

#### **Turkish Journal of Mathematics**

http://journals.tubitak.gov.tr/math/

Research Article

Turk J Math (2022) 46: 377 – 386 © TÜBİTAK doi:10.3906/mat-2104-69

# Determination of a differential pencil from interior spectral data on a union of two closed intervals

## İbrahim ADALAR\*

Zara Veysel Dursun School of Applied Sciences, Sivas Cumhuriyet University, Sivas, Turkey

Received: 15.04.2021 • Accepted/Published Online: 02.06.2021 • Final Version: 21.01.2022

**Abstract:** In this paper, we consider a quadratic pencil of Sturm–Liouville operator on closed sets. We study an interior-inverse problem for this kind operator and give a uniqueness theorem with an appropriate example.

Key words: Inverse problem, differential pencil, time scale

#### 1. Introduction

The spectral problems appear in geophysics, mathematical physics, mechanics and other branches of natural sciences. Specifically, second order differential pencil, namely diffusion equation, originates from the problem of describing interactions between colliding particles in physics [17]. The reconstruction of an operator from its spectral characteristics is referred to as the inverse problem. For inverse problems of a differential pencil, such characteristics are one spectrum and normalizing constants, two given spectra, interior spectral data, nodal points, scattering data and the Weyl function.

The inverse problem for interior spectral data of the differential operator consists of reconstruction of this operator from the given eigenvalues and some information on eigenfunctions at an internal point. The interior-inverse Sturm-Liouville problem, which is one of the important subjects of inverse spectral theory, was first studied in by Mochizuki and Trooshin in [21]. They proved that the spectrum of the problem

$$-y'' + q(x)y = \lambda y, \ t \in (0,1)$$
$$y'(0) - hy(0) = y'(1) + Hy(1) = 0$$

and the logarithmic derivatives of the eigenfunctions at the point 1/2 uniquely determine the potential q(x) on the whole interval [0,1] almost everywhere. This kind of problems for the differential operators on a continuous interval were studied in [8, 14], [18, 22, 23], [28, 34].

The time scale theory unifies continuous and discrete dynamic equations. This approach is used in many important applied problems, for example, heat transfer, dynamics of population, string theory. The theory of dynamic equations on closed sets was introduced in [16]. Nowadays, there are a few studies about the inverse problems of differential operators defined on time scales (see [1, 2, 7, 19, 20],[24, 27] and [35]).

If  $\mathbb{T}$  is a closed subset of  $\mathbb{R}$  it called a time scale. The forward and backward jump operators  $\sigma$ ,  $\rho: \mathbb{T} \to \mathbb{T}$ 

 $2010\ AMS\ Mathematics\ Subject\ Classification:\ 31B20,\ 39A12,\ 34N05$ 

<sup>\*</sup>Correspondence: iadalar@cumhuriyet.edu.tr

are defined as follows:

$$\sigma(t) = \inf \{ s \in \mathbb{T} : s > t \} \text{ if } t \neq \sup \mathbb{T},$$
  
$$\rho(t) = \sup \{ s \in \mathbb{T} : s < t \} \text{ if } t \neq \inf \mathbb{T}.$$

The delta derivative  $f^{\Delta}$  and the nabla derivative  $f^{\nabla}$  for a function f defined on  $\mathbb{T}$  was introduced in [4]. For basic concepts of the time scale theory and time scales notation, we refer to the textbooks [4, 5] and [6].

In this paper, we consider an interior-inverse problem for a differential pencil on a time scale and give a uniqueness theorem. The main result of this article is a generalization of the classical result for an interior-inverse problem on a continuous interval.

Throughout this paper, we assume that  $\mathbb{T} = [0, a_1] \cup [a_2, a_3]$  is a bounded time scale for  $0 < a_1 \le a_2$  and  $a_1 + a_2 = a_3$ .

We consider following boundary value problem L = L(q, p, h, H) on  $\mathbb{T} = [0, a_1] \cup [a_2, a_3]$ ,

$$\ell y := -y^{\Delta \Delta}(t) + [q(t) + 2\lambda p(t)]y^{\sigma}(t) = \lambda^2 y^{\sigma}(t), \quad t \in \mathbb{T}, \tag{1.1}$$

$$U(y) := y^{\Delta}(0) - hy(0) = 0, \tag{1.2}$$

$$V(y) := y^{\Delta}(a_3) + Hy(a_3) = 0 \tag{1.3}$$

where q(t) and p(t) are real valued continuous functions on  $\mathbb{T}$ ;  $h, H \in \mathbb{R}$  and  $\lambda$  is a spectral parameter.

