMO$  spaces, Weighted Lebesgue spaces

### 1. Introduction

The Laplace-Bessel differential operator  $\Delta_B$  which is an important technical tool in Fourier-Bessel harmonic analysis is defined by

$$\Delta_B = \sum_{k=1}^{n-1} \frac{\partial^2}{\partial x_k^2} + \left( \frac{\partial^2}{\partial x_n^2} + \frac{2\nu}{x_n} \frac{\partial^2}{\partial x_n} \right), \quad (\nu > 0, x_n > 0).$$

This operator is a hybrid differential operator which is obtained by applying the Laplace differential operator in the first  $n - 1$  variable and the Bessel differential operator in the last variable.

The relevant Fourier-Bessel harmonic analysis associated with the Bessel differential operator  $B_t$  (or Laplace-Bessel differential operator  $\Delta_B$ ) has been a research area for many mathematicians such as Delsarte, Levitan, Kipriyanov, Klyuchantsev, Lyakhov, Stempak, Gadjiev, Aliev, Guliev, Bayrakci, Sezer, Hasanov and many others [1–5, 8, 11, 14–16, 19–22, 24–26].

The classical  $BMO$  space of functions of bounded mean oscillation, was first introduced by John and Nirenberg [18] in 1961. The space  $BMO$  shares similar properties with the space  $L_\infty$ , and it often serves as a substitute for it. For instance, classical singular integrals do not map  $L_\infty$  to  $L_\infty$  but  $L_\infty$  to  $BMO$ . Many interpolations between  $L_p$  and  $BMO$  work very well between  $L_p$  and  $L_\infty$ . This space has been studied by several authors, e.g., John, Nirenberg, Fefferman, Stein, Garnett, Carleson, Chang, Sadosky, Jones, Meyers, Janson, and others [6, 7, 9, 10, 12, 17, 18, 23].

The space of functions of bounded mean oscillation associated with the Laplace-Bessel differential operator  $\Delta_B$ , called the  $B$ - $BMO$  space was defined by Guliyev [14]. In recent years, the  $B$ - $BMO$  space has been used by many mathematicians such as Guliyev, Abasova, Aliyeva, Shirinova, Hasanov, Ayazoglu, and Bayrakci to obtain the boundedness of some integral operators in suitable function spaces [1, 14–16].

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In this article, we obtain that the characteristic properties provided by the classical BMO space are also valid in the  $B$ -BMO space. We also state the John-Nirenberg-type inequality in B-BMO spaces and prove the relationship between the B-Poisson integrals and the  $B$ -BMO functions.

The paper is organized as follows. Section 2 contains basic definitions and results. The characteristic properties of the  $B$ -BMO space (Theorem 3.3, Theorem 3.5), the relationship between the B-Poisson integral and the  $B$ -BMO function (Theorem 3.6, Theorem 3.7) and finally the John-Nirenberg type inequality and application (Theorem 3.8, Theorem 3.9) are given in Section 3.

**2. Definitions, notations, and preliminaries**

Let  $\mathbb{R}^n$  be the  $n$ -dimensional Euclidean space. Also let for  $x = (x_1, \dots, x_n) \in \mathbb{R}^n$ ,  $|x| = \left(\sum_{k=1}^n x_k^2\right)^{\frac{1}{2}}$ . Denote  $\mathbb{R}_+^n = \{x \in \mathbb{R}^n : x = (x_1, \dots, x_{n-1}, x_n), x_n > 0\}$ . The Lebesgue measure of a measurable set  $E \subset \mathbb{R}_+^n$  is denoted by  $|E|_\nu = \int_E x_n^{2\nu} dx$ ,  $dx = dx_1 dx_2 \dots dx_n$  and  $\nu > 0$  is a fixed parameter.

