# SOME RESULTS ON SPACE-LIKE LINE CONGRUENCES AND THEIR SPACE-LIKE PARAMETER RULED SURFACES 

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#### Abstract

Using dual vectors, space-like ruled surfaces was introduced in [2]. In this paper, we define space-like line congruence and their space-like parameter ruled surfaces $\mathbf{R}_{11}(\mathbf{t})$ and $\mathbf{R}_{\mathbf{2 1}}(\mathbf{t})$ of $R_{1}^{3}$. Then, by choosing space-like parameter ruled surfaces as space-like principle ruled surfaces, we obtain some relations among the magnitudes of space-like ruled surfaces $\mathbf{R}_{1}(\mathbf{t}), \mathbf{R}_{11}(\mathbf{t})$.


## 1. Introduction

Let $A=a+\epsilon a_{0}$ be a dual number, $A \in I D=\left\{\left(a, a_{0}\right) \mid a, a_{0} \in R\right\}$ and ID be a commutative ring with a unit element. We call the dual number $\epsilon=(0,1) \in I D$ dual unit and $\epsilon^{2}=(0,0) .\left(I D^{3},+\right)$ is a module on the dual number ring. We call it ID-module, and dual vectors are the elements of this modul. We denote dual unit vector $\mathbf{A}$ as

$$
\begin{equation*}
\mathbf{A}=\left(\mathbf{a}, \mathbf{a}_{\mathbf{0}}\right)=\mathbf{a}+\epsilon \mathbf{a}_{\mathbf{0}} \quad, \quad \mathbf{a} \mathbf{a}=\mathbf{1}, \mathbf{a} \mathbf{a}_{\mathbf{0}}=\mathbf{0} \tag{1}
\end{equation*}
$$

where $\mathbf{a}, \mathbf{a}_{\mathbf{0}} \in \mathbf{I R}^{\mathbf{3}}$. The dual vectors with unit length correspond to oriented lines of $E^{3}$ [1].

Theorem 1.1. The oriented lines in $I R^{3}$ are in one-to-one correspondence with the points of the dual unit sphere in $I D^{3}$.

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The scalar product of two dual vectors $\mathbf{A}=\left(\mathbf{a}, \mathbf{a}_{\mathbf{0}}\right)=\mathbf{a}+\epsilon \mathbf{a}_{\mathbf{0}}$ and $\mathbf{B}=\left(\mathbf{b}, \mathbf{b}_{\mathbf{0}}\right)=\mathbf{b}+\epsilon \mathbf{b}_{\mathbf{0}}$ is

$$
\begin{equation*}
\mathbf{A B}=\mathbf{a} \cdot \mathbf{b}+\epsilon\left(\mathbf{b a}_{\mathbf{0}}+\mathbf{a} \mathbf{b}_{\mathbf{0}}\right)=\cos \varphi-\epsilon \varphi^{\star} \sin \varphi, \tag{2}
\end{equation*}
$$

where, $\varphi$ is the real angle between dual unit vectors $\mathbf{A}$ and $\mathbf{B}$ and $\varphi^{\star}$ is the shortest distance between the lines.

The Blaschke trihedron $\left(A_{1}, A_{2}, A_{3}\right)$ depends on the ruled surface striction of $\mathbf{A}_{\mathbf{1}}(\mathbf{t})$ in dual space $D^{3}[1]$. According to this, the first axis $\mathbf{A}_{\mathbf{1}}$ of the trihedron is the generator which passes from the striction point of the ruled surface, the second axis $\mathbf{A}_{\mathbf{2}}$ is the surface normal at this point and finally the third axis $\mathbf{A}_{\mathbf{3}}$ is the tangent of the striction line at this point.

## 2. Dual Lorentzian Space $D_{1}^{3}$

Let we consider vector space $R_{1}^{3}$ of $R^{3}$ provided with Lorentzian inner product of signature $(+,+,-)$. For any vector $\mathbf{a}=\left(a_{1}, a_{2}, a_{3}\right)$ of $R_{1}^{3}$;
i) if $\langle a, a\rangle>0$, a is said to be space-like,
ii) if $\langle a, a\rangle<0$, a is said to be time-like,
iii) if $\langle a, a\rangle=0$, $\mathbf{a}$ is said to be light-like (null)

The Lorentzian and hyperbolic sphere of radius 1 in $R_{1}^{3}$ are defined by

$$
\begin{gather*}
S_{1}^{2}=\left\{a=\left(a_{1}, a_{2}, a_{3}\right) \in R_{1}^{3},\langle a, a\rangle=1\right\}  \tag{3}\\
H_{0}^{2}=\left\{a=\left(a_{1}, a_{2}, a_{3}\right) \in R_{1}^{3},\langle a, a\rangle=-1\right\} \tag{4}
\end{gather*}
$$

respectively.
By considering the Lorentzian inner product, we may write inner product of $\mathbf{A}$ and B as follows:

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$$
\begin{equation*}
\mathbf{A B}=\mathbf{a} \cdot \mathbf{b}+\epsilon\left(\mathbf{b} \mathbf{a}_{\mathbf{0}}+\mathbf{a} \mathbf{b}_{\mathbf{0}}\right) \tag{5}
\end{equation*}
$$

We call it dual Lorentzian space which is defined and denote by $I D_{1}^{3}$.