Together with L, we consider a boundary value problem  $\widetilde{L} = L(\widetilde{q}, \widetilde{p}, h, H)$  of the same form but with different coefficients  $\widetilde{q}(t)$  and  $\widetilde{p}(t)$ . We assume that if a certain symbol x denotes an object related to L, then  $\widetilde{x}$  will denote an analogous object related to  $\widetilde{L}$ . Also, we assume

$$h|y(0)|^2 + H|y(a_3)|^2 + \int_0^{a_3} (|y^{\Delta}(t)|^2 + q(t)|y^{\sigma}(t)|^2) \Delta t > 0.$$

A function y, defined on  $\mathbb{T}$ , is called a solution of equation (1) if  $y \in C^2_{rd}(\mathbb{T})$  and y satisfies (1) for all  $t \in \mathbb{T}$ . The values of the  $\lambda$  parameter—for which (1)–(3) has nonzero solutions are called eigenvalues, and the corresponding nontrivial solutions are called eigenfunctions.

We state the main result of this paper.

Let  $\Lambda := \{\lambda_n, n \in \mathbb{Z}\}$  and  $\widetilde{\Lambda} := \{\widetilde{\lambda}_n, n \in \mathbb{Z}\}$  be the eigenvalues sets of L and  $\widetilde{L}$ ,  $y_n(t)$  and  $\widetilde{y_n}(t)$  are eigenfunctions related to this eigenvalues, respectively.

**Theorem 1.1** If  $\Lambda = \widetilde{\Lambda}$ , and for any  $n \in \mathbb{N}$ ,

$$\frac{y_n^{\nabla}(a_2)}{y_n(a_2)} = \frac{\widetilde{y}_n^{\nabla}(a_2)}{\widetilde{y}_n(a_2)} \tag{1.4}$$

then  $q(t) = \widetilde{q}(t)$  and  $p(t) = \widetilde{p}(t)$  on  $\mathbb{T}$ .

In case  $a_1 = a_2$  and p(t) = 0, Theorem 1 coincides with the theorem in [21]. We note that, if it is taken  $a_1 = a_2$  and  $\mathbb{T} = [0, a_3]$  in L, the classical diffusion operator is obtained. In addition, since  $y_n^{\Delta}(a_1) = -y_n^{\nabla}(a_2)$ ,

the condition (1.4) can be replaced by

$$\frac{y_n^{\Delta}(a_1)}{y_n(a_1)} = \frac{\widetilde{y}_n^{\Delta}(a_1)}{\widetilde{y}_n(a_1)}.$$

#### 2. Preliminaries

Let  $\varphi(t,\lambda)$  and  $\psi(t,\lambda)$  be the solution of (1.1) under the initial conditions

$$\varphi(0,\lambda) = 1, \varphi^{\Delta}(0,\lambda) = h$$

$$\psi(a_3,\lambda) = 1, \psi^{\Delta}(a_3,\lambda) = -H$$

respectively. The following lemmas can be given by using a similar method to that in [24].

**Lemma 2.1**  $\varphi(t,\lambda), \ \psi(t,\lambda)$  and their  $\Delta$ -derivatives are entire functions on  $\lambda$  for each fixed t.

**Lemma 2.2** The eigenvalues of the problem (1.1)-(1.3) are real and algebraicly simple.

It is clear that the characteristic function of L can be given as follows

$$\Delta(\lambda) = W[\psi, \varphi] = \varphi^{\Delta}(a_3, \lambda) + H\varphi(a_3, \lambda).$$