Suppose that  $E(x, r) = \{y \in \mathbb{R}_+^n : |x - y| < r\}$  denotes the ‘‘ball’’ of radius  $r > 0$  centered at  $x \in \mathbb{R}_+^n$ . It is known that  $|E(0, r)|_\nu = r^{n+2\nu} \omega(n, \nu)$  where  $\omega(n, \nu) = |E(0, 1)|_\nu$ .

Denote by  $T^y$  the generalized translation operator, acting according to the law

$$T^y f(x) = \frac{\Gamma(\nu + \frac{1}{2})}{\Gamma(\nu)\Gamma(\frac{1}{2})} \int_0^\pi f(x' - y'; \sqrt{(x_n^2 - 2x_n y_n \cos \theta + y_n^2)}) \sin^{2\nu-1} \theta d\theta,$$

where  $x = (x', x_n), y = (y', y_n), x', y' \in \mathbb{R}^{n-1}$  and

$$\left(\frac{\Gamma(\nu + \frac{1}{2})}{\Gamma(\nu)\Gamma(\frac{1}{2})}\right)^{-1} = \int_0^\pi \sin^{2\nu-1} \theta d\theta.$$

Note that the generalized translation operator  $T^y$  is closely related with the Laplace-Bessel differential operator  $\Delta_B$  (see [8, 21, 22] for details).

Let  $L_{p,\nu}(\mathbb{R}_+^n), 1 \leq p < \infty$  be the space of all measurable functions on  $\mathbb{R}_+^n$  with the norm

$$\|f\|_{p,\nu} = \left(\int_{\mathbb{R}_+^n} |f(x)|^p x_n^{2\nu} dx\right)^{\frac{1}{p}} < \infty.$$

In the case of  $p = \infty$ , the space  $L_\infty(\mathbb{R}_+^n)$  is equipped with the norm  $\|f\|_\infty = \text{ess sup}_{x \in \mathbb{R}_+^n} |f(x)|$ . Also, we denote by  $L_1^{loc}(\mathbb{R}_+^n)$  the set of locally integrable functions on  $\mathbb{R}_+^n$ .

**Lemma 2.1** [16] *For all  $x \in \mathbb{R}_+^n$  the following equality is valid:*

$$\int_{E(0,r)} T^z f(x) z_n^{2\nu} dz = \frac{\Gamma(\nu + \frac{1}{2})}{\Gamma(\nu)\Gamma(\frac{1}{2})} \int_{E(\tilde{x},r)} f(y', \sqrt{y_n^2 + y_{n+1}^2}) d\tilde{y}$$

where  $E(\tilde{x}, r) = E((x, 0), r) = \{\tilde{y} = (y, y_{n+1}) \in \mathbb{R}^n \times (0, \infty) : |\tilde{x} - \tilde{y}| < r\}$ ,  $y' \in \mathbb{R}^{n-1}$ ,  $dy = dy_1 \dots dy_{n-1} dy_n$ .

**Proof** Let  $\tilde{x} = (x, 0) = (x_1, \dots, x_n, 0)$ ,  $\tilde{y} = (y_1, \dots, y_n, y_{n+1})$ . We set  $\tilde{w} = \tilde{x} - \tilde{y} = (x_1 - y_1, \dots, x_n - y_n, -y_{n+1})$ . Thus we have

$$\int_{E(\tilde{x}, r)} f\left(y', \sqrt{y_n^2 + y_{n+1}^2}\right) y_{n+1}^{2\nu-1} dy_{n+1} dy = \int_{E(0, r)} f\left(x' - w', \sqrt{(x_n - w_n)^2 + w_{n+1}^2}\right) w_{n+1}^{2\nu-1} dw_{n+1} dw.$$