Definition 2.1. Let $\mathbf{A}=\left(\mathbf{a}, \mathbf{a}_{\mathbf{0}}\right)=\mathbf{a}+\epsilon \mathbf{a}_{\mathbf{0}} \in I D_{1}^{3}$. The dual vector $\mathbf{A}$ is said to be space-like if the vector $\mathbf{a}$ is space-like, time-like if the vector $\mathbf{a}$ is time-like, and light-like (or null) if the vector $\mathbf{a}$ is light-like.

We also defined the time orientation as follows:
A time-like vector $\mathbf{A}=\mathbf{a}+\epsilon \mathbf{a}_{\mathbf{0}}$ is future pointing if the vector $\mathbf{a}$ is future pointing.

Definition 2.2. Let $\mathbf{A}, \mathbf{B} \in I D_{1}^{3}$. We define the Lorentzian cross product of $\mathbf{A}$ and $\mathbf{B}$ by

$$
A \wedge B=\left|\begin{array}{lll}
E_{1} & E_{2} & -E_{3}  \tag{6}\\
A_{1} & A_{2} & A_{3} \\
B_{1} & B_{2} & B_{3}
\end{array}\right|
$$

where, $\mathbf{A}=\left(A_{1}, A_{2}, A_{3}\right), \mathbf{B}=\left(B_{1}, B_{2}, B_{3}\right)$ and $E_{1} \wedge E_{2}=E_{3}, E_{2} \wedge E_{3}=-E_{3}, E_{3}, E_{3} \wedge$ $E_{1}=-E_{2},[2]$.

## 3. Space-Like Congruence in Dual Lorentzian Space $I D_{1}^{3}$

A space-like line congruence in line space $R_{1}^{3}$ can be represented by a unit space like dual vector which is depending on two real parameters $u$ and $v$ as follows:

$$
\begin{equation*}
\mathbf{R}(u, v)=\mathbf{r}(u, v)+\epsilon \mathbf{r}^{\star}(u, v) \quad, \quad \mathbf{R}^{2}=1 \tag{7}
\end{equation*}
$$

The dual arc element of a space like ruled surface of space-like congruence can be given as

$$
\begin{align*}
d S^{2} & =d \mathbf{R}^{2}=\left(\mathbf{R}_{u} d u+\mathbf{R}_{\mathbf{v}} d v\right)^{2} \\
& =-E d u^{2}+2 F d u d v+G d v^{2} \tag{8}
\end{align*}
$$

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Where

$$
\begin{array}{rcl}
E=-R_{u}^{2} & F=R_{u} R_{v} & G=R_{v}^{2} \\
E=e+\epsilon e^{\star} & F=f+\epsilon f^{\star} & G=g+\epsilon g^{\star}  \tag{9}\\
& & \\
e=-r_{u}^{2} & f=r_{u} r_{v} & g=r_{v}^{2} \\
e^{\star}=2 r_{u} r_{u}^{\star} & f^{\star}=r_{u} r_{v}^{\star} & g^{\star}=2 r_{v} r_{v}^{\star}
\end{array}
$$

The differential forms I and II of the space-like line congruence are

$$
\begin{gather*}
I=-e d u^{2}+2 f d u d v+g d v^{2} \\
I I=-e^{\star} d u^{2}+2 f^{\star} d u d v+g^{\star} d v^{2} \tag{10}
\end{gather*}
$$

respectively. If we use the relations (8),(9) and (10), we have

$$
\begin{equation*}
d S^{2}=I+\epsilon I I \tag{11}
\end{equation*}
$$

and the dral of a space-like ruled surface of space like congruence can write as

$$
\begin{equation*}
\frac{1}{d}=\frac{I I}{2 I} \tag{12}
\end{equation*}
$$

Definition 3.3. Let $\frac{1}{d_{1}}$ and $\frac{1}{d_{2}}$ be extremum values of the dral. These values are called principle drals.

The principle drals can calculated following relation:

$$
\left|\begin{array}{ll}
-e d u+f d v & -e^{\star} d u+f^{\star} d v  \tag{13}\\
f d u+g d v & f^{\star} d u+g^{\star} d v
\end{array}\right|
$$

Thus, we may write mean dral and Gaussian dral as follows, respectively:

$$
\begin{equation*}
h=\frac{1}{2}\left(\frac{1}{d_{1}}+\frac{1}{d_{2}}\right) \tag{14}
\end{equation*}
$$

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$$
\begin{equation*}
k=\frac{1}{d_{1} d_{2}} \tag{15}
\end{equation*}
$$

Definition 3.4. The space-like ruled surfaces which are obtained by the relation (13) of the space like congruence are called the space-like principle ruled surfaces.