It follows from Lemma 1 that  $\Delta(\lambda)$  is also an entire function and so L has a discrete spectrum. It is known from [15] that  $\varphi(t,\lambda)$  satisfies the following representation on  $[0,a_1]$ 

$$\varphi(t,\lambda) = \cos(\lambda t - \alpha(t)) + \int_{0}^{t} A(t,x)\cos\lambda x dx + \int_{0}^{t} B(t,x)\sin\lambda x dx$$

where the kernels A(t,x) and B(t,x) are the solution of the problem

$$\frac{\partial^2 A(t,x)}{\partial t^2} - 2p(t)\frac{\partial B(t,x)}{\partial x} - q(t)A(t,x) = \frac{\partial^2 A(t,x)}{\partial x^2},$$

$$\frac{\partial^2 B(t,x)}{\partial t^2} + 2p(t)\frac{\partial A(t,x)}{\partial x} - q(t)B(t,x) = \frac{\partial^2 B(t,x)}{\partial x^2},$$

$$A(0,0) = h, \ B(t,0) = 0, \ \frac{\partial A(t,x)}{\partial x}\Big|_{x=0} = 0$$

where

$$q(t) + p^2(t) = 2\frac{d}{dt} \left[ A(t,t)\cos\alpha(t) + B(t,t)\sin\alpha(t) \right], \ \alpha(t) = \int_0^t p(x)dx.$$

On the other hand,  $\varphi^{\Delta}(t,\lambda)$  is continuous at  $a_1$ , and so the relation

$$a\varphi^{\Delta}(a_1 - 0, \lambda) = \varphi(a_2, \lambda) - \varphi(a_1, \lambda)$$
(2.1)

holds, where  $a := a_2 - a_1$ . Therefore, we obtain that if a > 0,, then the following asymptotic formula holds for  $|\lambda| \to \infty$ ;

$$\varphi(a_2, \lambda) = -a\lambda \sin(\lambda a_1 - \alpha(a_1)) + O(\exp|\tau| a_1), \qquad (2.2)$$

where  $\tau = Im\lambda$ .

Since  $\varphi(t,\lambda)$  satisfies the equation (1.1) for  $t=a_1$ , it follows that

$$(a^2\lambda^2 + b\lambda + c)\varphi(a_2, \lambda) + a\varphi^{\Delta}(a_2 + 0, \lambda) + \varphi(a_1, \lambda) = 0,$$
(2.3)

where  $a = a_2 - a_1$ ,  $b = -2a^2p(a_1)$ ,  $c = -a^2q(a_1) - 1$ .

**Lemma 2.3** The following asymptotic formulas hold for  $|\lambda| \to \infty$ :

$$\varphi(t,\lambda) = \begin{cases} \cos \lambda t + \frac{h}{\lambda} \sin \lambda t + O\left(\frac{1}{\lambda} \exp|\tau| t\right), & t \in [0, a_1], \\ a^2 \lambda^2 \sin \lambda a_1 \sin \lambda (t - a_2) + O\left(\lambda \exp|\tau| (t - a_2 + a_1)\right), & t \in [a_2, a_3], \end{cases}$$
(2.4)

$$\varphi^{\Delta}(t,\lambda) = \begin{cases} -\lambda \sin \lambda t + h \cos \lambda t + O(\exp |\tau| t), & t \in [0, a_1), \\ a^2 \lambda^3 \sin \lambda a_1 \cos \lambda (t - a_2) + O(\lambda^2 \exp |\tau| (t - a_2 + a_1)), & t \in [a_2, a_3]. \end{cases}$$
(2.5)

**Proof** It is clear that

$$\varphi(t,\lambda) = \cos \lambda t + \frac{h}{\lambda} \sin \lambda t +$$

$$+ \frac{1}{\lambda} \int_{0}^{t} \sin \lambda (t - \xi) \left[ q(\xi) + 2\lambda p(\xi) \right] \varphi(\xi,\lambda) \Delta \xi, \ t \in [0, a_{1}]$$
(2.6)

On the other hand, the general solution of (1.1) on  $[a_2, a_3]$  is given as follow.

$$\varphi(t,\lambda) = c_1(\lambda)\cos\lambda t + c_2(\lambda)\sin\lambda t + \frac{1}{\lambda} \int_{a_2}^{t} \sin\lambda (t-\xi) \left[q(\xi) + 2\lambda p(\xi)\right] \varphi(\xi,\lambda) \Delta \xi.$$
(2.7)