Now, let us apply the substitutions for  $(z_1, \dots, z_n) \in \mathbb{R}_+^n$ :  $w' = z'$ ,  $w_n = z_n \cos \theta$ ,  $w_{n+1} = z_n \sin \theta$ ,  $0 \leq \theta < \pi$ ,  $z_n \geq 0$ . Finally, we get

$$\frac{\Gamma(\nu + \frac{1}{2})}{\Gamma(\nu)\Gamma(\frac{1}{2})} \int_{E(0, r)} \int_0^\pi f\left(x' - z', \sqrt{(x_n - z_n \cos \theta)^2 + z_n^2 \sin^2 \theta}\right) z_n^{2\nu-1} \sin^{2\nu-1} \theta z_n d\theta dz = \int_{E(0, r)} T^z f(x) z_n^{2\nu} dz.$$

□

### 3. Main definitions and results

**Definition 3.1** (Guliev [14]) *The B-BMO space, generated by the generalized translation operator is defined as the space of locally integrable functions f with the norm (see Remark below):*

$$\|f\|_{B-BMO} = \sup_{\substack{r>0 \\ x \in \mathbb{R}_+^n}} \frac{1}{|E(0, r)|_\nu} \int_{E(0, r)} |T^y f(x) - f_{E(0, r)}(x)| y_n^{2\nu} dy < \infty$$

where

$$f_{E(0, r)}(x) = \frac{1}{|E(0, r)|_\nu} \int_{E(0, r)} T^y f(x) y_n^{2\nu} dy.$$

It is a simple fact that B-BMO is a linear space, that is, if  $f, g \in B-BMO$  and  $\alpha \in \mathbb{C}$  then,  $f + g$  and  $\alpha f$  are also in B-BMO and

$$\begin{aligned} \|f + g\|_{B-BMO} &\leq \|f\|_{B-BMO} + \|g\|_{B-BMO}, \\ \|\alpha f\|_{B-BMO} &= |\alpha| \|f\|_{B-BMO}. \end{aligned}$$

**Remark 3.2**  $\|\cdot\|_{B-BMO}$  is not a norm. The problem is that if  $\|f\|_{B-BMO} = 0$ , this does not imply that  $f = 0$  but that  $f$  is a constant. Moreover, every constant function  $c$  satisfies  $\|c\|_{B-BMO} = 0$ . Although  $\|\cdot\|_{B-BMO}$  is only a seminorm, it can be taken as a norm when there is no possibility of confusion.

Now, let us begin with a list of basic properties of the space B-BMO.

**Theorem 3.3** *Let  $f \in L_1^{loc}(\mathbb{R}_+^n)$ . The following properties of the space B-BMO are valid:*

- a)  $L_\infty(\mathbb{R}_+^n) \not\subseteq B-BMO$  and  $\|f\|_{B-BMO} \leq 2 \|f\|_\infty$ .
- b) Suppose that there exists an  $A > 0$  such that for all  $x \in \mathbb{R}_+^n$  and all  $r > 0$  there exists a constant  $c_{x,r}$  such that

$$\sup_{\substack{r>0 \\ x \in \mathbb{R}_+^n}} \frac{1}{|E(0, r)|_\nu} \int_{E(0, r)} |T^y f(x) - c_{x,r}| y_n^{2\nu} dy \leq A. \tag{3.1}$$

Then,  $f \in B\text{-BMO}$  and  $\|f\|_{B\text{-BMO}} \leq 2A$ .

c) If  $f \in B\text{-BMO}$  and  $\lambda > 0$  then, the function  $\delta^\lambda f$  defined by  $\delta^\lambda f(x) = f(\lambda x)$  is also in  $B\text{-BMO}$  and