Definition 3.5. The space-like ruled surfaces $u=$ constant and $v=$ constant of a space-like line congruence are called the space-like parameter ruled surfaces.
4. The Relations Among The Magnitudes of the Space Like Ruled Surfaces $\mathbf{R}_{1}, \mathbf{R}_{11}$ and $\mathbf{R}_{21}$

Let we consider a space-like ruled surface $\mathbf{R}=\mathbf{R}(\mathrm{t})$ of the space-like congruence $\mathbf{R}=\mathbf{R}(u, v)$, Where, $u$ and $v$ are functions of $t$.

Let we write the space-like parameter ruled surfaces as

$$
\begin{equation*}
\mathbf{R}_{\mathbf{1 1}}=\mathbf{R}_{\mathbf{1 1}}\left(u, v_{0}\right) \text { and } \mathbf{R}_{\mathbf{2 1}}=\mathbf{R}_{\mathbf{2 1}}\left(u_{0}, v\right) \tag{16}
\end{equation*}
$$

The space-like ruled surfaces $\mathbf{R}_{\mathbf{1}}, \mathbf{R}_{\mathbf{1 1}}$ and $\mathbf{R}_{\mathbf{2 1}}$ have common space-like line which is defined by the following relation:

$$
\begin{equation*}
\mathbf{R}_{\mathbf{0}}=\mathbf{R}\left(u_{0}, v_{0}\right)=\mathbf{R}_{\mathbf{1}}\left(u_{0}, v_{0}\right)=\mathbf{R}_{\mathbf{2 1}}\left(u_{0}, v_{0}\right) \tag{17}
\end{equation*}
$$

Blaschke trihedrons of these space-like ruled surfaces are in the following form:

$$
\begin{equation*}
\left(\mathbf{R}_{0}=\mathbf{R}_{1}, \mathbf{R}_{2}, \mathbf{R}_{3}\right) \quad,\left(\mathbf{R}_{0}=\mathbf{R}_{11}, \mathbf{R}_{12}, \mathbf{R}_{13}\right) \quad,\left(\mathbf{R}_{0}=\mathbf{R}_{21}, \mathbf{R}_{22}, \mathbf{R}_{23}\right) \tag{18}
\end{equation*}
$$

Where $\mathbf{R}_{\mathbf{2}}, \mathbf{R}_{\mathbf{1 2}}$ and $\mathbf{R}_{\mathbf{2 3}}$ are time-like, $\mathbf{R}_{\mathbf{1 3}}, \mathbf{R}_{\mathbf{3}}, \mathbf{R}_{\mathbf{2 2}}$ and $\mathbf{R}_{\mathbf{0}}$ are space-like. Thus we may write
$\mathbf{R}_{1}^{2}=\mathbf{R}_{3}^{2}=1, \mathbf{R}_{2}^{2}=-1 \mathbf{R}_{3} \wedge \mathbf{R}_{1}=\mathbf{R}_{2}, \mathbf{R}_{2} \wedge \mathbf{R}_{3}=-\mathbf{R}_{1}, \mathbf{R}_{1} \wedge \mathbf{R}_{2}=-\mathbf{R}_{3}$
$\mathbf{R}_{11}^{2}=\mathbf{R}_{13}^{2}=\mathbf{1}, \mathbf{R}_{12}^{2}=-1 \mathbf{R}_{13} \wedge \mathbf{R}_{11}=+\mathbf{R}_{12}, \mathbf{R}_{12} \wedge \mathbf{R}_{13}=-\mathbf{R}_{11}, \mathbf{R}_{11} \wedge \mathbf{R}_{12}=-\mathbf{R}_{13}$
$\mathbf{R}_{21}^{2}=\mathbf{R}_{22}^{2}=1, \mathbf{R}_{23}^{2}=-1 \mathbf{R}_{23} \wedge \mathbf{R}_{21}=-\mathbf{R}_{22}, \mathbf{R}_{22} \wedge \mathbf{R}_{23}=-\mathbf{R}_{21}, \mathbf{R}_{21} \wedge \mathbf{R}_{22}=\mathbf{R}_{23}$