Substituting (2.6) and (2.7) in (2.1) and (2.3) we obtain a linear system in which  $c_1(\lambda)$  and  $c_2(\lambda)$  are unknowns. Using the solution to the system, we show that the following integral equation is valid on  $[a_2, a_3]$ :

$$\varphi(t,\lambda) = A(\lambda)\sin\lambda (t - a_2) + B(\lambda)\cos\lambda (t - a_2) +$$

$$+ \frac{1}{\lambda} \int_{a_2}^{t} \sin\lambda (t - \xi) \left[ q(\xi) + 2\lambda p(\xi) \right] \varphi(\xi,\lambda) \Delta \xi$$
(2.8)

where

$$A(\lambda) = \left( (a^2 \lambda^2 + b\lambda + c) - \frac{(a^2 \lambda + b\lambda + c + 1)h}{a\lambda^2} \right) \sin \lambda a_1 +$$

$$- \left( \frac{(a^2 \lambda + b\lambda + c)(ah + 1) + 1}{a\lambda} \right) \cos \lambda a_1 +$$

$$- \left( \frac{a^2 \lambda^2 + b\lambda + c}{\lambda} \right) \int_0^{a_1} \cos \lambda (a_1 - \xi) \left[ q(\xi) + 2\lambda p(\xi) \right] \varphi(\xi) \Delta \xi +$$

$$- \left( \frac{a^2 \lambda^2 + b\lambda + c + 1}{a\lambda^2} \right) \int_0^{a_1} \sin \lambda (a_1 - \xi) \left[ q(\xi) + 2\lambda p(\xi) \right] \varphi(\xi) \Delta \xi$$

and

$$B(\lambda) = \left(\frac{h}{\lambda} - a\lambda\right) \sin \lambda a_1 + (ah+1)\cos \lambda a_1 +$$

$$+ a \int_0^{a_1} \cos \lambda (a_1 - \xi) \left[q(\xi) + 2\lambda p(\xi)\right] \varphi(\xi) \Delta \xi +$$

$$+ \frac{1}{\lambda} \int_0^{a_1} \sin \lambda (a_1 - \xi) \left[q(\xi) + 2\lambda p(\xi)\right] \varphi(\xi) \Delta \xi.$$

Thus, we calculate from (2.6) and (2.8) our desired relations.

It follows from  $\Delta(\lambda) = \varphi^{\Delta}(a_3, \lambda) + H\varphi(a_3, \lambda)$  that the asymptotic relation

$$\Delta(\lambda) = \frac{a^2}{2} \lambda^3 \sin 2\lambda a_1 + O\left(\lambda^2 \exp|\tau| \, 2a_1\right) \tag{2.9}$$

is valid for  $|\lambda| \to \infty$ .

Now we are ready to prove our main result. We give proof in the case a > 0. The other case is easier and similar.

#### 3. Proof of the main result

**Proof** [Proof of the Theorem 1.] We give the proof by four steps.

**Step 1:** Let us write the equation (1.1) for  $\varphi$  and  $\widetilde{\varphi}$ 

$$-\varphi^{\Delta\Delta}(t,\lambda) + [q(t) + 2\lambda p(t)]\varphi^{\sigma}(t,\lambda) = \lambda^{2}\varphi^{\sigma}(t,\lambda)$$
(3.1)

$$-\widetilde{\varphi}^{\Delta\Delta}(t,\lambda) + [\widetilde{q}(t) + 2\lambda \widetilde{p}(t)]\widetilde{\varphi}^{\sigma}(t,\lambda) = \lambda^2 \widetilde{\varphi}^{\sigma}(t,\lambda). \tag{3.2}$$

It can be obtained from (3.1) and (3.2) that

$$\left[\varphi^{\Delta}(t,\lambda)\widetilde{\varphi}(t,\lambda) - \varphi(t,\lambda)\widetilde{\varphi}^{\Delta}(t,\lambda)\right]^{\Delta} = \left[q(t) - \widetilde{q}(t) + 2\lambda(p(t) - \widetilde{p}(t))\right]\varphi^{\sigma}(t,\lambda)\widetilde{\varphi}^{\sigma}(t,\lambda). \tag{3.3}$$

Taking  $P(t) = p(t) - \widetilde{p}(t)$ ,  $Q(t) = q(t) - \widetilde{q}(t)$  and  $\Delta$  integrating both sides of (3.3) on  $[0, a_1]$ ,

$$\begin{split} \left[\varphi^{\Delta}(t,\lambda)\widetilde{\varphi}(t,\lambda) - \varphi(t,\lambda)\widetilde{\varphi}^{\Delta}(t,\lambda)\right]_{0}^{a_{1}} &= \int_{0}^{a_{1}}\left[Q(t) + 2\lambda P(t)\right]\varphi^{\sigma}(t,\lambda)\widetilde{\varphi}^{\sigma}(t,\lambda)\Delta t \\ &= \int_{0}^{a_{1}}\left[Q(t) + 2\lambda P(t)\right]\varphi(t,\lambda)\widetilde{\varphi}(t,\lambda)\Delta t. \end{split}$$