$$\|\delta^\lambda f\|_{B\text{-BMO}} = \|f\|_{B\text{-BMO}}. \tag{3.2}$$

**Proof** a) Let  $f \in L_\infty(\mathbb{R}_+^n)$ . Since

$$\begin{aligned} |T^y f(x)| &\leq \frac{\Gamma(\nu + \frac{1}{2})}{\Gamma(\nu)\Gamma(\frac{1}{2})} \int_0^\pi \left| f\left(x' - y'; \sqrt{x_n^2 - 2x_n y_n \cos \alpha + y_n^2}\right) \right| \sin^{2\nu-1} \alpha d\alpha \\ &\leq \operatorname{ess\,sup}_{x \in \mathbb{R}_+^n} |f(x)| \frac{\Gamma(\nu + \frac{1}{2})}{\Gamma(\nu)\Gamma(\frac{1}{2})} \int_0^\pi \sin^{2\nu-1} \alpha d\alpha = \|f\|_\infty \end{aligned}$$

we have

$$|f_{E(0,r)}(x)| \leq \frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} |T^y f(x)| y_n^{2\nu} dy \leq \|f\|_\infty.$$

Hence,

$$\begin{aligned} &\frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} |T^y f(x) - f_{E(0,r)}(x)| y_n^{2\nu} dy \leq \\ &\leq \frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} |T^y f(x)| y_n^{2\nu} dy + \frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} |f_{E(0,r)}(x)| y_n^{2\nu} dy \\ &\leq \|f\|_\infty + \|f\|_\infty = 2\|f\|_\infty. \end{aligned}$$

Thus,  $\|f\|_{B\text{-BMO}} \leq 2\|f\|_\infty$  is obtained by taking the supremum over all  $x \in \mathbb{R}_+^n$ ,  $r > 0$ .

b) Firstly, since

$$|T^y f(x) - f_{E(0,r)}(x)| \leq |T^y f(x) - c_{x,r}| + |c_{x,r} - f_{E(0,r)}(x)|$$

and

$$\begin{aligned} |c_{x,r} - f_{E(0,r)}(x)| &= \left| c_{x,r} - \frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} T^y f(x) y_n^{2\nu} dy \right| \\ &\leq \frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} |T^y f(x) - c_{x,r}| y_n^{2\nu} dy \end{aligned}$$

we have

$$\frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} |T^y f(x) - f_{E(0,r)}(x)| y_n^{2\nu} dy \leq \frac{2}{|E(0,r)|_\nu} \int_{E(0,r)} |T^y f(x) - c_{x,r}| y_n^{2\nu} dy.$$

Finally, by taking supremum over all  $x \in \mathbb{R}_+^n$ ,  $r > 0$  we have  $f \in B\text{-}BMO$  and  $\|f\|_{B\text{-}BMO} \leq 2A$ .

c) For the property (3.2), let us show  $(\delta^\lambda f)_{E(0,r)}(x) = f_{E(0,\lambda r)}(\lambda x)$ . For this, we have

$$\begin{aligned} T^y(\delta^\lambda f)(x) &= \frac{\Gamma(\nu + \frac{1}{2})}{\Gamma(\nu)\Gamma(\frac{1}{2})} \int_0^\pi (\delta^\lambda f)\left(x' - y', \sqrt{x_n^2 - 2x_n y_n \cos \alpha + y_n^2}\right) \sin^{2\nu-1} \alpha d\alpha \\ &= \frac{\Gamma(\nu + \frac{1}{2})}{\Gamma(\nu)\Gamma(\frac{1}{2})} \int_0^\pi f\left(\lambda x' - \lambda y', \sqrt{(\lambda x_n)^2 - 2(\lambda x_n)(\lambda y_n) \cos \alpha + (\lambda y_n)^2}\right) \sin^{2\nu-1} \alpha d\alpha \\ &= T^{\lambda y} f(\lambda x). \end{aligned} \tag{3.3}$$