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On the other hand, if we choose the space-like paremeter ruled surfaces as space like principle ruled surfaces, we may write $\mathrm{f}=0$ and $f^{\star}=0$. Thus,

$$
\begin{equation*}
F=0, \quad \mathbf{R}_{\mathbf{u}} \cdot \mathbf{R}_{\mathbf{v}}=0 \tag{20}
\end{equation*}
$$

can write.
The dual arc elements of the space-like ruled surfaces $\mathbf{R}_{\mathbf{1}}, \mathbf{R}_{\mathbf{1 1}}$ and $\mathbf{R}_{\mathbf{2 1}}$ can be given respectively as

$$
\begin{equation*}
d S=P d t \quad, d S_{1}=P_{1} d u \quad, d S_{2}=P_{2} d v \tag{21}
\end{equation*}
$$

where

$$
\begin{gather*}
P_{1}=\sqrt{\left|\mathbf{R}_{\mathbf{u}}^{2}\right|}=\sqrt{-\mathbf{R}_{\mathbf{u}}^{2}}=\sqrt{E} P_{2}=\sqrt{\mathbf{R}_{\mathbf{v}}^{2}}=\sqrt{\mathbf{G}} \quad P=\sqrt{\mathbf{R}_{\mathbf{I}}^{\prime 2}}  \tag{22}\\
\mathbf{R}_{\mathbf{1}}^{\prime}=\mathbf{R}_{u} \frac{d u}{d t}+\mathbf{R}_{\mathbf{v}} \frac{d v}{d t}
\end{gather*}
$$

Using (21) and (22), we have

$$
\begin{equation*}
d S_{1}=\sqrt{E} d u a n d d S_{2}=\sqrt{G} d v \tag{23}
\end{equation*}
$$

The derivative formulas of these Blaschke trihedrons, defined by (18), are

$$
\begin{align*}
& \mathbf{R}_{1}^{\prime}=\mathbf{P} \mathbf{R}_{2} \mathbf{R}_{2}^{\prime}=\mathbf{P} \mathbf{R}_{\mathbf{1}}+\mathbf{Q} \mathbf{R}_{3} \mathbf{R}_{3}^{\prime}=\mathbf{Q} \mathbf{R}_{2} \\
& \mathbf{R}_{11}^{\prime}=\mathbf{P}_{1} \mathbf{R}_{12} \mathbf{R}_{12}^{\prime}=\mathbf{P}_{1} \mathbf{R}_{12}+\mathbf{Q}_{1} \mathbf{R}_{13} \mathbf{R}_{13}^{\prime}=\mathbf{Q}_{1} \mathbf{R}_{12}  \tag{24}\\
& \mathbf{R}_{21}^{\prime}=\mathbf{P}_{2} \mathbf{R}_{22} \mathbf{R}_{22}^{\prime}=-\mathbf{P}_{2} \mathbf{R}_{21}-\mathbf{Q}_{2} \mathbf{R}_{23} \mathbf{R}_{23}^{\prime}=-\mathbf{Q}_{2} \mathbf{R}_{22}
\end{align*}
$$

Thus, the Blaschke vectors of the Blaschke trihedrons can be given by the following relations:

$$
\begin{equation*}
\mathbf{B}=-\mathbf{Q} \mathbf{R}_{0}+\mathbf{P} \mathbf{R}_{3} \quad, \mathbf{B}_{1}=-\mathbf{Q}_{1} \mathbf{R}_{0}+\mathbf{P}_{1} \mathbf{R}_{13} \quad, \mathbf{B}_{2}=-\mathbf{Q}_{2} \mathbf{R}_{0}-\mathbf{P}_{2} \mathbf{R}_{23} \tag{25}
\end{equation*}
$$

respectively.

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The dual unit vectors $\mathbf{R}_{\mathbf{2}}, \mathbf{R}_{\mathbf{1 2}}$ and $\mathbf{R}_{\mathbf{2 2}}$ which are the second edges of the trihedrons (18), can be written as

$$
\begin{array}{r}
\mathbf{R}_{2}=\frac{\mathbf{R}_{1}^{\prime}}{\mathbf{P}}=\frac{\mathbf{1}}{\mathbf{P}}\left(\mathbf{R}_{\mathbf{u}} \frac{\mathbf{d u}}{\mathbf{d t}}+\mathbf{R}_{\mathbf{v}} \frac{\mathbf{d v}}{\mathbf{d t}}\right) \\
\mathbf{R}_{12}=\frac{\mathbf{R}_{11}^{\prime}}{\mathbf{P}_{1}}=\frac{\mathbf{R}_{\mathbf{u}}}{\mathbf{P}_{1}}=\frac{\mathbf{R}_{\mathbf{u}}}{\sqrt{\mathbf{E}}} \\
\mathbf{R}_{22}=\frac{\mathbf{R}_{21}^{\prime}}{\mathbf{P}_{2}}=\frac{\mathbf{R}_{\mathbf{v}}}{\mathbf{P}_{2}}=\frac{\mathbf{R}_{\mathbf{v}}}{\sqrt{\mathbf{G}}} \tag{28}
\end{array}
$$

Using the relation

$$
\begin{equation*}
\mathbf{R}=\frac{\mathbf{R}_{\mathbf{u}} \times \mathbf{R}_{\mathbf{v}}}{\left\|\mathbf{R}_{\mathbf{u}} \times \mathbf{R}_{\mathbf{v}}\right\|} \tag{29}
\end{equation*}
$$