It is obvious that

$$\varphi^{\Delta}(a_1,\lambda)\widetilde{\varphi}(a_1,\lambda) - \varphi(a_1,\lambda)\widetilde{\varphi}^{\Delta}(a_1,\lambda) = \int_{0}^{a_1} \left[Q(t) + 2\lambda P(t)\right]\varphi(t,\lambda)\widetilde{\varphi}(t,\lambda)\Delta t.$$

Let

$$K(\lambda) := \int_{0}^{a_1} \left[ Q(t) + 2\lambda P(t) \right] \varphi(t, \lambda) \widetilde{\varphi}(t, \lambda) \Delta t.$$

Since

$$\frac{y_n^{\Delta}(a_1)}{y_n(a_1)} = \frac{\widetilde{y}_n^{\Delta}(a_1)}{\widetilde{y}_n(a_1)},$$

we get that  $K(\lambda_n)=0$  for all  $\lambda_n\in\Lambda$  and so  $\chi(\lambda):=\frac{K(\lambda)}{\Delta(\lambda)}$  is an entire function on  $\lambda$ .

On the other hand, we obtain

$$K(\lambda) = O(\lambda \exp 2 |\tau| a_1)$$

for all complex  $\lambda$  by using the asymptotics (2.4). From (2.9) it can be calculated that

$$|\chi(\lambda)| \le C |\lambda|^{-2}$$
.

for suffciently large  $|\lambda|$ . By the Liouville's theorem,  $\chi(\lambda) = 0$  for all  $\lambda$ . Hence,  $K(\lambda) \equiv 0$ .

**Step 2:** By integrating again both sides of the equality (3.3) on  $(0, a_1)$ , we get

$$\varphi^{\Delta}(a_1,\lambda)\widetilde{\varphi}(a_1,\lambda) - \varphi(a_1,\lambda)\widetilde{\varphi}^{\Delta}(a_1,\lambda) = K(\lambda) = 0$$

and so

$$\varphi^{\Delta}(a_1,\lambda)\widetilde{\varphi}(a_1,\lambda) = \varphi(a_1,\lambda)\widetilde{\varphi}^{\Delta}(a_1,\lambda). \tag{3.4}$$

Put  $\psi(t,\lambda) := \varphi(a_1 - t,\lambda)$ . It is clear that  $\psi(t,\lambda)$  is the solution of the following initial value problem

$$-y^{\Delta\Delta} + [q(a_1 - t) + 2\lambda p(a_1 - t)]y^{\sigma} = \lambda y^{\sigma}, \ t \in (0, a_1)$$
$$y(a_1) = 1, y^{\Delta}(a_1) = -h$$

It follows from (3.4) that

$$\psi^{\Delta}(0,\lambda)\widetilde{\psi}(0,\lambda) = \psi(0,\lambda)\widetilde{\psi}^{\Delta}(0,\lambda).$$

Taking into account Theorem 4.1. in [32], it is concluded that  $q(t) = \widetilde{q}(t)$  and  $p(t) = \widetilde{p}(t)$  on  $[0, a_1]$ .

**Step 3:** To prove that  $q(t) = \widetilde{q}(t)$  and  $p(t) = \widetilde{p}(t)$  on  $[a_2, a_3]$ , we will consider the supplementary problem  $L_1$ :

$$-y^{\nabla\nabla} + [q_1(t) + 2\lambda p_1(t)]y^{\rho} = \lambda y^{\rho}, \ t \in \mathbb{T},$$
$$y^{\nabla}(0) - Hy(0) = y^{\nabla}(a_3) + hy(a_3) = 0,$$

where  $q_1(t) = q(a_3 - t)$  and  $p_1(t) = p(a_3 - t)$ .