Thus we get

$$\begin{aligned} (\delta^\lambda f)_{E(0,r)}(x) &= \frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} T^y(\delta^\lambda f)(x) y_n^{2\nu} dy \\ &= \frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} T^{\lambda y} f(\lambda x) y_n^{2\nu} dy \\ &= \dots(z = \lambda y, dz = \lambda^n dy, z_n^{2\nu} = \lambda^{2\nu} y_n^{2\nu}) \dots \\ &= \frac{1}{|E(0,r)|_\nu} \frac{1}{\lambda^{n+2\nu}} \int_{\lambda E(0,r)} T^z f(\lambda x) z_n^{2\nu} dz \\ &= (f)_{\lambda E(0,r)}(\lambda x) = (f)_{E(0,\lambda r)}(\lambda x). \end{aligned} \tag{3.4}$$

Hence by taking into account (3.3) and (3.4), we obtain  $\|\delta^\lambda f\|_{B\text{-}BMO} = \|f\|_{B\text{-}BMO}$ . □

**Remark 3.4** We indicate that  $L_\infty(\mathbb{R}_+^n)$  is a proper subspace of  $B\text{-}BMO$ . As in case of the classical  $BMO$ -space, we claim that the function  $f(x) = \log|x|$  is in  $B\text{-}BMO$  but not in  $L_\infty(\mathbb{R}_+^n)$ , (cf. ([25], p.141)).

Our next goal is to determine the connection between the  $B\text{-}BMO$  function and  $B$ -Poisson integral and present John-Nirenberg-type inequality for  $B\text{-}BMO$  functions. These results have promising applications for the classical  $BMO$ -functions, such as determining the characterization of  $BMO$ -space and the connection between  $BMO$ -functions and the Carleson measure (see [7, 9, 10, 13, 25] for details).

**Theorem 3.5** Let  $f$  be in  $B\text{-}BMO$ . Then,

a)

$$|f_{E(0,r)}(x) - f_{E(0,s)}(x)| \leq \left(\frac{s}{r}\right)^{n+2\nu} \|f\|_{B\text{-}BMO}, \quad r \leq s. \tag{3.5}$$

b)

$$|f_{E(0,r)}(x) - f_{E(0,2^m r)}(x)| \leq 2^{n+2\nu} m \|f\|_{B\text{-}BMO}, \quad m \in \mathbb{N}. \tag{3.6}$$

**Proof** a) For  $r \leq s$ , we have

$$\begin{aligned} |f_{E(0,r)}(x) - f_{E(0,s)}(x)| &\leq \frac{1}{|E(0,r)|_\nu} \int_{E(0,r)} |T^y f(x) - f_{E(0,s)}(x)| y_n^{2\nu} dy \\ &\leq \frac{|E(0,s)|_\nu}{|E(0,r)|_\nu} \frac{1}{|E(0,s)|_\nu} \int_{E(0,s)} (T^y f(x) - f_{E(0,s)}(x)) y_n^{2\nu} dy \\ &\leq \left(\frac{s}{r}\right)^{n+2\nu} \|f\|_{B-BMO}. \end{aligned}$$

b) Firstly let  $m = 1$ . By using (3.5) we get

$$|f_{E(0,r)}(x) - f_{E(0,2r)}(x)| \leq \left(\frac{2r}{r}\right)^{n+2\nu} \|f\|_{B-BMO} = 2^{n+2\nu} \|f\|_{B-BMO}.$$

Finally using this inequality, we obtain

$$\begin{aligned} |f_{E(0,r)}(x) - f_{E(0,2^m r)}(x)| &\leq |f_{E(0,r)}(x) - f_{E(0,2r)}(x)| + \dots + |f_{E(0,2^{m-1}r)}(x) - f_{E(0,2^m r)}(x)| \\ &\leq 2^{n+2\nu} m \|f\|_{B-BMO}, \quad m \in \mathbb{N}. \end{aligned}$$

□

**Theorem 3.6** *There exists a constant  $c(n, \nu) > 0$  such that for all  $f \in B-BMO$  we have*

$$\sup_{\substack{x \in \mathbb{R}_+^n \\ t > 0}} \int_{\mathbb{R}_+^n} |T^y f(x) - V_t f(x)| P_\nu(y, t) y_n^{2\nu} dy \leq c(n, \nu) \|f\|_{B-BMO}. \tag{3.7}$$

Here  $P_\nu(y, t)$  denotes the B-Poisson kernel introduced in [2] and

$$(V_t f)(x) = \int_{\mathbb{R}_+^n} T^y f(x) P_\nu(y, t) y_n^{2\nu} dy$$

the B-Poisson integral of  $f$ .