, we can express the relation (29) for the common space-like line $\mathbf{R}_{\mathbf{0}}$ as

$$
\begin{equation*}
\mathbf{R}_{0}=\frac{\mathbf{R}_{u}\left(u_{0}, v_{0}\right) x \mathbf{R}_{v}\left(u_{0}, v_{0}\right)}{\left\|\mathbf{R}_{u}\left(u_{0}, v_{0}\right) x \mathbf{R}_{v}\left(u_{0}, v_{0}\right)\right\|}=\mathbf{R}\left(u_{0}, v_{0}\right) \tag{30}
\end{equation*}
$$

In our study, we will take the space-like parameter ruled surfaces as the space-like principle ruled surfaces. Thus, if we consider (20), (27) and (28), we have

$$
\begin{equation*}
\mathbf{R}_{12} \cdot \mathbf{R}_{22}=0 \tag{31}
\end{equation*}
$$

From the relation below

$$
\begin{equation*}
\mathbf{R}_{12} \times \mathbf{R}_{22}=\frac{\mathbf{R}_{\mathbf{u}} \times \mathbf{R}_{\mathbf{v}}}{\sqrt{\mathbf{E G}}}=\mathbf{R}_{0} \tag{32}
\end{equation*}
$$

and from (26), (27) and (28), we obtain

$$
\begin{equation*}
\mathbf{R}_{2}=\frac{\mathbf{P}_{1}}{\mathbf{P}} \frac{\mathrm{du}}{\mathrm{dt}} \mathbf{R}_{12}+\frac{\mathbf{P}_{2}}{\mathbf{P}} \frac{\mathrm{dv}}{\mathrm{dt}} \mathbf{R}_{22} \tag{33}
\end{equation*}
$$

Thus, using the dual angle between the unit dual vectors $\mathbf{R}_{\mathbf{2}}$ and $\mathbf{R}_{12}$, we may write

$$
\begin{equation*}
\mathbf{R}_{\mathbf{2}}=\operatorname{ch} \Phi \mathrm{R}_{12}+\operatorname{sh} \Phi \mathrm{R}_{22} \tag{34}
\end{equation*}
$$

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where

$$
\begin{gather*}
\operatorname{ch\Phi }=\frac{d S_{1}}{d S}=\sqrt{E} \frac{d u}{d S} \quad, \operatorname{sh} \Phi=\frac{d S_{2}}{d S}=\sqrt{G} \frac{d v}{d S}  \tag{35}\\
t g \Phi=\frac{d S_{2}}{d S_{1}}=\sqrt{\frac{G}{E}} \frac{d v}{d u} \tag{36}
\end{gather*}
$$

and from (19) and (34)

$$
\begin{align*}
& -\mathbf{R}_{\mathbf{3}}=\mathbf{R}_{\mathbf{0}} \times \mathbf{R}_{\mathbf{2}}=\mathbf{R}_{\mathbf{0}} \times\left(c h \Phi \mathbf{R}_{\mathbf{1 2}}+s h \Phi \mathbf{R}_{\mathbf{2 2}}\right)  \tag{37}\\
& -\mathbf{R}_{\mathbf{3}}=\operatorname{ch} \Phi \mathbf{R}_{\mathbf{1 3}}-\operatorname{sh\Psi \mathbf {R}_{\mathbf {23}}} \tag{38}
\end{align*}
$$

are obtained. Then, by the relations (19) and (32), we have

$$
\begin{align*}
& \mathbf{R}_{12} \times \mathbf{R}_{13}=-\mathbf{R}_{\mathbf{0}}, \quad \mathbf{R}_{22} \times \mathbf{R}_{23}=-\mathbf{R}_{\mathbf{0}} \quad, \mathbf{R}_{12} \times \mathbf{R}_{22}=\mathbf{R}_{\mathbf{0}} \\
& \mathbf{R}_{12} \times\left(\mathbf{R}_{13}+\mathbf{R}_{22}\right)=\mathbf{0} \rightarrow \mathbf{R}_{13}+\mathbf{R}_{22}=\mathbf{M} \mathbf{R}_{12}  \tag{39}\\
& \mathbf{R}_{22} \times\left(\mathbf{R}_{23}-\mathbf{R}_{12}\right)=\mathbf{0} \rightarrow \mathbf{R}_{23}-\mathbf{R}_{12}=\mathbf{N} \mathbf{R}_{22} \tag{40}
\end{align*}
$$

where M and N are dual scalars. Taking dot product of the both sides of (39) and (40) by the unit dual vectors $\mathbf{R}_{\mathbf{1 2}}$ and $\mathbf{R}_{\mathbf{2 2}}$, respectively, and considering the relations (18) and (31), we have $M=0$ and $N=0$. Then, if we insert the values $M$ and $N$ into (39) and (40), we have

$$
\begin{align*}
& \mathbf{R}_{12}=-\mathbf{R}_{22}  \tag{41}\\
& \mathbf{R}_{23}=\mathbf{R}_{12}
\end{align*}
$$

Finally, we obtain following teorem:

Theorem 4.2. The edges of Blascke trihedrons of the space-like parameter ruled surfaces coincide with each other under the condition that their directions and orders are not the same.