By using chain rule for nabla deriative in [3], we have that  $\varphi_1(t,\lambda) = \varphi(a_3 - t,\lambda)$  satisfies the equation

$$-\varphi_1^{\nabla\nabla} + [q_1(t) + 2\lambda p_1(t)]\varphi_1^{\rho} = \lambda \varphi_1^{\rho}$$

and the initial conditions

$$\varphi_1(a_3,\lambda) = 1, \varphi_1^{\nabla}(a_3,\lambda) = -h.$$

Furthermore, the assumption of the theorem

$$\frac{\varphi_1^{\nabla}(a_2, \lambda_n)}{\varphi_1(a_2, \lambda_n)} = \frac{\tilde{\varphi}_1^{\nabla}(a_2, \lambda_n)}{\tilde{\varphi}_1(a_2, \lambda_n)}$$

holds.

If we repeat the calculations in the Step 1, then we replace equation (3.3) by

$$\left[\varphi_1(t,\lambda)\widetilde{\varphi}_1^{\nabla}(t,\lambda) - \varphi_1^{\nabla}(t,\lambda)\widetilde{\varphi}_1(t,\lambda)\right]^{\nabla} = \left[q_1(t) - \widetilde{q}_1(t) + 2\lambda(p_1(t) - \widetilde{p}_1(t))\right]\varphi_1^{\rho}(t,\lambda)\widetilde{\varphi}_1^{\rho}(t,\lambda). \tag{3.5}$$

By integrating (in the sense of  $\nabla$ -integral) both sides of this equality on  $[0, a_2]$ , we obtain

$$\begin{split} \left[ \varphi_1(t,\lambda) \widetilde{\varphi}_1^{\nabla}(t,\lambda) - \varphi_1^{\nabla}(t,\lambda) \widetilde{\varphi}_1(t,\lambda) \right]_0^{a_2} &= \int_0^{a_2} \left[ q_1(t) - \widetilde{q}_1(t) + 2\lambda (p_1(t) - \widetilde{p}_1(t)) \right] \varphi_1^{\rho}(t,\lambda) \widetilde{\varphi}_1^{\rho}(t,\lambda) \nabla t \\ &= \int_0^{a_1} \left[ q_1(t) - \widetilde{q}_1(t) + 2\lambda (p_1(t) - \widetilde{p}_1(t)) \right] \varphi_1(t,\lambda) \widetilde{\varphi}_1(t,\lambda) \nabla t \\ &+ \int_{a_1}^{a_2} \left[ q_1(t) - \widetilde{q}_1(t) + 2\lambda (p_1(t) - \widetilde{p}_1(t)) \right] \varphi_1^{\rho}(t,\lambda) \widetilde{\varphi}_1^{\rho}(t,\lambda) \nabla t \end{split}$$

From Step 2, since  $q(a_1) = \widetilde{q}(a_1)$  and  $p(a_1) = \widetilde{p}(a_1)$ , then  $q_1(a_2) = \widetilde{q}_1(a_2)$  and  $p_1(a_2) = \widetilde{p}_1(a_2)$ . Thus we get

$$\int_{a_1}^{a_2} \left[q_1(t) - \widetilde{q}_1(t)\right] \varphi^{\rho}(t,\lambda) \widetilde{\varphi}^{\rho}(t,\lambda) \nabla t = \left[q_1(a_2) - \widetilde{q}_1(a_2)\right] \varphi_1^{\rho}(a_2,\lambda) \widetilde{\varphi}_1^{\rho}(a_2,\lambda) (a_2 - a_1) = 0.$$

Therefore we have

$$\varphi_1^{\nabla}(a_2,\lambda)\widetilde{\varphi}_1(a_2,\lambda) - \varphi_1(a_2,\lambda)\widetilde{\varphi}_1^{\nabla}(a_2,\lambda) = \int_0^{a_1} \left[q_1(t) - \widetilde{q}_1(t) + 2\lambda(p_1(t) - \widetilde{p}_1(t))\right] \varphi_1(t,\lambda)\widetilde{\varphi}_1(t,\lambda)\nabla t.$$

Let

$$K_1(\lambda) := \int_0^{a_1} \left[ q_1(t) - \widetilde{q}_1(t) + 2\lambda (p_1(t) - \widetilde{p}_1(t)) \right] \varphi_1(t, \lambda) \widetilde{\varphi}_1(t, \lambda) \nabla t.$$