**Proof** We have

$$\begin{aligned} &|T^y f(x) - V_t f(x)| P_\nu(y, t) y_n^{2\nu} dy \leq \\ &\leq \int_{\mathbb{R}_+^n} |T^y f(x) - f_{E(0,t)}(x)| P_\nu(y, t) y_n^{2\nu} dy + \int_{\mathbb{R}_+^n} |V_t f(x) - f_{E(0,t)}(x)| P_\nu(y, t) y_n^{2\nu} dy \\ &= I_1 + I_2. \end{aligned} \tag{3.8}$$

Now, let us calculate the estimates  $I_1$  and  $I_2$ . Since

$$P_\nu(y, t) = d_\nu(n) \frac{t}{(t^2 + |y|^2)^{(n+2\nu+1)/2}}$$

and

$$d_\nu(n) = \frac{(2\pi)^{1-n} 2^{1-\nu+n/2} \Gamma(n+2\nu+1/2)}{\sqrt{\pi} \Gamma^2(n+1/2)}, \quad (\text{see [2] for details})$$

we obtain

$$\begin{aligned} I_1 &= \int_{\mathbb{R}_+^n} |T^y f(x) - f_{E(0,t)}(x)| P_\nu(y,t) y_n^{2\nu} dy = d_\nu(n) \int_{\mathbb{R}_+^n} \frac{t |T^y f(x) - f_{E(0,t)}(x)|}{(t^2 + |y|^2)^{\frac{n+2\nu+1}{2}}} y_n^{2\nu} dy \\ &\leq \int_{E(0,t)} \frac{t |T^y f(x) - f_{E(0,t)}(x)|}{(t^2 + |y|^2)^{\frac{n+2\nu+1}{2}}} y_n^{2\nu} dy + \\ &\quad + \sum_{k=0}^\infty \int_{E(0,2^{k+1}t) \setminus E(0,2^k t)} \frac{t (|T^y f(x) - f_{E(0,2^{k+1}t)}(x)| + |f_{E(0,2^{k+1}t)} - f_{E(0,t)}(x)|)}{(t^2 + |y|^2)^{(n+2\nu+1)/2}} y_n^{2\nu} dy \\ &\leq \frac{1}{t^{n+2\nu}} \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)| y_n^{2\nu} dy + \sum_{k=0}^\infty 2^{-k(n+2\nu+1)} \times \\ &\quad \left( \frac{1}{t^{n+2\nu}} \int_{2^{k+1}E(0,t)} |T^y f(x) - f_{2^{k+1}E(0,t)}(x)| y_n^{2\nu} dy + \frac{1}{t^{n+2\nu}} |f_{2^{k+1}E(0,t)}(x) - f_{E(0,t)}(x)| 2^{(k+1)(n+2\nu)} \omega(n,\nu) t^{n+2\nu} \right) \\ &\leq \omega(n,\nu) \|f\|_{B-BMO} + \omega(n,\nu) \|f\|_{B-BMO} \left( 2^{n+2\nu} \sum_{k=0}^\infty \frac{1}{2^k} + 2^{2(n+2\nu)} \sum_{k=0}^\infty \frac{k+1}{2^k} \right) \\ &\leq c(n,\nu) \|f\|_{B-BMO}. \end{aligned} \tag{3.9}$$