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Result 1. The Blaschke vectors $\mathbf{B}, \mathbf{B}_{\mathbf{1}}$ and $\mathbf{B}_{\mathbf{2}}$ can be expressed following form by the vectors $\mathbf{R}_{\mathbf{1 2}}, \mathbf{R}_{\mathbf{2 2}}$ and $\mathbf{R}_{\mathbf{0}}$.

It is clear that

$$
\begin{align*}
\mathbf{B} & =-\mathbf{Q} \mathbf{R}_{0}+\mathbf{P}\left(-\operatorname{ch} \Phi \mathbf{R}_{22}-\operatorname{sh} \Phi \mathbf{R}_{12}\right) \\
\mathbf{B}_{1} & =-\mathbf{Q}_{1} \mathbf{R}_{0}-\mathbf{P}_{1} \mathbf{R}_{22}  \tag{42}\\
\mathbf{B}_{2} & =-\mathbf{Q}_{2} \mathbf{R}_{0}-\mathbf{P}_{2} \mathbf{R}_{12}
\end{align*}
$$

can be written easily from (38) and (41)
Result 2. If the trihedron $\left(\mathbf{R}_{\mathbf{0}}, \mathbf{R}_{\mathbf{1 2}}=\mathbf{R}_{\mathbf{2 3}}, \mathbf{R}_{\mathbf{1 3}}=-\mathbf{R}_{\mathbf{2 2}}\right)$ moves on the striction curves of the space-like parameter ruled surface $\mathbf{R}_{\mathbf{1 1}}$, it changes as a function of dual arc $S_{1}$ of $\mathrm{v}=$ constant ruled surface. If the $\left(\mathbf{R}_{\mathbf{0}}, \mathbf{R}_{\mathbf{1 2}}=\mathbf{R}_{\mathbf{2 3}}, \mathbf{R}_{\mathbf{1 3}}=-\mathbf{R}_{\mathbf{2 2}}\right)$ moves on the striction curves of the space-likeparameter ruled surface $\mathbf{R}_{\mathbf{1 1}}$, it changes as a function of dual arc $S_{2}$ of $\mathrm{u}=$ constant ruled surface. Thus, the edges of this trihedron are depend on two parameters

Theorem 4.3. If we consider Blaschke trihedrons and their derivative formulaes of the space-like ruled surface which are determined by (18) and (24), we have

$$
\begin{align*}
& \frac{\partial \mathbf{R}_{\mathbf{0}}}{\partial u}=\mathbf{B}_{\mathbf{1}} \times \mathbf{R}_{\mathbf{0}}, \frac{\partial \mathbf{R}_{2 \mathbf{3}}}{\partial u}=\mathbf{B}_{\mathbf{1}} \times \mathbf{R}_{23}, \frac{\partial \mathbf{R}_{\mathbf{2 2}}}{\partial u}=\mathbf{B}_{\mathbf{1}} \times \mathbf{R}_{2 \mathbf{2}}  \tag{43}\\
& \frac{\partial \mathbf{R}_{0}}{\partial v}=\mathbf{B}_{\mathbf{2}} \times \mathbf{R}_{\mathbf{0}}, \frac{\partial \mathbf{R}_{13}}{\partial v}=\mathbf{B}_{\mathbf{2}} \times \mathbf{R}_{13}, \frac{\partial \mathbf{R}_{12}}{\partial v}=\mathbf{B}_{\mathbf{2}} \times \mathbf{R}_{12}
\end{align*}
$$

Proof. If we write derivative formulas of the edges of Blaschke trihedrons by the Blaschke vectors, we have

$$
\begin{align*}
& \frac{\partial \mathbf{R}_{\mathbf{0}}}{\partial u}=\mathbf{B}_{\mathbf{1}} \times \mathbf{R}_{\mathbf{0}}, \frac{\partial \mathbf{R}_{\mathbf{1 2}}}{\partial u}=\mathbf{B}_{\mathbf{1}} \times \mathbf{R}_{12}, \frac{\partial \mathbf{R}_{\mathbf{1 3}}}{\partial u}=\mathbf{B}_{\mathbf{1}} \times \mathbf{R}_{13}  \tag{44}\\
& \frac{\partial \mathbf{R}_{\mathbf{0}}}{\partial v}=\mathbf{B}_{\mathbf{2}} \times \mathbf{R}_{\mathbf{0}}, \frac{\partial \mathbf{R}_{\mathbf{2 2}}}{\partial v}=\mathbf{B}_{\mathbf{2}} \times \mathbf{R}_{\mathbf{2 2}}, \frac{\partial \mathbf{R}_{\mathbf{2 3}}}{\partial v}=\mathbf{B}_{\mathbf{2}} \times \mathbf{R}_{\mathbf{2 3}}
\end{align*}
$$

If we insert (41) into (44), we get (43).