It is obvious that  $K_1(\lambda_n) = 0$  for all  $\lambda_n \in \Lambda$  and so  $\omega(\lambda) := \frac{K_1(\lambda)}{\Delta(\lambda)}$  is entire on  $\lambda$ . Similar to the calculations in the last part of Step 1, we obtain

$$K_1(\lambda) \equiv 0.$$

**Step 4:** By integrating again both sides of the equality (3.5) on  $(0, a_1)$ , we get

$$\varphi_1^{\nabla}(a_1,\lambda)\widetilde{\varphi}_1(a_1,\lambda) - \varphi_1(a_1,\lambda)\widetilde{\varphi}_1^{\nabla}(a_1,\lambda) = K_1(\lambda) = 0.$$

Repeating the Step 2 for the supplementary problem  $L_1$  and  $\varphi_1(t,\lambda)$ , it is concluded that  $q_1(t) = \widetilde{q}_1(t)$  and  $p_1(t) = \widetilde{p}_1(t)$  on  $[0,a_1]$ , that is  $q(t) = \widetilde{q}(t)$  and  $p(t) = \widetilde{p}(t)$  on  $[a_2,a_3]$ . This completes the proof.

**Example 3.1** Consider the following problems on  $\mathbb{T} = [0, 1/2] \cup [1, 3/2]$ ,

$$L_0: \begin{cases} -y^{\Delta\Delta}(t) = \lambda y^{\sigma}(t), & t \in [0, 1/2] \cup [1, 3/2] \\ y^{\Delta}(0) = 0 \\ y^{\Delta}(3/2) = 0 \end{cases}$$

and

$$\widetilde{L}_0: \left\{ \begin{array}{c} -y^{\Delta\Delta}(t) + [q(t) + 2\lambda p(t)]y^{\sigma}(t) = \lambda y^{\sigma}(t), & t \in [0, 1/2] \cup [1, 3/2] \\ y^{\Delta}(0) = 0 \\ y^{\Delta}(3/2) = 0. \end{array} \right.$$

Let  $\Lambda_0 := \{\lambda_n, n \in \mathbb{Z}\}$  and  $\widetilde{\Lambda}_0 := \{\widetilde{\lambda}_n, n \in \mathbb{Z}\}$  be the eigenvalues sets of L and  $\widetilde{L}_0$ ,  $y_n(t)$  and  $\widetilde{y}_n(t)$  are eigenfunctions related to this eigenvalues, respectively. According to Theorem-1, if  $\Lambda_0 = \widetilde{\Lambda}_0$  and for any  $n \in \mathbb{N}$ ,

$$\frac{y_n^{\nabla}(1)}{y_n(1)} = \frac{\widetilde{y}_n^{\nabla}(1)}{\widetilde{y}_n(1)}$$

then  $q(t) \equiv p(t) = 0$  on  $\mathbb{T}$ .

### References

- [1] Adalar I. Hochstadt-Lieberman type theorems for a diffusion operator on a time scale. Quaestiones Mathematicae 2021; doi: 10.2989/16073606.2021.1895351
- [2] Adalar I, Ozkan AS. An interior inverse Sturm-Liouville problem on a time scale. Analysis and Mathematical Physics 2020: 10 (4): 1-10.
- [3] Bartosiewicz Z, Piotrowska E. Lyapunov functions in stability of nonlinear systems on time scales. Journal of Difference Equations and Applications. 2020; 17 (3): 309-325.
- [4] Bohner M, Peterson A. Dynamic Equations on Time Scales. Birkha"user, Boston MA, 2001.