Also, by taking into account

$$\int_{\mathbb{R}_+^n} P_\nu(y,t) y_n^{2\nu} dy = 1, \quad (\text{see [2]})$$

we obtain

$$\begin{aligned} I_2 &= \int_{\mathbb{R}_+^n} |V_t f(x) - f_{E(0,t)}(x)| P_\nu(y,t) y_n^{2\nu} dy \\ &= |V_t f(x) - f_{E(0,t)}(x)| \\ &\leq \int_{\mathbb{R}_+^n} |T^y f(x) - f_{E(0,t)}(x)| P_\nu(y,t) y_n^{2\nu} dy \\ &\leq c(n,\nu) \|f\|_{B-BMO} \end{aligned}$$

which combined with the inequality in (3.9). Therefore, this estimate combined with (3.8) yields the desired result and finishes the proof.

The following theorem gives a sufficient condition for the inverse inequality in (3.7). □

**Theorem 3.7** *There is a constant  $\tilde{c}(n, \nu)$  such that for all  $f \in L_1^{loc}(\mathbb{R}_+^n)$  for which*

$$\int_{\mathbb{R}_+^n} \frac{|T^y f(x)|}{(1 + |y|^2)^{n+2\nu+1}} y_n^{2\nu} dy < \infty$$

*we have*

$$\sup_{\substack{x \in \mathbb{R}_+^n \\ t > 0}} \int_{\mathbb{R}_+^n} |T^y f(x) - V_t f(x)| P_\nu(y, t) y_n^{2\nu} dy \geq \tilde{c}(n, \nu) \|f\|_{B-BMO}.$$

**Proof** Let

$$A = \sup_{\substack{x \in \mathbb{R}_+^n \\ t > 0}} \int_{\mathbb{R}_+^n} |T^y f(x) - V_t f(x)| P_\nu(y, t) y_n^{2\nu} dy.$$

For  $y \in E(0, t)$ ,  $|y| < t$  we have  $(t^2 + |y|^2)^{(n+2\nu+1)/2} \leq c_1(n, \nu) t^{n+2\nu+1}$  and  $P_\nu(y, t) \geq c_2(n, \nu) t^{-(n+2\nu)}$  which gives

$$A \geq \frac{c_2(n, \nu)}{t^{n+2\nu}} \int_{\mathbb{R}_+^n} |T^y f(x) - V_t f(x)| y_n^{2\nu} dy.$$

Taking into account (3.1) implies that

$$\|f\|_{B-BMO} \leq \frac{2A}{c_2(n, \nu)}.$$

This concludes the theorem. □

**Theorem 3.8** *For all  $f \in B-BMO$ , all  $t > 0$ ,  $x \in \mathbb{R}_+^n$  and all  $\alpha > 0$  there exist  $c_1 = c_1(n, \nu)$ ,  $c_2 = c_2(n, \nu) > 0$  such that*

$$|\{x \in E(0, t) : |T^y f(x) - f_{E(0,t)}(x)| > \alpha\}|_\nu \leq c_1 e^{-\alpha c_2 / \|f\|_{B-BMO}}. \tag{3.10}$$

**Proof** For the sake of simplicity, we assume that  $\|f\|_{B-BMO} = 1$ . We apply the Calderon-Zygmund decomposition to the function  $T^y f - f_{E(0,t)}$  inside the  $E(0, t)$ . The rest part of the proof is similar to the proof of the well-known John-Nirenberg inequality for the classical  $BMO$  functions, (see [18], (cf. [13], p.124) for details). □

By using Theorem 3.8, we obtain the following important  $L_{p,\nu}$  characterization of  $B-BMO$ .