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Theorem 4.4. If we use space-like paremeter ruled surfaces, we can write below form:

$$
\begin{align*}
& \mathbf{R}_{12} \frac{\partial \mathbf{R}_{22}}{\partial S_{1}}=-\mathbf{R}_{22} \frac{\partial \mathbf{R}_{12}}{\partial S_{1}}=-\frac{(\sqrt{E})_{v}}{\sqrt{E G}} \\
& \mathbf{R}_{22} \frac{\partial \mathbf{R}_{12}}{\partial S_{2}}=-\mathbf{R}_{12} \frac{\partial \mathbf{R}_{22}}{\partial S_{2}}=\frac{(\sqrt{G})_{u}}{\sqrt{E G}} \tag{45}
\end{align*}
$$

Proof. From (22) and by taking derivatives of $(\sqrt{E})^{2}=-\mathbf{R}_{u}^{2},(\sqrt{G})^{2}=\mathbf{R}_{v}^{2}$, we have

$$
\begin{equation*}
\sqrt{E}(\sqrt{E})_{v}=-\mathbf{R}_{u} \mathbf{R}_{u v}, \sqrt{G}(\sqrt{G})_{u}=\mathbf{R}_{v} \mathbf{R}_{v u} \tag{46}
\end{equation*}
$$

and taking derivatives of (27) and (28)

$$
\begin{equation*}
\frac{\partial \mathbf{R}_{12}}{\partial v}=\frac{\mathbf{R}_{u v} \sqrt{E}-(\sqrt{E})_{v} \mathbf{R}_{u}}{E}, \quad \frac{\partial \mathbf{R}_{22}}{\partial u}=\frac{\mathbf{R}_{u v} \sqrt{G}-(\sqrt{G})_{u} \mathbf{R}_{v}}{G} \tag{47}
\end{equation*}
$$

are obtained respectively. Then, from the relations above (27), (28), (46) and (20).

$$
\begin{align*}
& \mathbf{R}_{12} \frac{\partial \mathbf{R}_{22}}{\partial u}=\frac{\mathbf{R}_{u}}{\sqrt{E}} \frac{\partial \mathbf{R}_{22}}{\partial u}=\frac{\mathbf{R}_{u} \mathbf{R}_{u v}}{\sqrt{E G}}=-\frac{(\sqrt{E})_{v}}{\sqrt{G}} \\
& \mathbf{R}_{22} \frac{\partial \mathbf{R}_{12}}{\partial v}=\frac{\mathbf{R}_{v}}{\sqrt{G}} \frac{\partial \mathbf{R}_{12}}{\partial v}=\frac{\mathbf{R}_{v} \mathbf{R}_{u v}}{\sqrt{E G}}=\frac{(\sqrt{G})_{u}}{\sqrt{E}} \tag{48}
\end{align*}
$$

are faund. If we use (23) in (48), we have

$$
\begin{align*}
& \mathbf{R}_{12} \frac{\partial \mathbf{R}_{22}}{\partial S_{1}}=R_{12} \frac{\partial \mathbf{R}_{22}}{\partial u} \frac{1}{\sqrt{E}}=-\frac{(\sqrt{E})_{v}}{\sqrt{E G}}  \tag{49}\\
& \mathbf{R}_{22} \frac{\partial \mathbf{R}_{12}}{\partial S_{2}}=R_{22} \frac{\partial \mathbf{R}_{22}}{\partial v} \frac{1}{\sqrt{G}}=\frac{(\sqrt{G})_{u}}{\sqrt{E G}} \\
& \mathbf{R}_{12} \frac{\partial \mathbf{R}_{22}}{\partial S_{1}}=-R_{22} \frac{\partial \mathbf{R}_{12}}{\partial S_{1}}  \tag{50}\\
& \mathbf{R}_{22} \frac{\partial \mathbf{R}_{22}}{\partial S_{2}}=-R_{12} \frac{\partial \mathbf{R}_{22}}{\partial S_{2}}
\end{align*}
$$

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Result 3. There is following relation between $\mathbf{R}_{12}$ and $\mathbf{R}_{\mathbf{2 2}}$ :

$$
\begin{equation*}
\mathbf{R}_{12} \cdot d \mathbf{R}_{22}=-d \mathbf{R}_{12} \cdot \mathbf{R}_{22}=-\frac{(\sqrt{E})_{v}}{\sqrt{G}} d u-\frac{(\sqrt{G})_{u}}{\sqrt{E}} d v \tag{51}
\end{equation*}
$$

Proof. If we differentiate (31) and consider (23) and (45), we obtain

$$
\begin{align*}
& -\mathbf{R}_{12} \cdot d \mathbf{R}_{22}=\mathbf{R}_{12} d \mathbf{R}_{22}=\mathbf{R}_{12}\left(\frac{\partial \mathbf{R}_{22}}{\partial S_{1}} d S_{1}+\frac{\partial \mathbf{R}_{22}}{\partial S_{2}} d S_{2}\right) \\
& =\mathbf{R}_{12} \frac{\partial \mathbf{R}_{22}}{\partial S_{1}} \sqrt{E} d u+\mathbf{R}_{12} \frac{\partial \mathbf{R}_{22}}{\partial S_{2}} \sqrt{G} d v  \tag{52}\\
& =-\frac{(\sqrt{E})_{v}}{\sqrt{E G}} \sqrt{E} d u-\frac{(\sqrt{G})_{u}}{\sqrt{E G}} \sqrt{G} d v
\end{align*}
$$

Thus, we have the result as below.