## ADALAR/Turk J Math

- [5] Bohner M, Peterson A. Advances in Dynamic Equations on Time Scales. Birkha"user, Boston MA, 2003.
- [6] Bohner M, Georgiev, SG. Multivariable dynamic calculus on time scales. Cham, Switzerland: Springer, 2016.
- [7] Bohner M, Koyunbakan H. Inverse problems for Sturm–Liouville difference equations. Filomat 2016; 30 (5): 1297-1304.
- [8] Bondarenko NP. Inverse problem for the differential pencil on an arbitrary graph with partial information given on the coefficients. Analysis and Mathematical Physics 2019 9 1393–1409.
- [9] Buterin SA, Yurko VA. Inverse spectral problem for pencils of differential operators on a finite interval, Vestnik Bashkir. Univ. 2006; 48–12.
- [10] Buterin SA, Shieh CT. Inverse nodal problem for differential pencils. Applied Mathematics Letters 2009; 22 (8): 1240-1247.
- [11] Buterin SA. On half inverse problem for differential pencils with the spectral parameter in boundary conditions. Tamkang journal of mathematics 2011; 42 (3): 355-364.
- [12] Buterin SA, Shieh CT. Incomplete inverse spectral and nodal problems for differential pencil, Results Math. 2012; 62 167–179.
- [13] Cakmak Y, Isık S. Half Inverse problem for the impulsive diffusion operator with discontinuous coeffcient, Filomat 2016; 30 (1): 157-168.
- [14] Freiling G, Yurko VA. Inverse Sturm-Liouville Problems and their Applications. New York, Nova Science, 2001.
- [15] Gasymov MG, Guseinov GS, Determination of a diffusion operator from the spectral data. Dokl. Akad. Nauk Azerbaijan SSSR 1981; 37 (2): 19-23.
- [16] Hilger S. Analysis on measure chains a unified approach to continuous and discrete calculus. Results in Math. 1990; 18 18–56.
- [17] Jaulent M, Jean C. The inverse s-wave scattering problem for a class of potentials depending on energy. Comm. Math. Phys. 1972; 28 177–220.
- [18] Koyunbakan H. Inverse problem for a quadratic pencil of Sturm-Liouville operator. J. Math. Anal. Appl. 2011; 378 (2): 549-554.
- [19] Kuznetsova MA. A Uniqueness Theorem on Inverse Spectral Problems for the Sturm-Liouville Differential Operators on Time Scales. Results Math 2020; 75 (2): 1-23.
- [20] Kuznetsova MA. On Reconstruction of Sturm-Liouville Differential Operators on Time Scales. Mat. notes. 2021; 109 (1): 82-100.
- [21] Mochizuki K, Trooshin I. Inverse problem for interior spectral data of Sturm-Liouville operator. J. Inverse Ill-posed Probl. 2001; 9 425-433.
- [22] Mosazadeh S. A new approach to uniqueness for inverse Sturm-Liouville problems on finite intervals. Turk. J. Math. 2017; 41 1224-1234.
- [23] Neamaty A, Khalili Y. Determination of a differential operator with discontinuity from interior spectral data. Inverse Problems in Science and Engineering 2014; 22 (6): 1002-1008.
- [24] Ozkan AS. Sturm-Liouville operator with parameter-dependent boundary conditions on time scales. Electron. J. Differential Equations 2017; 212 1-10.
- [25] Ozkan AS. Ambarzumyan-type theorems on a time scale. Journal of Inverse and Ill-posed Problems. 2018 26 (5): 633–637.
- [26] Ozkan, AS. Boundary-value problem for a class of second-order parameter-dependent dynamic equations on a time scale. Mathematical Methods in the Applied Sciences 2020; 43 (7): 4353-4359.
- [27] Ozkan AS, Adalar I. Half-inverse Sturm-Liouville problem on a time scale. Inverse Problems 2020; 36 025015.

## ADALAR/Turk J Math

- [28] Wang YP, Inverse problems for Sturm-Liouville operators with interior discontinuities and boundary conditions dependent on the spectral parameter. Mathematical Methods in the Applied Sciences 2013; 36 (7): 857-868.
- [29] Wang YP. An interior inverse problem for Sturm-Liouville operators with eigenparameter dependent boundary conditions. Tamkang Journal of Mathematics 2011; 42 (3): 395-403.
- [30] Wang YP. The inverse problem for differential pencils with eigenparameter dependent boundary conditions from interior spectral data. Applied Mathematics Letters 2012; 25 (7): 1061-1067.
- [31] Yang CF, Guo YX. Determination of a differential pencil from interior spectral data. Journal of Mathematical Analysis and Applications 2011; 375 (1): 284-293.
- [32] Yang CF. Uniqueness theorems for differential pencils with eigenparameter boundary conditions and transmission conditions. J. Differ. Equ. 2013; 255 2615-2635.
- [33] Yang CF, Yang XP. An interior inverse problem for the Sturm-Liouville operator with discontinuous conditions. Applied Mathematics Letters 2009; 22 (9): 1315-1319.
- [34] Yang CF, Bondarenko NP, Xu XC. An inverse problem for the Sturm-Liouville pencil with arbitrary entire functions in the boundary condition. Inverse Problems & Imaging 2020; 14 (1): 153.
- [35] Yurko VA. On determination of functional-differential pencils on closed sets from the Weyl-type function. Izv. Saratov Univ. (N.S.), Ser. Math. Mech. Inform. 2020; 20 (3): 343–350.