**Theorem 3.9** *Let  $f$  be in  $B-BMO$ . Then, for  $1 \leq p < \infty$  there exists  $c = c(n, \nu, p)$ ,  $d = d(n, \nu, p) > 0$  such that*

$$c \|f\|_{B-BMO} \leq \sup_{\substack{x \in \mathbb{R}_+^n \\ t > 0}} \left( \frac{1}{|E(0, t)|_\nu} \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)|^p y_n^{2\nu} dy \right)^{1/p} \leq d \|f\|_{B-BMO}. \tag{3.11}$$



**Proof** By using (3.10) and the Gamma function  $\Gamma(z) = \int_0^\infty t^{z-1} e^{-t} dt$  we have

$$\begin{aligned}
 & \frac{1}{|E(0,t)|_\nu} \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)|^p y_n^{2\nu} dy = \\
 &= \frac{1}{|E(0,t)|_\nu} \int_0^\infty p\alpha^{p-1} |\{x \in E(0,t) : |T^y f(x) - f_{E(0,t)}(x)|^p > \alpha\}|_\nu d\alpha \\
 &\leq \frac{pc_1}{|E(0,t)|_\nu} \int_0^\infty \alpha^{p-1} e^{-\alpha c_2/\|f\|_{B-BMO}} d\alpha \\
 &= \dots(\beta = \frac{\alpha c_2}{\|f\|_{B-BMO}}, d\beta = \frac{c_2}{\|f\|_{B-BMO}} d\alpha) \dots \\
 &= c_3 \|f\|_{B-BMO}^{p-1} \|f\|_{B-BMO} \int_0^\infty \beta^{p-1} e^{-\beta} d\beta \\
 &= d^p \|f\|_{B-BMO}^p, \quad d = d(n, \nu, p) > 0.
 \end{aligned} \tag{3.12}$$

Hence,

$$\sup_{\substack{x \in \mathbb{R}_+^n \\ t > 0}} \left( \frac{1}{|E(0,t)|_\nu} \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)|^p y_n^{2\nu} dy \right)^{\frac{1}{p}} \leq d \|f\|_{B-BMO}.$$

Also by taking into account the Hölder inequality  $\frac{1}{p} + \frac{1}{q} = 1, 1 < p < \infty$  we obtain

$$\begin{aligned}
 \int_{E(0,t)} |T^y f(x) - f_{E_t}(x)| y_n^{2\nu} dy &\leq \left( \int_{E(0,t)} |T^y f(x) - f_{E_t}(x)|^p y_n^{2\nu} dy \right)^{\frac{1}{p}} \left( \int_{E(0,t)} y_n^{2\nu} dy \right)^{\frac{1}{q}} \\
 &= \left( \int_{E_t} |T^y f(x) - f_{E_t}(x)|^p y_n^{2\nu} dy \right)^{\frac{1}{p}} |E(0,t)|_\nu^{\frac{1}{q}}.
 \end{aligned}$$

Therefore,

$$\begin{aligned}
 \frac{1}{|E(0,t)|_\nu} \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)| y_n^{2\nu} dy &\leq \frac{|E(0,t)|_\nu^{\frac{1}{q}}}{|E(0,t)|_\nu} \left( \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)|^p y_n^{2\nu} dy \right)^{\frac{1}{p}} \\
 &= \left( \frac{1}{|E(0,t)|_\nu} \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)|^p y_n^{2\nu} dy \right)^{\frac{1}{p}}
 \end{aligned}$$

and

$$\|f\|_{B-BMO} \leq \sup_{\substack{x \in \mathbb{R}_+^n \\ t > 0}} \left( \frac{1}{|E(0,t)|_\nu} \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)|^p y_n^{2\nu} dy \right)^{\frac{1}{p}}, \quad 1 \leq p < \infty. \quad (3.13)$$

Finally, combining (3.12) and (3.13) yields that (3.11).  $\square$

**Corollary 3.10** *For all  $1 \leq p < \infty$  we have*

$$\|f\|_{B-BMO} \approx \sup_{\substack{x \in \mathbb{R}_+^n \\ t > 0}} \left( \frac{1}{|E(0,t)|_\nu} \int_{E(0,t)} |T^y f(x) - f_{E(0,t)}(x)|^p y_n^{2\nu} dy \right)^{\frac{1}{p}}.$$

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