Result 4. There is following relation between $\mathbf{R}_{12}$ and $\mathbf{R}_{22}$ :

$$
\begin{equation*}
\frac{\partial}{\partial v}\left(\mathbf{R}_{12} \cdot \frac{\partial \mathbf{R}_{22}}{\partial u}\right)-\frac{\partial}{\partial u}\left(\mathbf{R}_{12} \cdot \frac{\partial \mathbf{R}_{22}}{\partial v}\right)=-\frac{\partial}{\partial v}\left(\frac{(\sqrt{E})_{v}}{\sqrt{G}}\right)+\frac{\partial}{\partial u}\left(\frac{(\sqrt{G})_{u}}{\sqrt{E}}\right) \tag{53}
\end{equation*}
$$

Proof. Taking derivative of (31) with respect to the parameters $v$ and $u$

$$
\begin{equation*}
\mathbf{R}_{12} \cdot \frac{\partial \mathbf{R}_{22}}{\partial v}=-\mathbf{R}_{22} \cdot \frac{\partial \mathbf{R}_{12}}{\partial v}=-\frac{(\sqrt{G})_{u}}{\sqrt{E}}, R_{12} \frac{\partial R_{22}}{\partial u}=-\frac{(\sqrt{E})_{v}}{\sqrt{G}} \tag{54}
\end{equation*}
$$

is written. Then, if we consider (48), (53) is obtain.

Theorem 4.5. There are following relations for the magnitudes $Q_{1}, Q_{1}$ and $Q$ of the space-like ruled surfaces $\mathbf{R}_{\mathbf{1 1}}, \mathbf{R}_{\mathbf{2 1}}$ and $\mathbf{R}_{\mathbf{1}}$, respectively.

$$
\begin{align*}
Q_{1} & =-\frac{(\sqrt{E})_{v}}{\sqrt{G}} \quad, \quad Q_{2}=-\frac{(\sqrt{G})_{u}}{\sqrt{E}}  \tag{55}\\
Q^{2} & =-\operatorname{sh}^{2} \Phi\left(\Phi^{\prime}-Q_{2}\right)^{2}+\operatorname{ch}^{2} \Phi\left(\Phi^{\prime}-Q_{1}\right)^{2} \tag{56}
\end{align*}
$$

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Proof. If we consider relations (43) and (23), we have

$$
\frac{\partial \mathbf{R}_{22}}{\partial S_{1}}=\frac{\partial \mathbf{R}_{22}}{\sqrt{E}}=\frac{1}{\sqrt{E}} \mathbf{B}_{1} \times \mathbf{R}_{22}
$$

Then, using the relation above and (45), (32) and (25), we may write

$$
-\frac{(\sqrt{E})_{v}}{\sqrt{E G}}=\mathbf{R}_{12} \cdot \frac{\partial \mathbf{R}_{22}}{\partial S_{1}}+\mathbf{R}_{12} \frac{\mathbf{B}_{1} \times \mathbf{R}_{\mathbf{1 2}}}{\sqrt{E}}=-\frac{\mathbf{R}_{12} \times \mathbf{R}_{22}}{\sqrt{E}} \cdot \mathbf{B}_{1}=-\frac{\mathbf{R}_{0} \cdot \mathbf{B}_{1}}{\sqrt{E}}=\frac{Q_{1}}{\sqrt{E}}
$$

By the same way and from the relations (45), (23), (44), (32) and (25), we have

$$
\frac{(\sqrt{G})_{u}}{\sqrt{E G}}=-\mathbf{R}_{12} \cdot \frac{\partial \mathbf{R}_{22}}{\partial S_{2}}=-\frac{1}{\sqrt{G}} \mathbf{R}_{12} \cdot \frac{\mathbf{R}_{22}}{\partial v}=-\frac{\mathbf{R}_{12} \cdot\left(\mathbf{B}_{2} \times \mathbf{R}_{22}\right)}{\sqrt{G}}=\frac{\mathbf{R}_{0} \cdot \mathbf{B}_{2}}{\sqrt{G}}=-\frac{Q_{2}}{\sqrt{G}}
$$

Finally, if we take derivative of the relation (38) by using the derivative formulas (24) and consider (41) and (31), it can be reached (55) and (56).

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